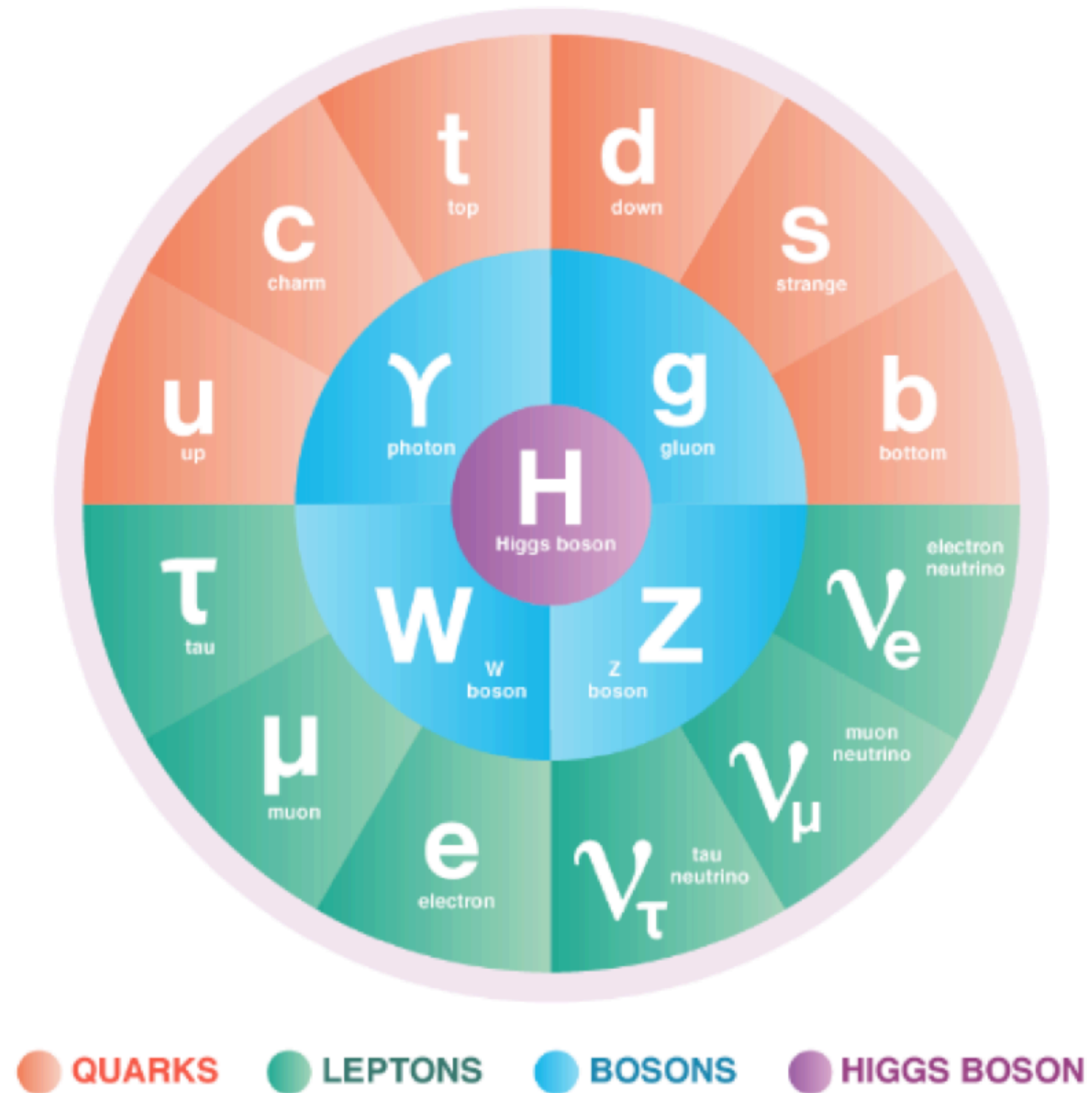


Physics of quarkonia at Belle & Belle II

Speaker: Junhao Yin

The standard model



We know 6 quarks & 6 leptons

FERMIONS

matter constituents
spin = 1/2, 3/2, 5/2, ...

Leptons spin = 1/2

Flavor	Mass GeV/c ²	Electric charge
ν_L lightest neutrino*	$(0-0.13)\times 10^{-9}$	0
e electron	0.000511	-1
ν_M middle neutrino*	$(0.009-0.13)\times 10^{-9}$	0
μ muon	0.106	-1
ν_H heaviest neutrino*	$(0.04-0.14)\times 10^{-9}$	0
τ tau	1.777	-1

Quarks spin = 1/2

Flavor	Approx. Mass GeV/c ²	Electric charge
u up	0.002	2/3
d down	0.005	-1/3
c charm	1.3	2/3
s strange	0.1	-1/3
t top	173	2/3
b bottom	4.2	-1/3

