

Diffractive physics at ALICE, LHC.

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Physics motivation – Mesons not allowed by CQM

- Constituent Quark Model (CQM)
 - Goes back over 40 years to Gell-Mann and Zweig
 - Constituent quarks: Quasi particles with additional effective mass due to interaction with gluon field.
- Mesons in constituent quark model
 - Color-singlet $|q\bar{q}\rangle$ states, groups into $SU(N)_{flavour}$ multiplets
 - Table of mesons with quantum numbers

ℓ	S	J	$P=(-1)^{1+\ell}$	I	$G=(-1)^{I+\ell+S}$	$C=(-1)^{\ell+S}$	$I^G J^{PC}$	nomenclature
0	0	0	-1	0	+1	+1	0^+0^{-+}	η
0	0	0	-1	1	-1	+1	1^-0^{-+}	π
0	1	1	-1	0	-1	-1	0^-1^{--}	ω
0	1	1	-1	1	+1	-1	1^+1^{--}	ρ
1	1	0	+1	0	+1	+1	0^+0^{++}	f_0
1	1	0	+1	1	-1	+1	1^-0^{++}	a_0
1	0	1	+1	0	-1	-1	0^-1^{++}	h_1
1	0	1	+1	1	+1	-1	1^+1^{+-}	b_1
1	1	1	+1	0	+1	+1	0^+1^{++}	f_1
1	1	1	+1	1	-1	+1	1^-1^{++}	a_1
1	1	2	+1	0	+1	+1	0^+2^{++}	f_2
1	1	2	+1	1	-1	+1	1^-2^{++}	a_2
2	0	2	-1	0	+1	+1	0^+2^{-+}	η_2
2	0	2	-1	1	-1	+1	1^-2^{-+}	π_2

- Not allowed J^{PC} with CQM mesons : $0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, 3^{-+}, \dots$ (spin-exotic)





Purposes of this study

- Finding states beyond Constituent Quark Model (CQM)
 - Physical mesons = linear superposition of all allowed basis states: $|q\bar{q}\rangle, |q\bar{q}g\rangle, |gg\rangle, |q^2\bar{q}^2\rangle, \dots$
 - Quarkonia  $|q\bar{q}\rangle$
 - Tetra-quarks : is not my sight  $|q^2\bar{q}^2\rangle$
 - Hybrids : The first purpose of this study  $|q\bar{q}g\rangle$
 - Glueballs : The second purpose  $|gg\rangle$

Hybrids $|q\bar{q}g\rangle$

- Resonances with excited glue and quarks.
- Angular momentum of glue component \rightarrow **all J^{PC} possible**
- Lightest predicted hybrid : **spin-exotic $J^{PC} = 1^{-+}$**
 - ✓ Mass 1.3 to 2.2 GeV/c²
 - ✓ Experimental candidates $\pi_1(1400, 1500, 2000)$ controversial

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Glueballs $|gg\rangle$

- Bound states consisting purely of gluons
- Lightest predicted glueball : ordinary $J^{PC} = 0^{++}$
 - ✓ Will strongly mix with nearby conventional $J^{PC} = 0^{++}$
 - ✓ Mass 1.5 to 2.0 GeV/c²
 - ✓ Experimental candidate $f_0(1500)$; glueball interpretation disputed



Systems for the study?

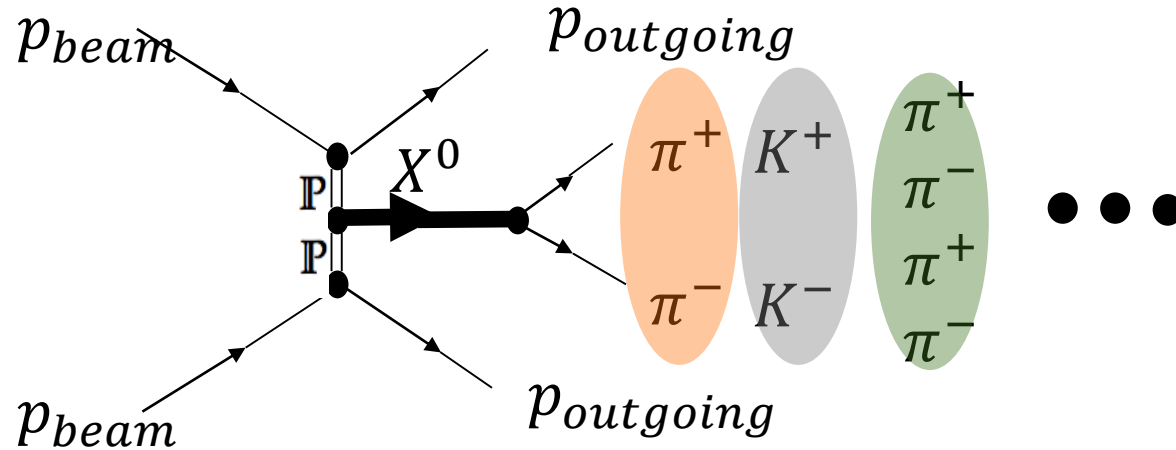
- Gluon-rich environment is needed
 - Gluons play a key role for Hybrids & glueballs
 - More possibility to search them with gluon-rich environment

→ Centrally diffracted system(CD)

