

# Search for $B \rightarrow X_s \nu \bar{\nu}$ decay at Belle II experiment

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#### **Belle II Experiment**

- Electrons and positrons are accelerated up to 7 GeV and 4 GeV respectively by SuperKEKB accelerator
- Its energy correspond to the resonance of  $\Upsilon(4S)$  which mainly decays into B meson pair
- Belle II detector consists of several sub-detector components



#### **Belle II Experiment**

- Entire information of initial states is known in Belle II experiment.
  - $\rightarrow$  Belle II has an advantage on the decay modes with invisible particles, like neutrinos
- $B \qquad e^+ \qquad B \qquad e^+$

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• The target integrated luminosity:  $50 ab^{-1}$ 

• Currently, 575.5  $fb^{-1}$  data is recorded. The target sample is 364  $fb^{-1} \Upsilon(4S)$  on-resonance data for this analysis



### Motivation

- $B \rightarrow X_s \nu \bar{\nu}$  decay
  - Flavour-changing neutral currents process
  - BR =  $(2.9 \pm 0.3) \times 10^{-5}$  at SM [JHEPO2(2015)184]
- This decay can give a clue for the new physics
  - Ieptoquark [PhysRevD.95.035027]
  - invisible light scalar [Eur. Phys. J. C (2017) 77: 650]
  - fermion dark matter [arXiv:2405.06742]
- Previous study
  - $UL(B \to X_s \nu \bar{\nu}) = 6.4 \times 10^{-4}$  (90% CL) by ALEPH<sup>†</sup> [Eur.Phys.J.C 19 (2001) 213-227]
    - <sup>D</sup> There is no  $B \to X_S \nu \bar{\nu}$  results from Belle or Belle II because  $\nu \bar{\nu}$  makes it challenging
  - $BR(B^+ \to K^+ \nu \bar{\nu}) = (2.3 \pm 0.7) \times 10^{-5}$  by BELLE II <sup>‡</sup> [PhysRevD.109.112006]



‡ combined result for inclusive and hadronic tagging analysis





# Analysis Strategy

- One side of B meson is hadronically reconstructed
  - Hierarchical reconstruction technique is used [Comput Softw Big Sci 3, 6 (2019)]
  - This B meson is called as tag side B meson, B<sub>tag</sub>
- For signal side B meson, sum of exclusive method is used
  - 30 decay modes are reconstructed for X<sub>s</sub>

	$B^0 \bar{B}^0$			$B^{\pm}$		
$\overline{K}$	$K_S^0$			$K^{\pm}$		
$K\pi$	$K^{\pm}\pi^{\mp}$	$K^0_S \pi^0$		$K^{\pm}\pi^0$	$K^0_S \pi^{\pm}$	
$K2\pi$	$K^{\pm}\pi^{\mp}\pi^{0}$	$K^0_S \pi^{\pm} \pi^{\mp}$	$K^0_S\pi^0\pi^0$	$K^{\pm}\pi^{\mp}\pi^{\pm}$	$K^0_S \pi^{\pm} \pi^0$	$K^{\pm}\pi^0\pi^0$
$K3\pi$	$K^{\pm}\pi^{\mp}\pi^{\pm}\pi^{\mp}$	$K^0_S \pi^{\pm} \pi^{\mp} \pi^0$	$K^{\pm}\pi^{\mp}\pi^{0}\pi^{0}$	$K^{\pm}\pi^{\mp}\pi^{\pm}\pi^{0}$	$K^0_S \pi^{\pm} \pi^{\mp} \pi^{\pm}$	$K^0_S \pi^\pm \pi^0 \pi^0$
$K4\pi$	$K^{\pm}\pi^{\mp}\pi^{\pm}\pi^{\mp}\pi$	${}^{0}K^{0}_{S}\pi^{\pm}\pi^{\mp}\pi^{\pm}\pi^{\pm}$	$^{\mp}K^0_S\pi^{\pm}\pi^{\mp}\pi^0\pi^0$	$K^{\pm}\pi^{\mp}\pi^{\pm}\pi^{\mp}\pi$	${}^{\pm}\!K^0_S \pi^{\pm} \pi^{\mp} \pi^{\pm} \pi^0$	$K^{\pm}\pi^{\mp}\pi^{\pm}\pi^{0}\pi^{0}$
3K	$K^{\pm}K^{\mp}K^0_S$			$K^{\pm}K^{\mp}K^{\pm}$		
$3K\pi$	$ K^{\pm}K^{\mp}K^{\pm}\pi^{\mp}$	$K^{\pm}K^{\mp}K^0_S\pi^0$		$K^{\pm}K^{\mp}K^{\pm}\pi^{0}$	$K^0_S K^\pm K^\mp \pi^\pm$	