We know four types of interactions

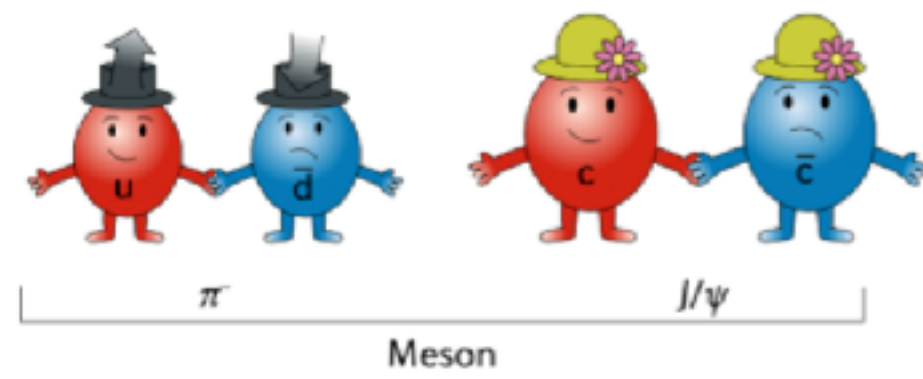
Properties of the Interactions				
The strengths of the interactions (forces) are shown relative to the strength of the electromagnetic force for two u quarks separated by the specified distances.				
Property	Gravitational Interaction	Weak Interaction (Electroweak)	Electromagnetic Interaction	Strong Interaction
Acts on:	Mass – Energy	Flavor	Electric Charge	Color Charge
Particles experiencing:	All	Quarks, Leptons	Electrically Charged	Quarks, Gluons
Particles mediating:	Graviton (not yet observed)	W^+ W^- Z^0	γ	Gluons
Strength at	10^{-18} m	10^{-41}	0.8	25
	3×10^{-17} m	10^{-41}	10^{-4}	60

- Gravity is responsible for the structure of the Universe
- Electromagnetic interaction → the molecules and atoms
- Weak interaction → the stars shine
- Strong interaction → the structure of the nuclei, nucleons, hadronic matters from the building blocks --- quarks!

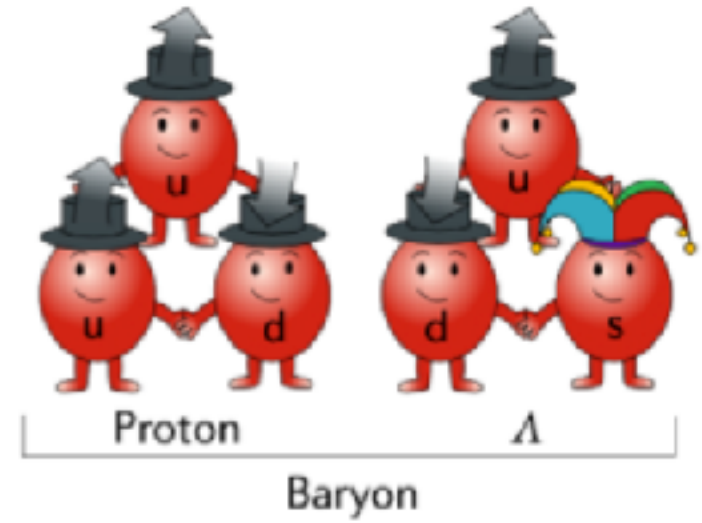
but HOW?