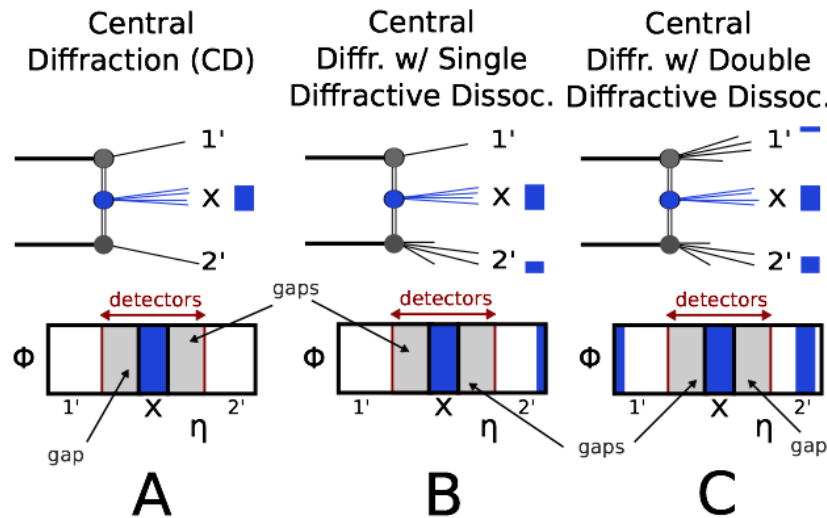
- Central diffraction system
 - **Central diffraction system is caused by two \mathbb{P} (pomeron) exchange**
 - Pomeron is also quasi-particle(at least $\{gg\}$ (first order) in and around proton)
 - Gluon-rich environment
 - More possible to find a glueball and hybrid
 - The best thing is we can **exclude many quarkonia from the system**
 - $J^{PC} = \text{even}^{++}(J^{PC})$ states are only possible with two \mathbb{P} exchange.
 - Thus many quantum states cannot survive in CD system
 - We can search glueballs or hybrids without backgrounds

Central diffraction and double-gap topology

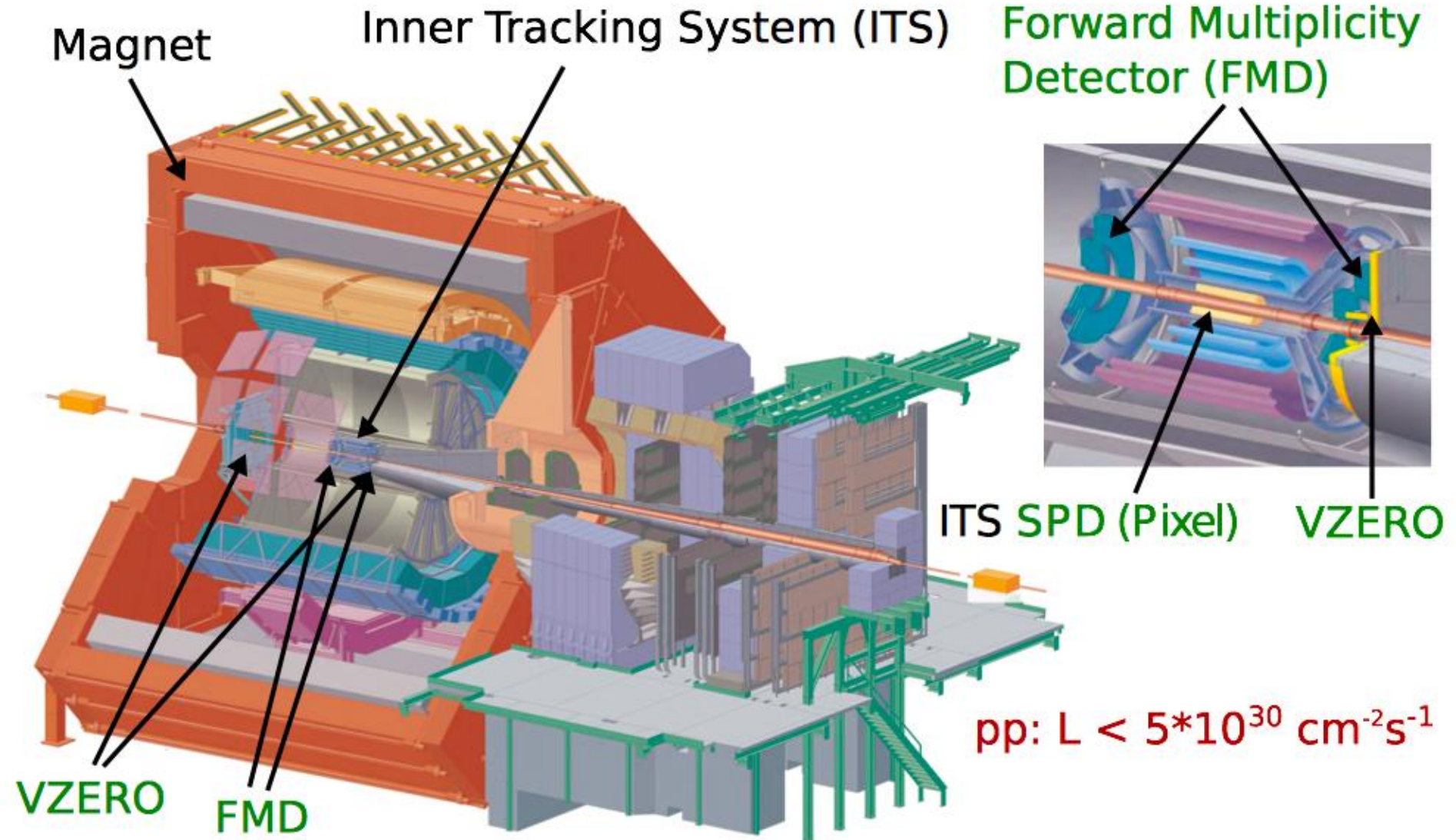
- Schematic view of the Central Diffraction (**double \mathbb{P} exchange**)