• It covers ~92% of entire  $X_s$  sample, with assuming  $K^0$  equally decays into  $K_s^0$  or  $K_L^0$ 



#### **Event Generation**

- We produce  $B \to K \nu \bar{\nu}$ ,  $B \to K^* \nu \bar{\nu}$ , and non-resonant  $B \to X_s \nu \bar{\nu}$  MC sample separately
- For  $B \to K \nu \bar{\nu}$  and  $B \to K^* \nu \bar{\nu}$  Monte-Carlo (MC) samples, the SM form factors are used [Phys.Rev.D 107 (2023) 1, 014511] [JHEP02(2015)184] [JHEP02(2015)184]

Decay amplitude for  $B \to K \nu \bar{\nu}$ :  $\mathcal{M}(b \to s \nu \bar{\nu}) \propto f_+(s) \left\{ (p_B + p)_\mu - \frac{m_B^2 - m_K^2}{s} q_\mu \right\} (\bar{\nu} \gamma^\mu (1 - \gamma_5) \nu)$ 

Decay amplitude for  $B \to K^* \nu \bar{\nu}$ :  $\mathcal{M}(b \to s \nu \bar{\nu}) \propto \mathcal{T}_{\mu} \left( \bar{\nu} \gamma^{\mu} (1 - \gamma_5) \nu \right)$ 

$$\mathcal{T}_{\mu} = (m_B + m_{K^*})A_1(s)\epsilon_{\mu}^* - A_2(q^2)\frac{\epsilon^* \cdot q}{m_B + m_{K^*}}(p + p_{K^*})_{\mu} + i\frac{2V(s)}{m_B + m_{K^*}}\epsilon_{\mu\nu\rho\sigma}\epsilon^{*\nu}p^{\rho}p_{K^*}^{\sigma}$$





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#### **Event Generation**

- Non-resonant  $B \to X_s \nu \bar{\nu}$  MC sample is produced by the following distribution  $(\mu EPO4(2009)022)$  $\frac{d\Gamma}{dq^2} \propto \sqrt{\lambda(1, \hat{m}_s^2, s_b)} \left[ 3s_b(1 + \hat{m}_s^2 - s_b - 4\hat{m}_s) + \lambda(1, \hat{m}_s^2, s_b) \right]$
- To determine the mass of  $X_s$ , Fermi motion model is adopted [PhysRevD.55.4105]
  - In Fermi motion model, quarks have some velocity inside B meson
  - Also,  $M_{X_s} > 1.1 \text{ GeV/c}^2$  is required for non-resonant  $X_s \nu \bar{\nu}$  MC sample
- Hadronization is done by PYTHIA



- Preselection
  - To reduce combinatorial backgrounds
  - reconstructed  $X_s$  mass,  $M_{X_s}^{\text{reco}} < 2.0 \text{ GeV/c}^2$

#### To reject backgrounds with additional particles, outside of $B_{tag}X_s$ candidates

- the number of tracks <sup>‡</sup> = 0
- the number of  $\pi^0$  candidates = 0
- the number of  $K_S^0$  candidates = 0





- Several selections are applied
  - $M_{bc}^{\text{tag}} > 5.27 \text{ GeV/c}^2$
  - $|\Delta E^{tag}| < 0.2 \text{ GeV}$
  - $E_{ecl} < 1.3 \, {\rm GeV}$

 $\approx E_{ecl}$ : remaining energy deposited in the calorimeter 9



 $X M_{X_c}^{reco}$ : mass of reconstructed  $X_s$  candidate

schematic view of Belle II detector

- Boosted Decision Tree (BDT) is used as multivariate analysis (MVA) method
- 32 variables are used
  - Variables are selected based on discriminant power and goodness of data/MC agreement on sideband
  - Example of powerful variables: E<sub>ecl</sub> and event shape variable [Nucl. Phys. B 149, 413 (1979)]



- BDT has several hyperparameters
  - Grid search is done to achieve high performance and low overtraining

hyper parameter	tested values	selected value
nTrees	100, 500, 1000, 1500, 2000	2000
$\operatorname{depth}$	2,3,4	3
$\operatorname{shrinkage}$	0.05,0.1,0.15,0.2	0.05
subsample	0.3,0.4,0.5,0.6,0.7	0.5
binning	6,7,8,9	6

- AUC and BDT output is checked
  - AUC value for training sample: 0.966
  - AUC value for test sample: 0.964