Hadrons

mesons



baryons

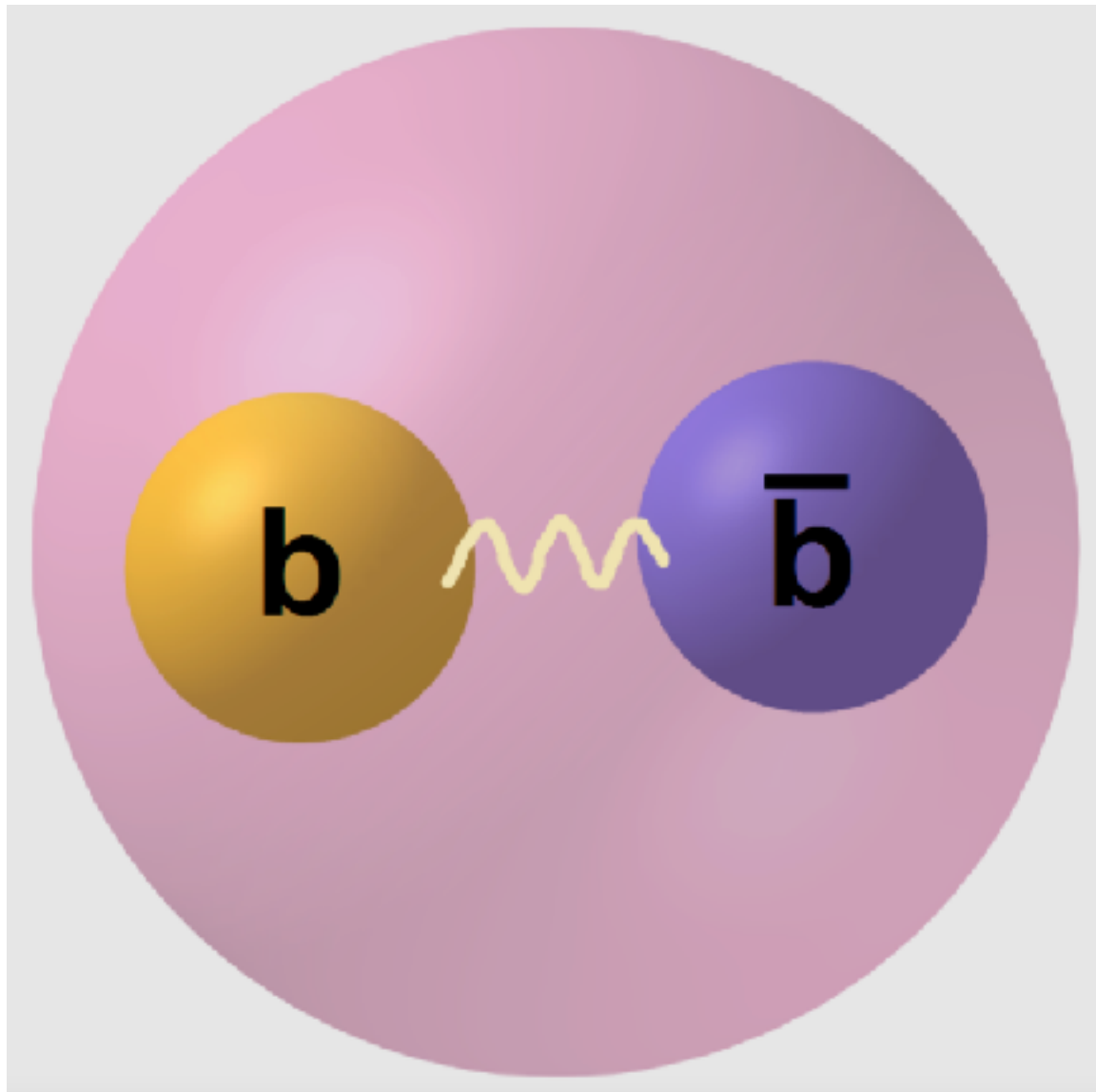
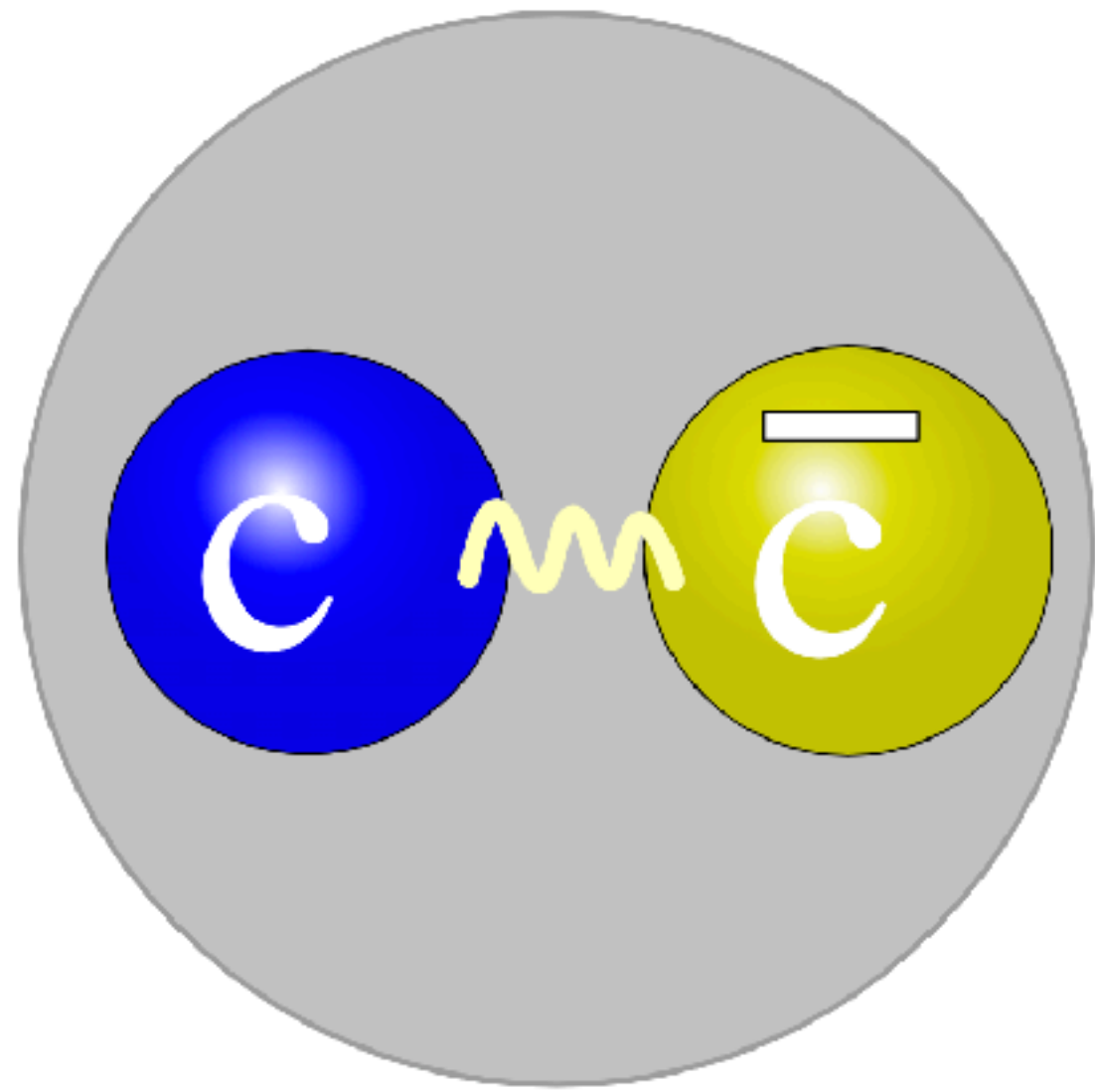