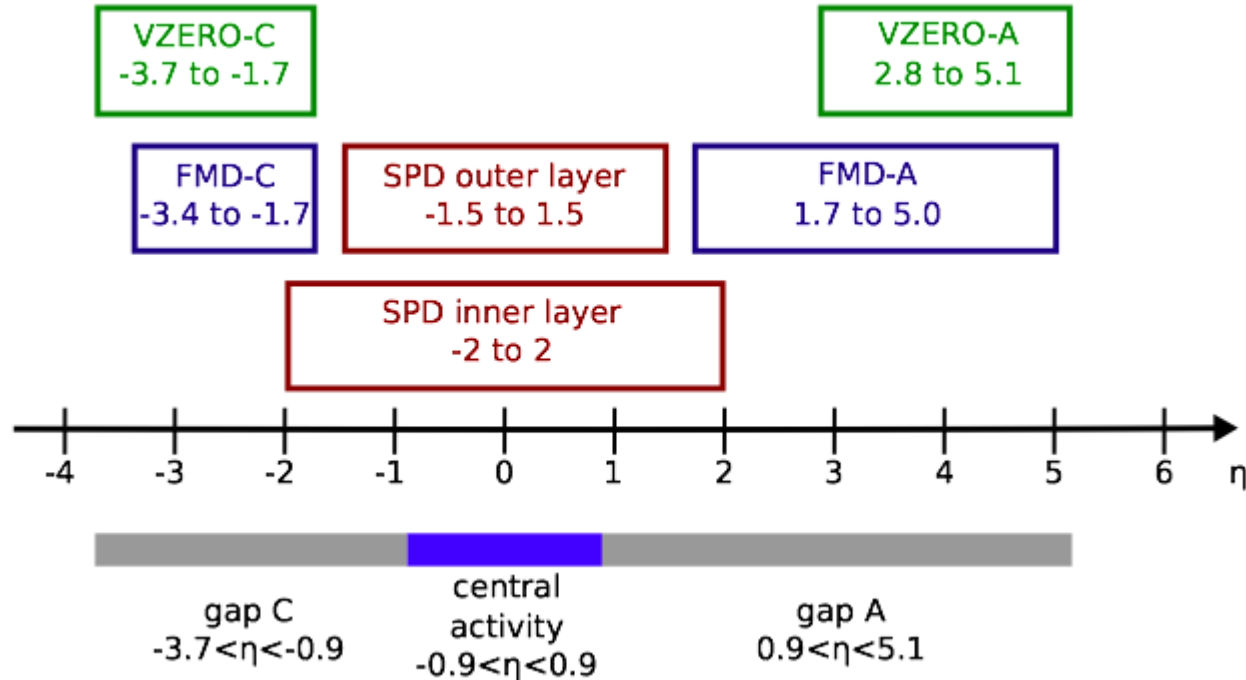
- Of particular interest is double-gap topology **as filter for central diff(CD)**



A Large Ion Collider Experiment

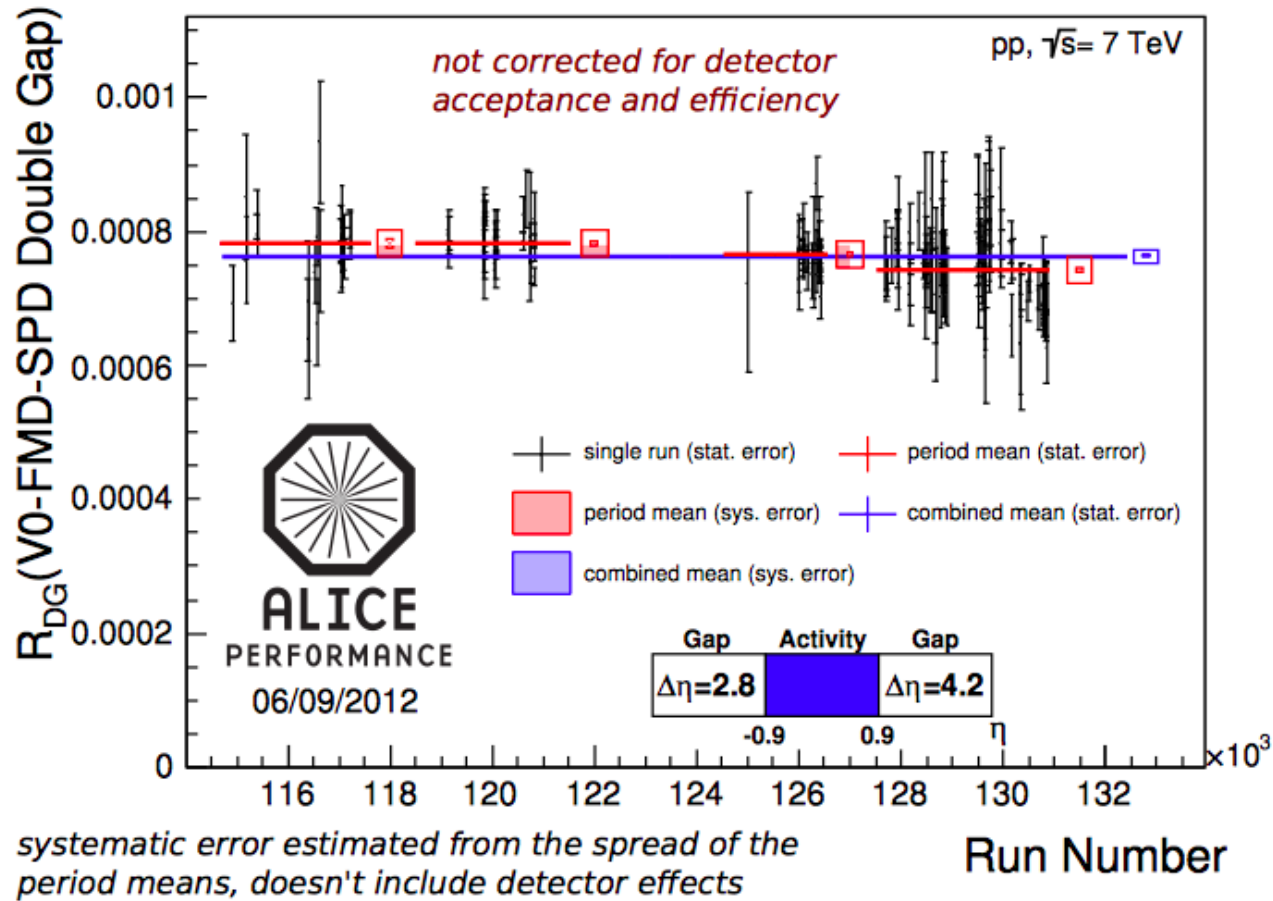


Double-Gap topology in ALICE



- Selection criteria for the double-gap topology
 - Activity in the central barrel ($-0.9 < \eta < 0.9$) \rightarrow detector : SPD
 - Two gaps (no activity) outside of central barrel
 - A-side : $0.9 < \eta < 5.1$, C-side : $-3.7 < \eta < -0.9$
 - Detectors : VZERO, FMD, SPD requiring no activity(no tracks and signals)