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- BDT output > 0.86 is applied
  - Because it is used as a fitting variable, deliberately loose cut is applied
  - Figure of merit  $\left(=\frac{S}{\sqrt{S+B}}\right)$  is used to select this criteria



• Finally, signal efficiency is about 0.11%

\*  $M_{X_c}^{\text{reco}}$ : mass of reconstructed  $X_s$  candidate

# Corrections - $B \rightarrow K^{(*)} n \overline{n}$

- $B \rightarrow K^{(*)} n \overline{n}$  mismodeling correction
  - $B \to K^{(*)} n \bar{n}$  MC sample is not modeled well, even though it can mimic  $B \to K^{(*)} \nu \bar{\nu}$
  - Produce  $B \to K^{(*)}n\bar{n}$  MC with flat  $M_{n\bar{n}}$  distribution and reweight it!
  - use  $B \to K^{(*)} p \bar{p}$  study result to obtain  $M_{n\bar{n}}$  distribution for  $B \to K^{(*)} n \bar{n}$  decay
  - use exponential function or  $2^{nd}$  order polynomial to get a form of  $M_{n\bar{n}}$  distribution



# Corrections - $B \rightarrow F$

- $B \rightarrow K K_L^0 K_L^0$  mismodeling correction
  - $B \rightarrow K K_L^0 K_L^0$  MC sample is not modeled well
  - need to correct Dalitz plot
  - use  $B \rightarrow KK_S^0K_S^0$  study result to obtain PDF [Physical Review D 85.11 (2012): 112010] [Physical Review D 85.5 (2012): 054023]

$$F_{j}^{L}(s_{12}, s_{23}) = R_{j}(m) X_{L}(|\overrightarrow{p}^{*}|r') T_{j}(L, \overrightarrow{p}, \overrightarrow{q})$$
  
$$F_{j}^{L}(s_{\min}, s_{\max}, L) = R_{j}(m) X_{L}(|\overrightarrow{p}^{*}|r') X_{L}(|\overrightarrow{q}|r) T_{j}(L, \overrightarrow{p}, \overrightarrow{q})$$



2 4 6 8 10 12 14 16 18 20 22

2 4 6 8 10 12 14 16 18 20 2

-0.0

0.0

0.0 0.0 0.0

0.0

0.0

0.0

0.0

0.0

0.0

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#### Off-resonance sample

- Off-resonance
  - Off-resonance samples are analyzed for the correction and validation
- $M_{bc}^{tag}$  and  $\Delta E^{tag}$  are calculated as follows, when analyze the off-resonance sample:

$$M_{bc,corrected}^{tag} = \sqrt{\left(E_{beam}^{off}\right)^2 - \left(p_B^*\right)^2} + \left(E_{beam}^{on} - E_{beam}^{off}\right)$$

 $\Delta E^{tag} = E_B^* - E_{beam}^{off}$ 

• this correction is applied to use the same  $M_{bc}$  and  $\Delta E$  cut

### Off-resonance sample

#### • Off-resonance

Off-resonance samples are analyzed



- This result is used to correct the continuum MC sample in on-resonance
  - Overall normalization factor is obtained for each  $M_{X_s}^{reco}$  region
- Also, systematic uncertainty from the shape of variables is estimated by off-resonance sample

#### **Control Channel**

- $B \rightarrow X_s J/\psi$  analysis is also done
  - Analysis method is almost same
  - after reconstruction of  $B \rightarrow X_s J/\psi$  , ignore  $\mu$



# **Control Channel**

•  $B \rightarrow X_s J/\psi$  analysis is also done



- This result is used to obtain correction/systematic uncertainty
  - BDT Efficiency correction
  - Systematic uncertainty from the Efficiency for B<sub>tag</sub>

#### Systematics

Several systematics are estimated

some major sources:

#### **Background normalization**

- apply  $\pm 20\%$  uncertainties on each backgrounds
- motivated by the data/MC of sideband

#### **MC Statistics**

- comes from statistical uncertainty of MC sample

#### BR of main B meson decays

- uncertainty of BR which is obtained from PDG
- for major B decays, uncertainties are applied