A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon b if we assign to the triplet t the following properties: spin $\frac{1}{2}$, $z = -\frac{1}{3}$, and baryon number $\frac{1}{3}$. We then refer to the members $u^{\frac{2}{3}}$, $d^{-\frac{1}{3}}$, and $s^{-\frac{1}{3}}$ of the triplet as "quarks" q and the members of the anti-triplet as anti-quarks \bar{q} . **Baryons** can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while **mesons** are made out of $(q\bar{q})$, $(qq\bar{q}\bar{q})$, etc. It is assuming that the lowest baryon configuration (qqq) gives just the representations 1, 8, and 10 that have been observed, while the lowest meson configuration $(q\bar{q})$ similarly gives just 1 and 8.



M. Gell-Mann

Physics Letters 8, 214 (1964);



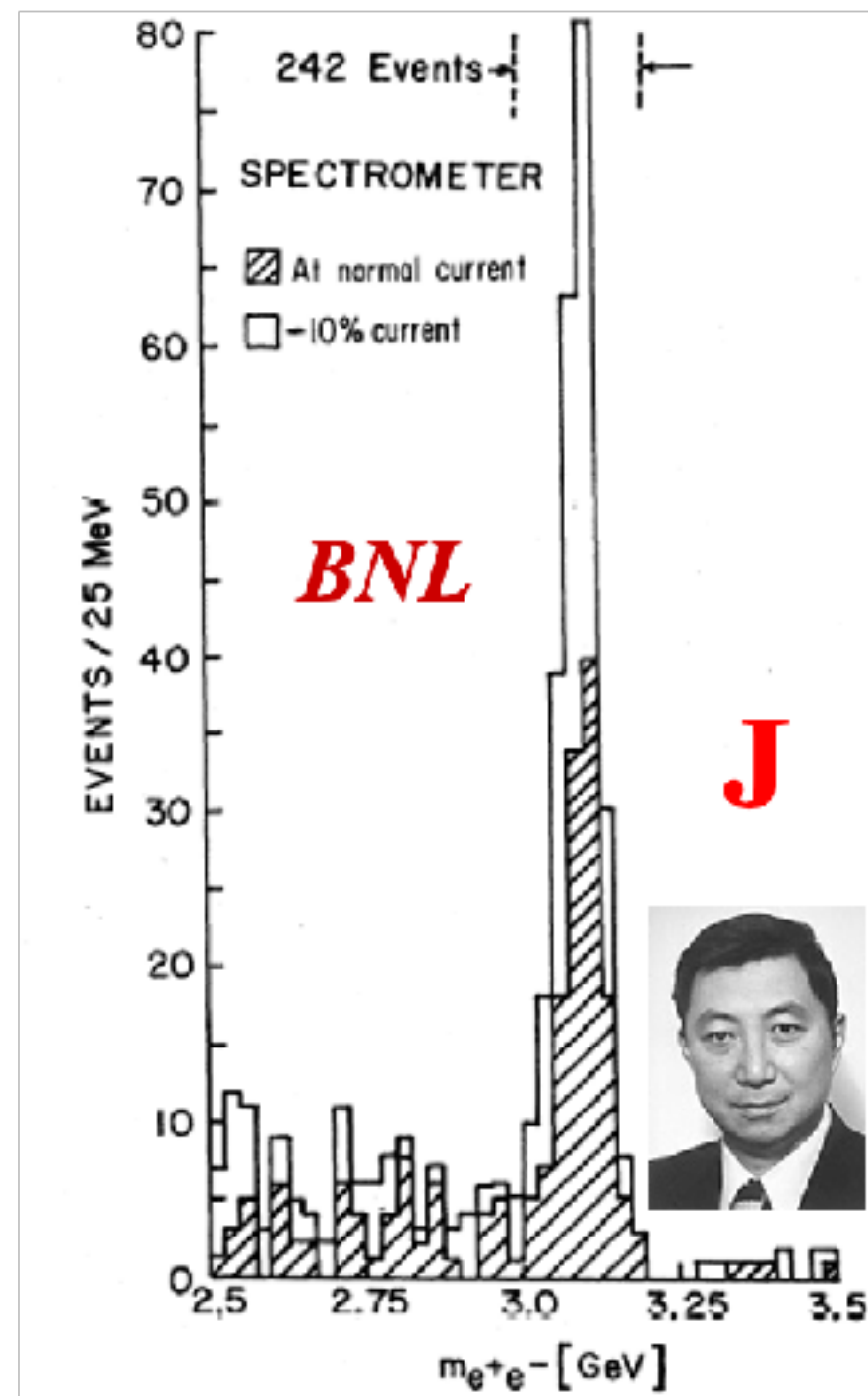
The flavor-changing decays would be strongly suppressed if there were a fourth quark (now called the *charm quark*) that was a complementary counterpart to the *strange quark*.