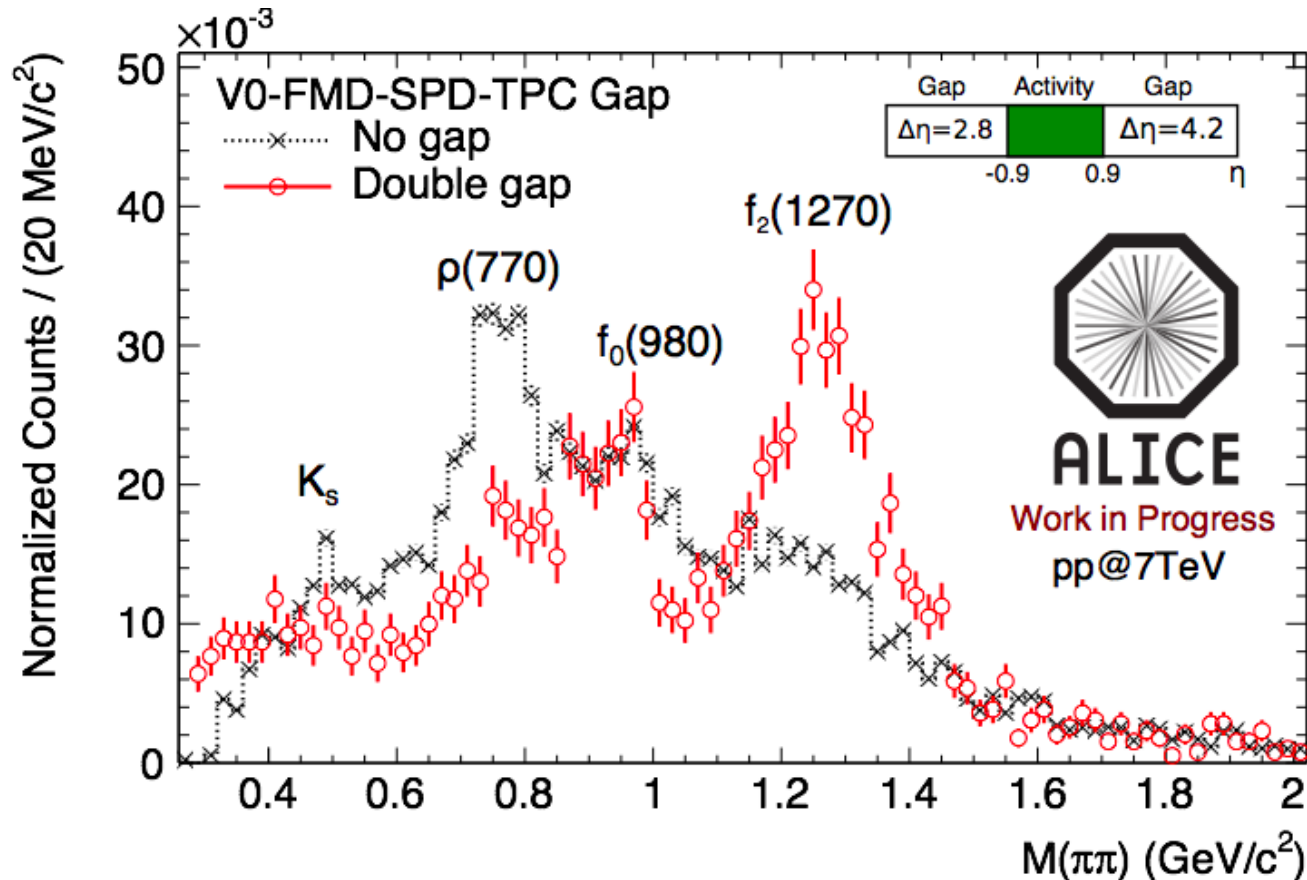
Double-Gap fraction



- $R_{DG}(2.8, \pm 0.9, 4.2) = \frac{N_{DG}}{N_{MB}} = (7.63 \pm 0.02_{stat} \pm 0.87_{sys}) \times 10^{-4}$
 - Uniform behavior at ALICE, stable with various run-periods.