#### $q\overline{q}$ background shape

- estimated by training another BDT with 60 MeV lower  $E_{beam}$  sample

$\sigma_{\!\mu}$ at each	$M_{X_s}^{\text{true}}$	region
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	mass region $[GeV]$		
Source	$0.0 < M_{X_s}^{\rm true} < 0.6$	$0.6 < M_{X_s}^{ m true} < 1.0$	$1.0 < M_{X_s}^{\rm true}$
Background normalization	$^{+1.19}_{-1.14}$	$^{+2.29}_{-2.13}$	$^{+4.01}_{-4.07}$
MC statistics	$^{+0.98}_{-0.77}$	$^{+1.68}_{-1.37}$	$^{+3.93}_{-3.31}$
BR of main B meson decays	$^{+0.30}_{-0.14}$	$^{+0.62}_{-0.49}$	$^{+1.28}_{-0.69}$
$q\bar{q}$ shape	$^{+0.29}_{-0.27}$	$^{+0.24}_{-0.24}$	$^{+0.71}_{-0.69}$
statistical uncertainty	$^{+1.78}_{-1.65}$	$+2.98 \\ -2.79$	$+5.94 \\ -5.68$

# binning and fitting range



• Probability density function is constructed from each bins

$$\mathcal{P} = \prod_{b \in \text{bins}} \text{Pois}(n_b | v_b) \cdot \prod_p f_p(a_p | \alpha_p)$$

Poisson distributionConstraint term for systematic uncertaintyfor each bin/channel(nuisance parameters)

- extended maximum likelihood fit is done
- Fitting parameter μ (signal strength):
   a factor relative to the SM expectation

# Toy MC study

- Toy MC study is done
  - Generate 10000 toy MC sample with fluctuating nuisance parameter
  - Fit and obtain the MINOS asymmetric error for each toys
  - Pull is calculated from the fitting result:



if (fit result)  $\leq$  (true value)

otherwise,

(true value)-(fit result)

(positive MINOS error)

(fit result) – (true value)

(negative MINOS error)

pull =

### Toy MC study

- Toy MC study is done
  - Signal strength through all  $M_{X_s}^{\text{true}}$  can be easily obtained:  $\mu = \frac{BR_1 \times \mu_1 + BR_2 \times \mu_2 + BR_3 \times \mu_3}{BR}$
  - Pull is calculated with the signal strength through all  $M_{X_s}^{true}$  region



### Linearity Test

- Linearity is done
  - do toy MC study for different input µ values
  - check the output µ distributions and median value is selected, because error is asymmetric
  - Some bias can be found. This fitter bias is included in systematic uncertainty



### Linearity Test

- Linearity is done
  - Signal strength through all  $M_{X_s}^{\text{true}}$  can be easily obtained:  $\mu = \frac{BR_1 \times \mu_1 + BR_2 \times \mu_2 + BR_3 \times \mu_3}{BR}$
  - The linearity test is done with the signal strength through all  $M_{X_s}^{\text{true}}$  region
  - Some bias can be found. This fitter bias is included in systematic uncertainty



# **Upper Limit**

- Upper limit of branching ratio is calculated by CLs method (with  $364 fb^{-1}$ )
  - MC sample is used to calculate the upper limit of the branching ratio



 $\underset{X_s}{\times} M_{X_s}^{\text{true}}$ : mass of  $X_s$  recorded by MC information

# **Upper Limit**

- Upper limit of branching ratio is calculated by CLs method (with  $364 fb^{-1}$ )
  - MC sample is used to calculate the upper limit of the branching ratio



 $\overset{\text{true}}{\times} M_{X_s}^{\text{true}}$ : mass of  $X_s$  recorded by MC information

#### Conclusion

- $B \rightarrow X_s \nu \overline{\nu}$  is interesting decay because it can give a clue for several new physics
  - Ieptoquark [PhysRevD.95.035027]
  - invisible light scalar [Eur. Phys. J. C (2017) 77: 650]
  - fermion dark matter [arXiv:2405.06742]
- Several selection, including MVA, are used to suppress background
- Background normalization and MC statistic are dominant systematic sources
- Most of analysis procedures are done
  - With 364 *f b*<sup>-1</sup>, MC expectation is

$$UL(B \to X_{s}\nu\bar{\nu}) = \begin{cases} 2.4 \times 10^{-5} & (0.0 < M_{X_{s}}^{\text{true}} < 0.6 \text{ GeV/c}^{2}), \\ 7.2 \times 10^{-5} & (0.6 < M_{X_{s}}^{\text{true}} < 1.0 \text{ GeV/c}^{2}), \\ 28.3 \times 10^{-5} & (1.0 \text{ GeV/c}^{2} < M_{X_{s}}^{\text{true}}). \end{cases}$$

$$UL(B \to X_s \nu \bar{\nu}) = 27.9 \times 10^{-5}$$
 (all  $M_{X_s}^{\text{true}}$  region) at 90% confidence level

#### Backup