Discovery of the J/psi

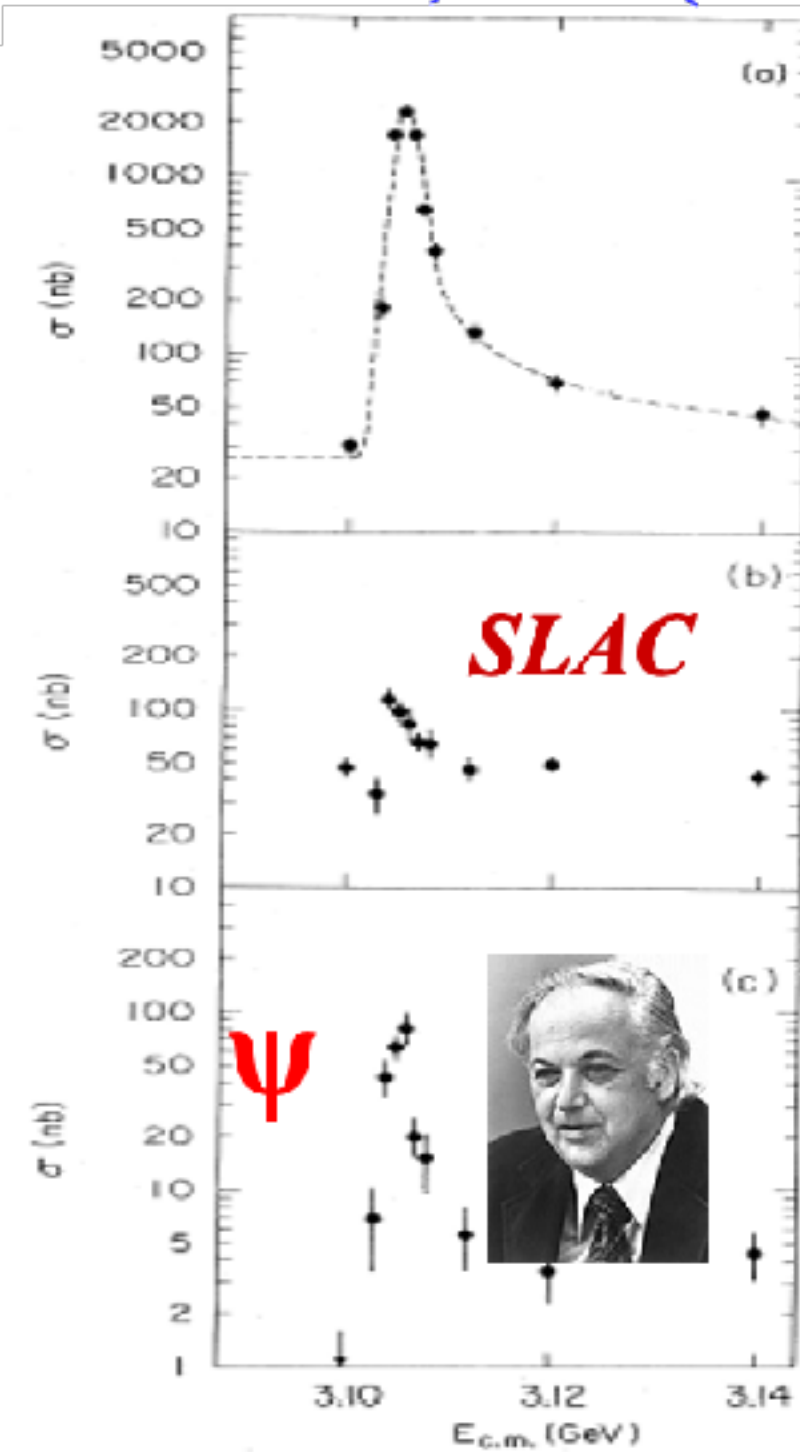
“November Revolution of Particle Physics!”

Charm quark was proposed in 1964, first application in 1970!

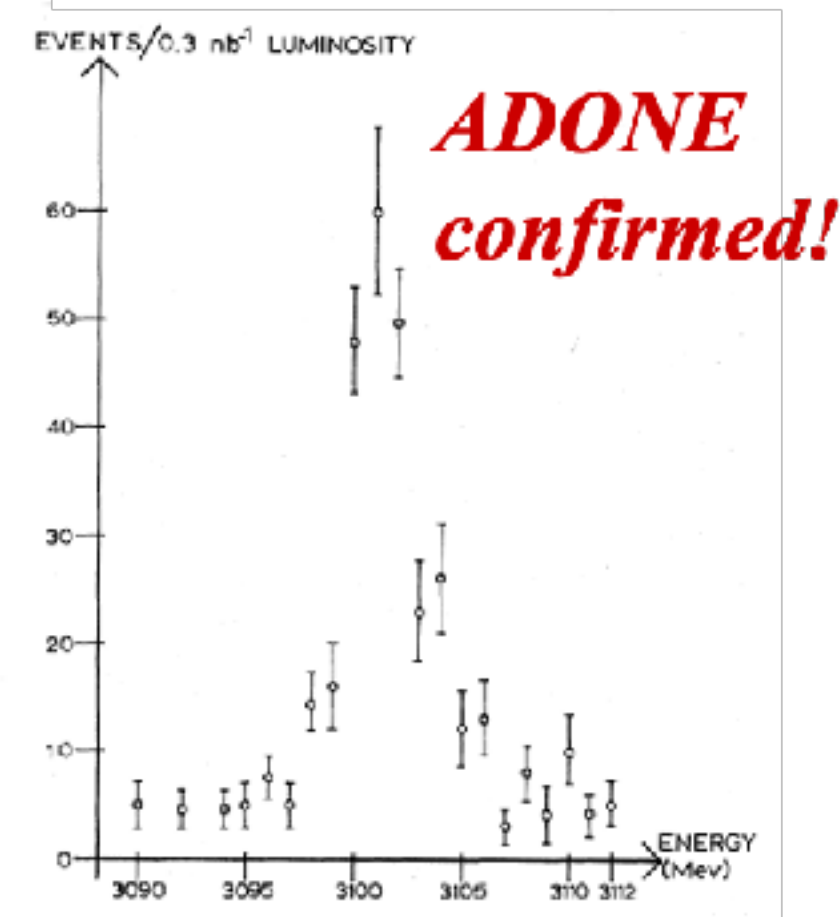
PRL33, 1404 (1974)