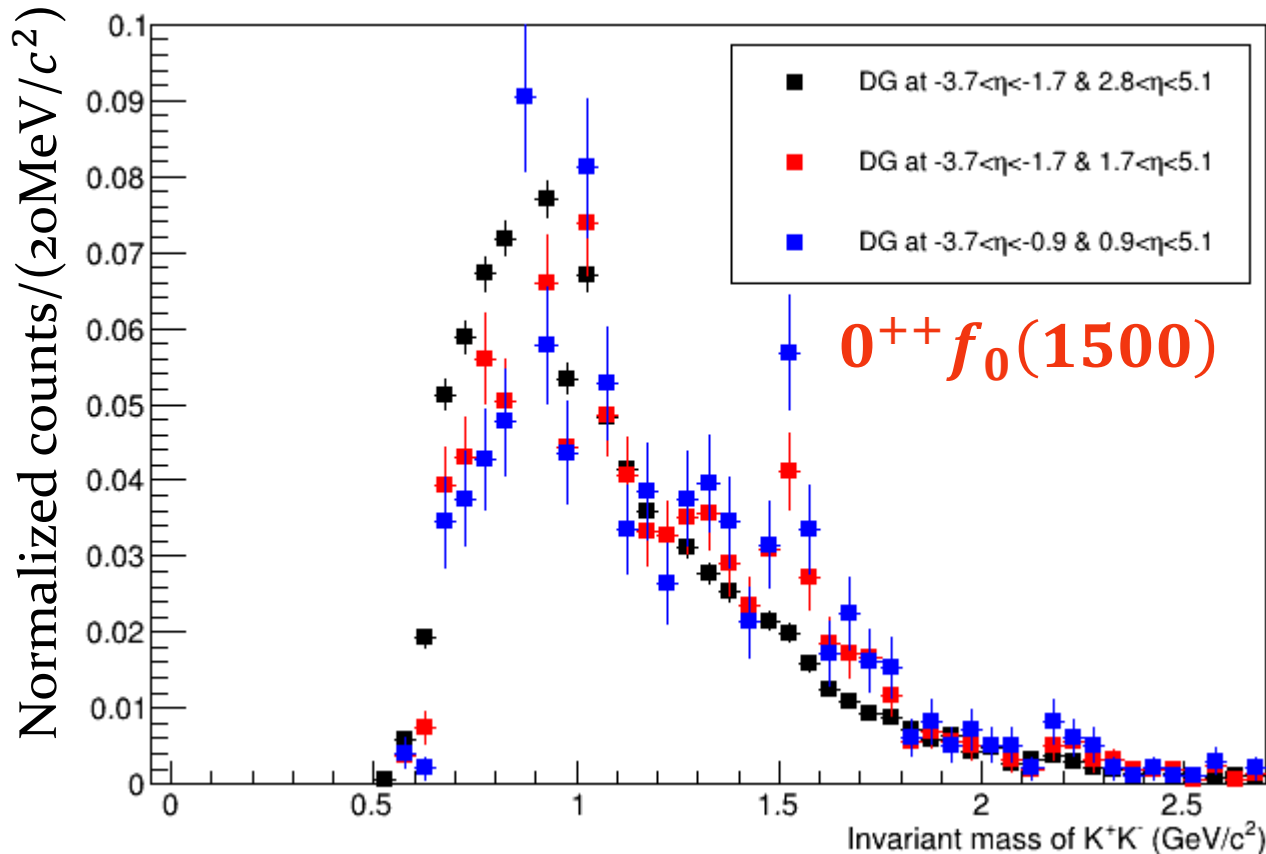
Are we correctly taking CD events by DG topology?

- CD (**two IP exchange**) system confines quantum numbers
 - Double **IP** exchange only allows quantum numbers, $J^{PC} = \text{even}^{++} (0^{++}, 2^{++}, \dots)$
 - With double-gap condition
 - $1^{--} \rho_0(770)$ was restricted, $0^{++} f_0(980)$ and $2^{++} f_2(1270)$ appear significantly



Results $X^0 \rightarrow K^+ K^-$

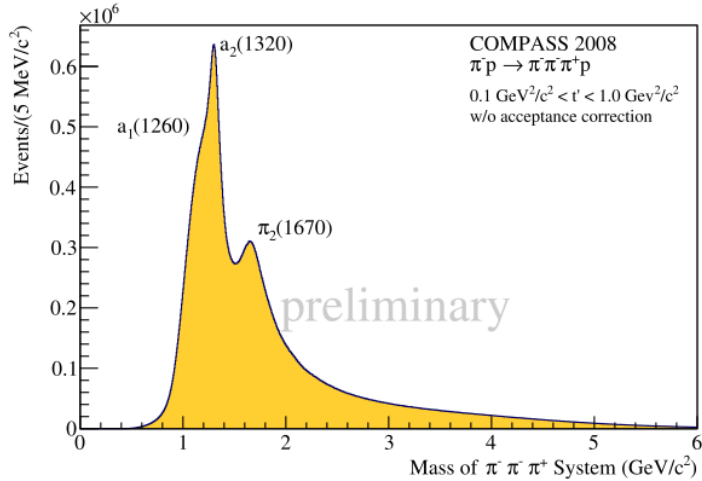
- Full statistics at 7TeV
- Clear $f_0(1500)$ (*gluball candidate*) or $f_2'(1525)$ and $f_0(1710)$
 - Need to study about $J=0$ and p and $c=+$ (by partial wave fit)
- $f_2(1270)$, $f_0(1370)$ also can be seen however cannot be distinguished.



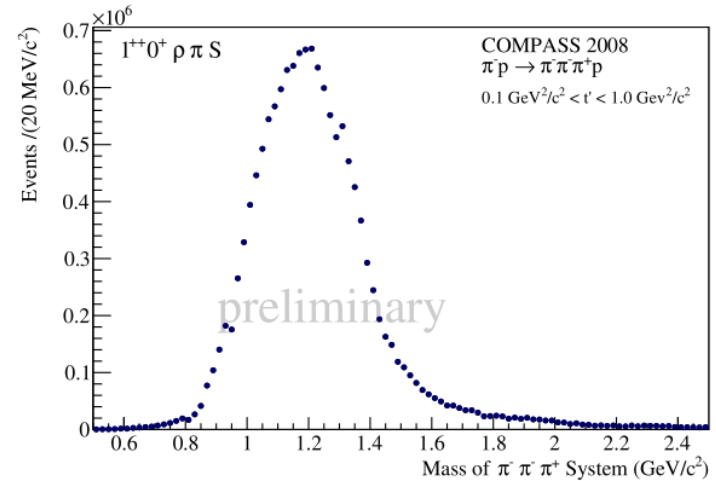
Stronger
Double-gap condition

COMPASS measurement for hybrid $1^{-+}(\pi_0(1600))$

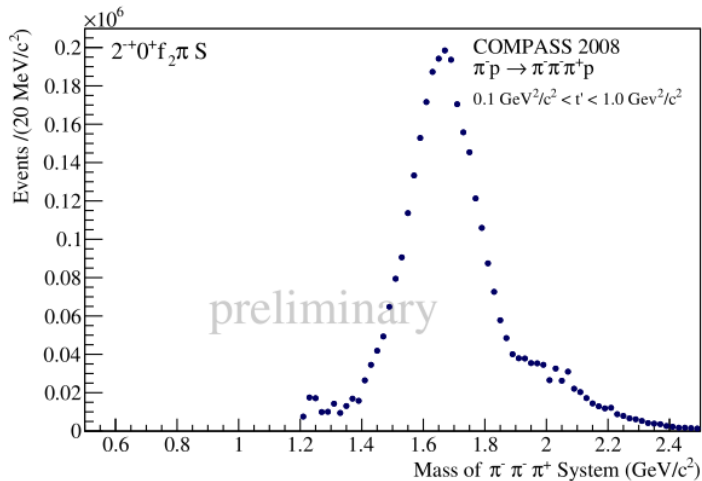
$\pi^- \pi^+ \pi^-$ invariant mass spectrum