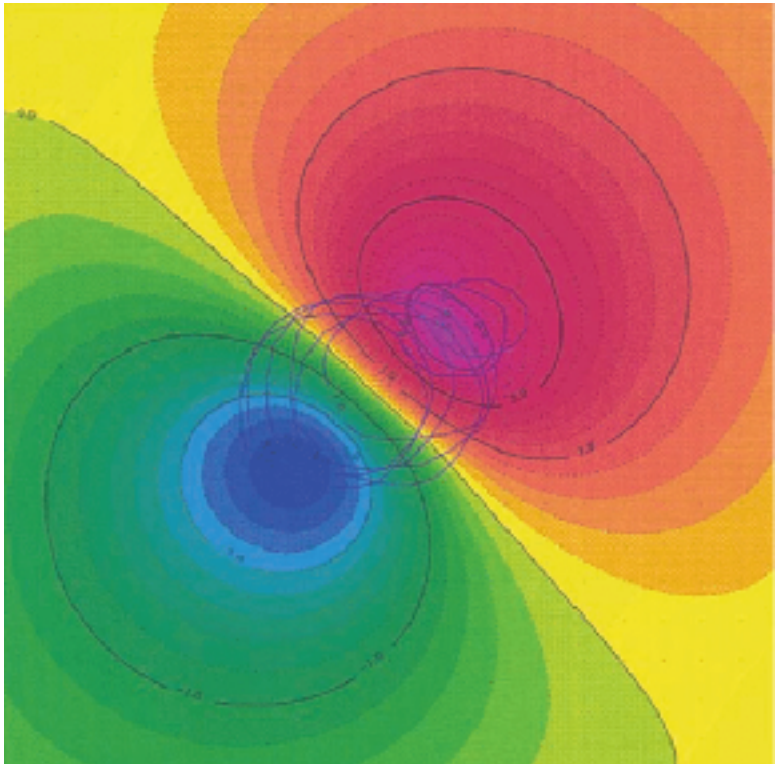
PRL33, 1406 (1974)



PRL33, 1408 (1974)