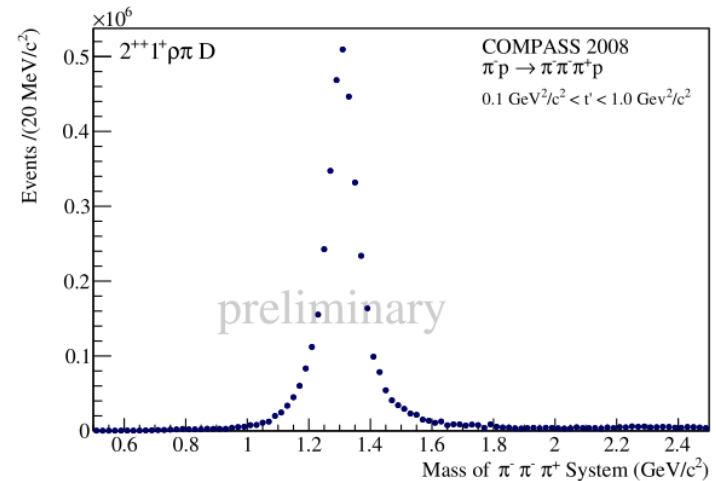
$1^{++} 0^+ [\rho\pi] S: a_1(1260)$



$2^{-+} 0^+ [f_2\pi] S: \pi_2(1670)$

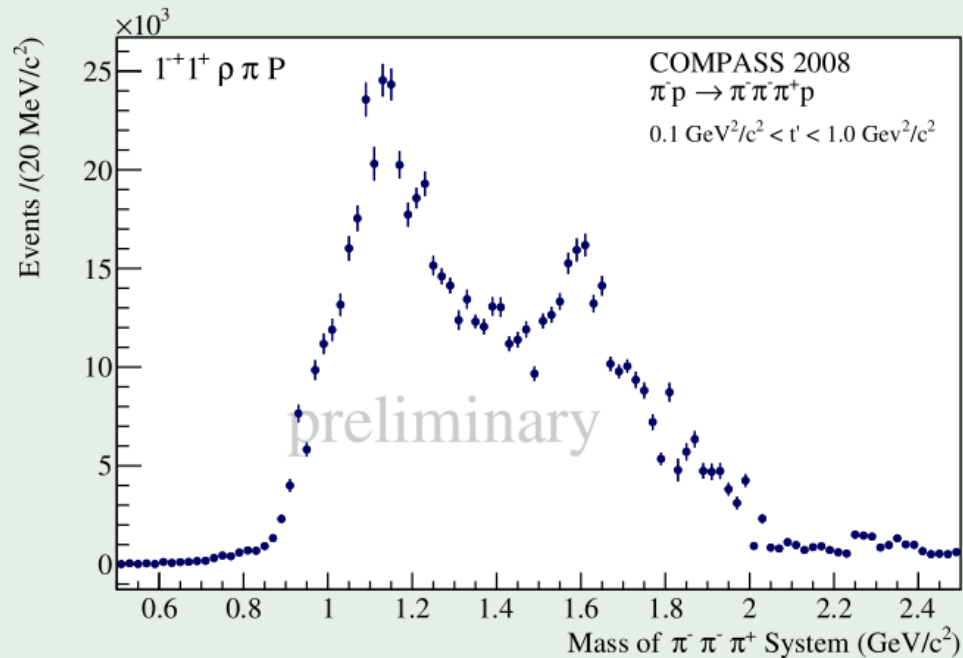


$2^{++} 1^+ [\rho\pi] D: a_2(1320)$



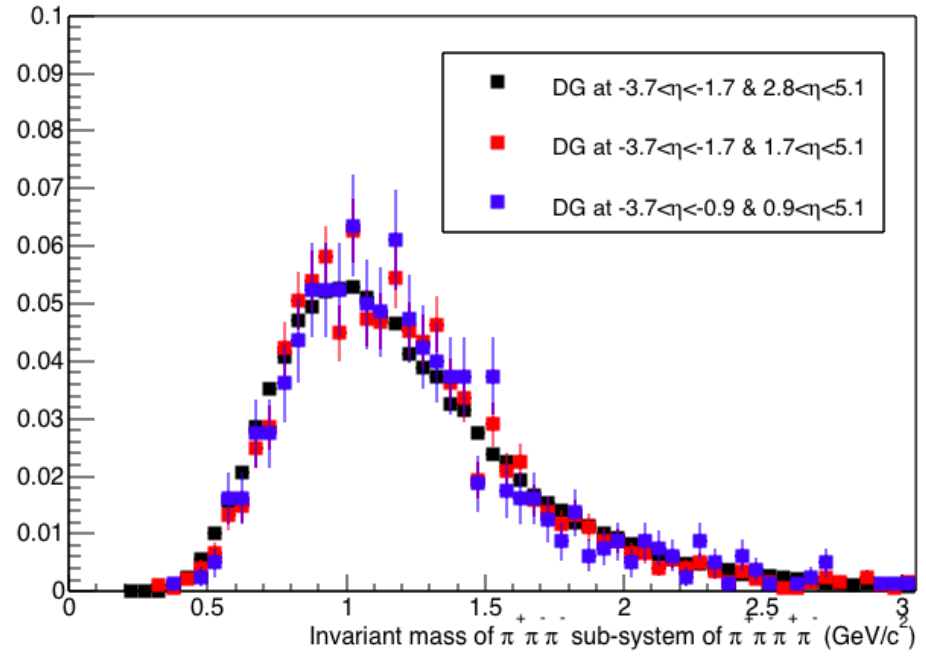
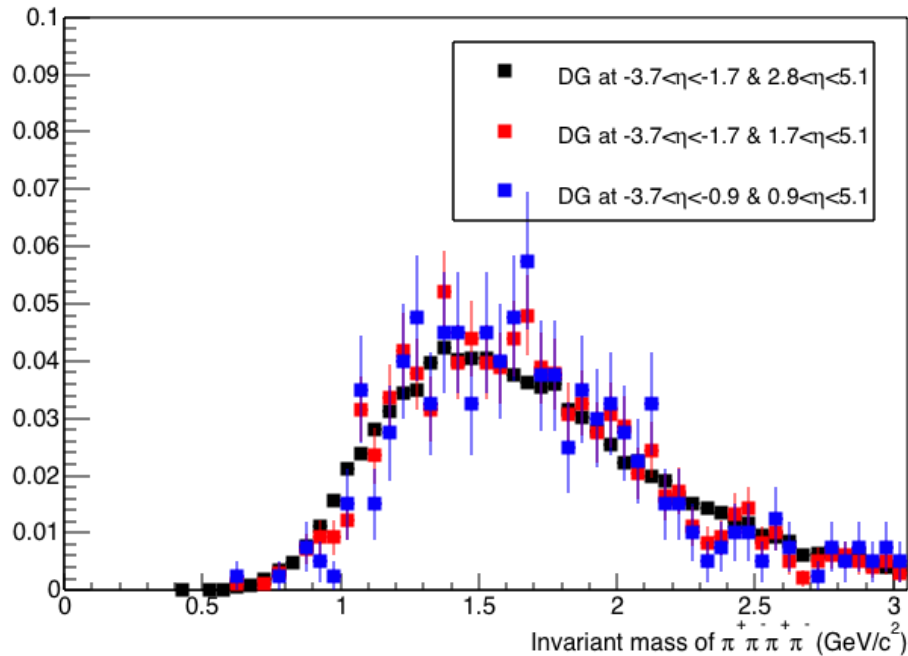
COMPASS measurement for hybrid 1^{-+} ($\pi_0(1600)$)