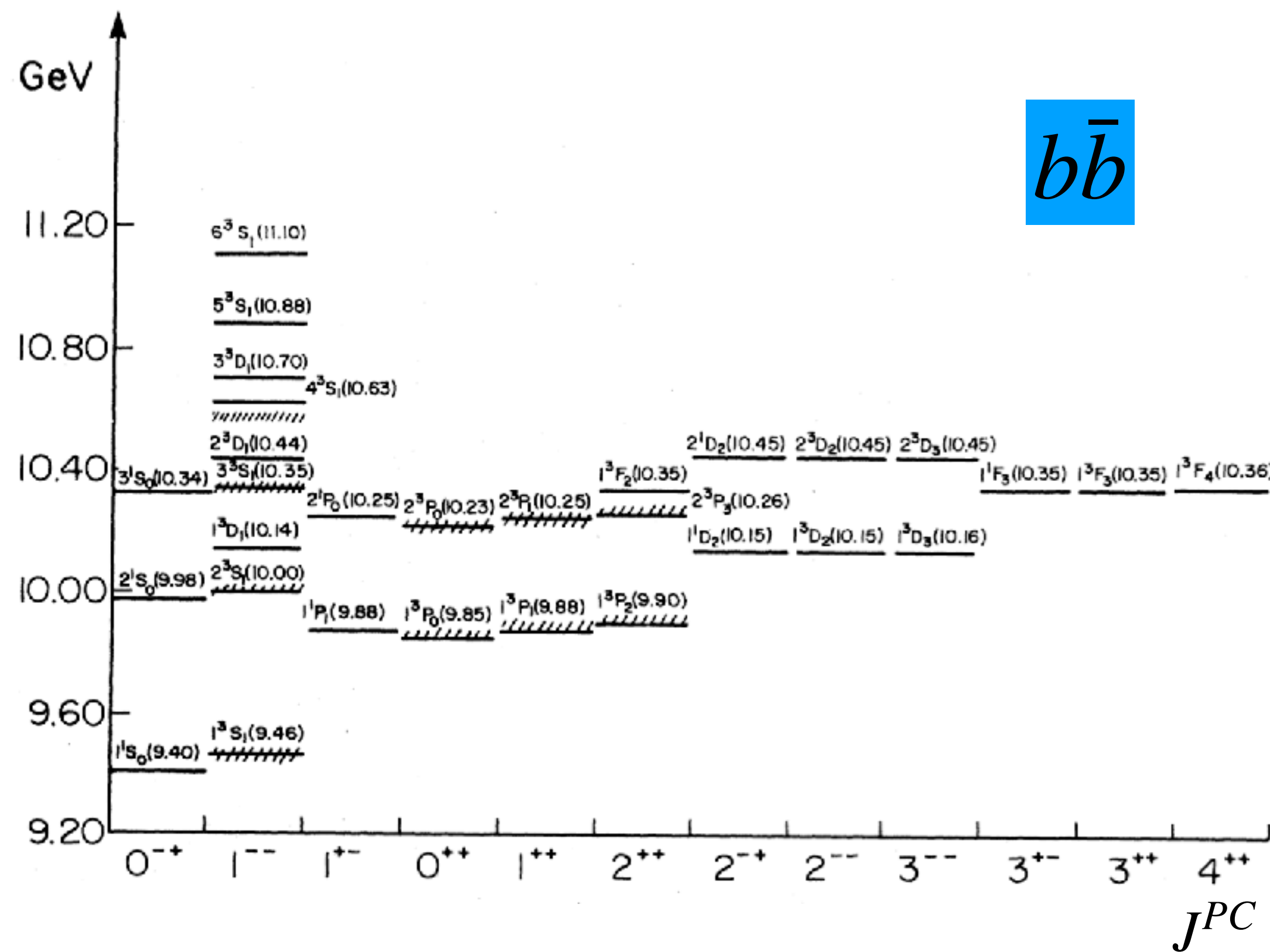
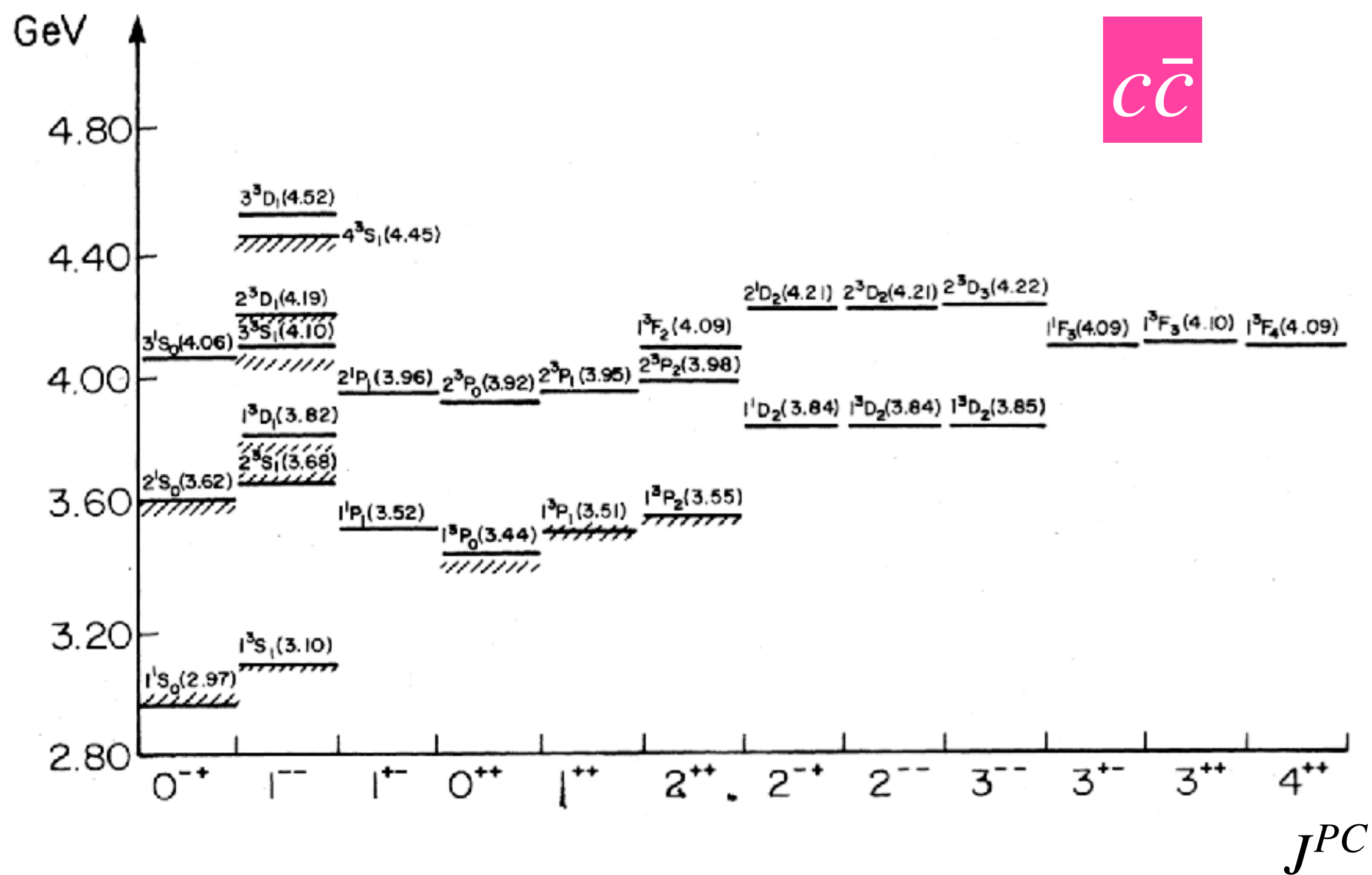


Design:
Max. E_{cm} ~ 3 GeV!