COMPASS



- 190 GeV/c π beam
- p target
- $50 \cdot 10^6$ events
- $0.1 < t' < 1.0$ (GeV/c)²
- Rank-2 fit with 53 waves

4 Pion events and 3 pion sub-system at ALICE



- We are expecting to find a hybrid 3 pion sub-system of 4 pion events from $m_{DG}(1^{--} \rightarrow \pi^+\pi^-\pi^-)$
- Developing “partial wave analysis” which was used for the COMPASS result in order to pull out signals from data
- Analysis are ongoing for glueball and hybrids.



Summary and promising points for the physics

- We found that we can choose central production by control of the double rapidity gap condition. ($f_0(980)$, $f_2(1270)$ showed up as we expected)
 - Search for glueball& hybrids is ongoing. Remaining jobs are just “fitting”
 - At RUN2(2015~) period, pp collisions with $\sqrt{s} = 13$ to 14 TeV are reserved . We expect that there would be more statistics with central production since cross-section of the double pomeron exchange goes up exponentially with energy.
 - With the help of the forward new detectors, ADA & ADC, we can measure the system more efficiently with extended pseudo-rapidity coverage
 - ALICE has the capability to measure low p_T particles compared to other LHC experiment(ATLAS, CMS) which is very important for this study
- Promising subject for RUN2, RUN3 with ALICE



Residuals



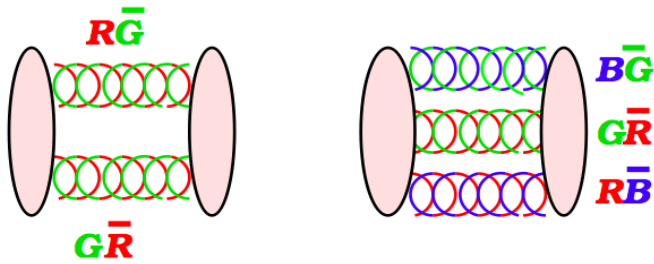
Physics motivation

- Exotic meson : mesons which have quantum numbers not possible in the quark model
 - QCD suggests : in addition to mesons made by quarks, **new states** where colored gluons play an essential role.
 - New state 1 : “glueballs” resonating states purely made by gluons
 - New state 2 : “Hybrids” resonating states made by gluons and quarks.
 - This study : search for this glueballs or Hybrids
- Possible states by J^{PC} expression for the constituent model of $q\bar{q}$ mesons
$$0^{-+}, 0^{++}, 1^{--}, 1^{+-}, 1^{--}, 2^{--}, 2^{-+}, 2^{++}, 3^{--}, 3^{+-}, 3^{--}, \dots$$
- Not possible states
$$0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, 3^{-+}, \dots$$
- The latter quantum numbers are known as “*explicitly exotic*”. If this states are found, it implies that it should not be a normal $q\bar{q}$ meson

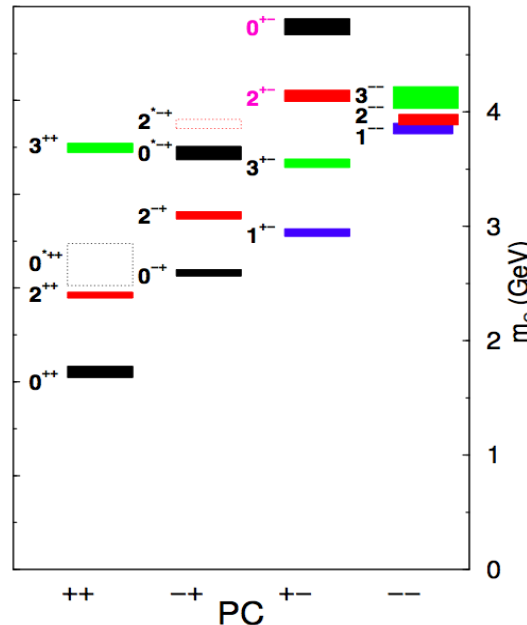
Glueballs in the exotic mesons

- Gluons

- Because gluons carry color charge, it is possible to bind into **color-singlet objects**
- The simplest glueballs are either two or three gluons confined together.
- The best prediction of the glueball spectrum comes from the lattice QCD



Two and three gluons bound into color singlet glueballs

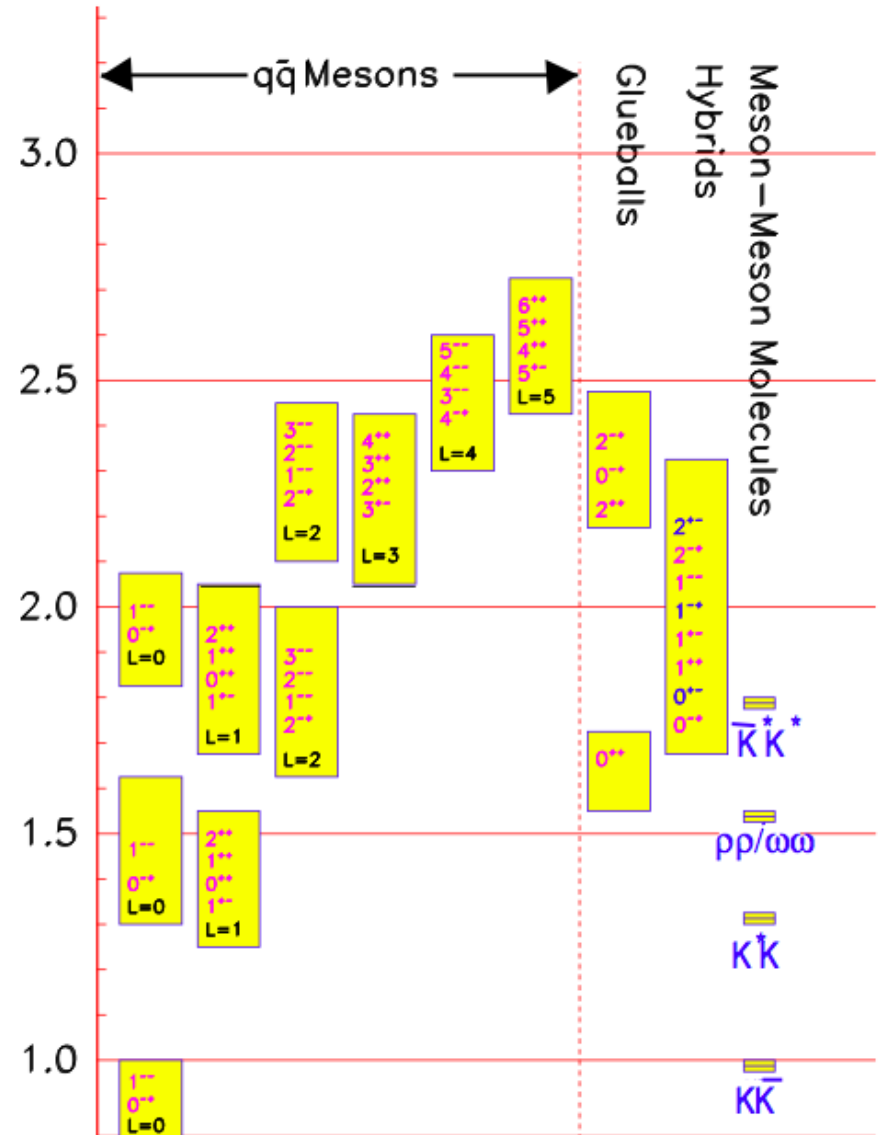


Glueball states which are predicted by L-QCD. The states described by black means that this states are also possible by $q\bar{q}$ meson. On the other hand, pink states are not possible by $q\bar{q}$

- The lightest glueball is expected to have $J^{PC} = 0^{++}$ followed by a 2^{++} and then a 0^{-+} . Unfortunately all these 3 states are allowed for the normal meson[$q\bar{q}$]
- The lightest glueball states with exotic are beyond $4\text{GeV}/c^2$. This exceeds our experimental mass regime.

Glueballs in this experiment

- By double-gap condition
 - $J^{PC} = \text{even}^{++}$ is only allowed
 - 0^{++} and 2^{++}
 - Two possible states of glueball
 - For 2^{++} , it is difficult to discriminate it from $q\bar{q}$ mesons
 - Target of this experiment
 - 0^{++} glueball



Mass independent partial wave fit

Components of the LogLikelihood function:

$$\ln L = \sum_{n=1}^{N_{\text{events}}} \ln \sum_{\epsilon} \sum_{i,j} \rho_{ij}^{\epsilon} \bar{\psi}_i^{\epsilon}(\tau_n) \bar{\psi}_j^{\epsilon}(\tau_n)^* - \sum_{\epsilon} \sum_{i,j} \rho_{ij}^{\epsilon} I A_{ij}^{\epsilon}$$

Decay amplitudes (points to ρ_{ij}^{ϵ})
 Kinematics (points to τ_n)
 Acceptance corrected Phase space integral (points to $I A_{ij}^{\epsilon}$)
 Spin density matrix (fit parameters) (points to ρ_{ij}^{ϵ})
 Coherent sum over waves (points to $\bar{\psi}_i^{\epsilon}(\tau_n) \bar{\psi}_j^{\epsilon}(\tau_n)^*$)
 Incoherent sum over reflectivities (points to \sum_{ϵ})

Production amplitudes \rightarrow Spin density matrix:

$$\rho_{ij}^{\epsilon} = \sum_r T_{ir}^{\epsilon} T_{jr}^{\epsilon*}$$

Normalized decay amplitudes:

$$\bar{\psi}_i^{\epsilon}(\tau) = \frac{\psi_i^{\epsilon}(\tau)}{\sqrt{\int |\psi_i^{\epsilon}(\tau')|^2 d\tau'}}$$

Phase space integrals (with acceptance):

$$I A_{ij}^{\epsilon} = \int \bar{\psi}_i^{\epsilon}(\tau_n) \bar{\psi}_j^{\epsilon}(\tau_n)^* \text{Acc}(\tau) d\tau$$

$$\text{Acc}(\tau) = \begin{cases} 0 \\ 1 \end{cases}$$