Heavy quark system can be well described by non-relativistic potential model.

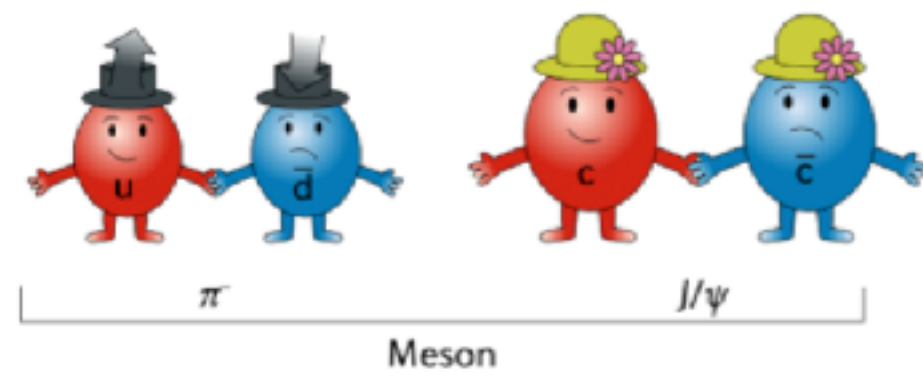
Godfrey & Isgur, PhysRevD.32.189 (1985)



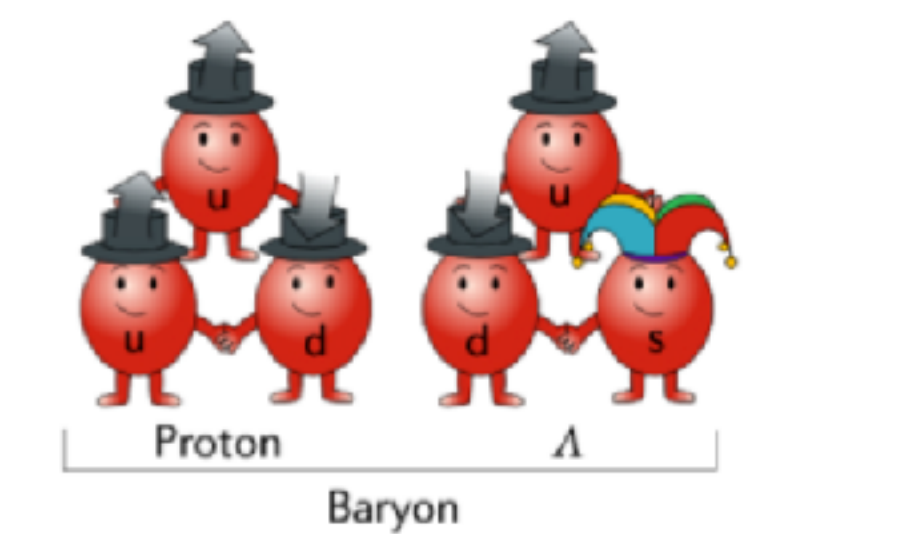
The End?

Hadrons

mesons



baryons



Normal
&
Exotic

A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon b if we assign to the triplet t the following properties: spin $\frac{1}{2}$, $z = -\frac{1}{3}$, and baryon number $\frac{1}{3}$. We then refer to the members $u^{\frac{2}{3}}$, $d^{-\frac{1}{3}}$, and $s^{-\frac{1}{3}}$ of the triplet as "quarks" q and the members of the anti-triplet as anti-quarks \bar{q} . **Baryons** can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while **mesons** are made out of $(q\bar{q})$, $(q\bar{q}q\bar{q})$, etc. It is assuming that the lowest baryon configuration (qqq) gives just the representations 1, 8, and 10 that have been observed, while the lowest meson configuration $(q\bar{q})$ similarly gives just 1 and 8.



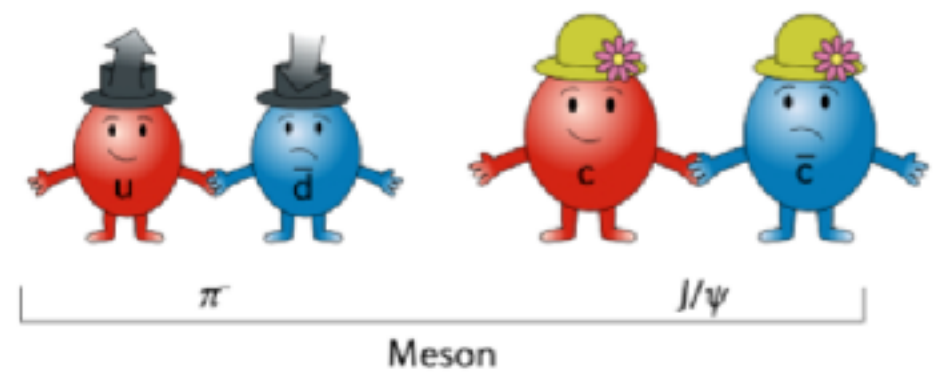
M. Gell-Mann

Physics Letters 8, 214 (1964);

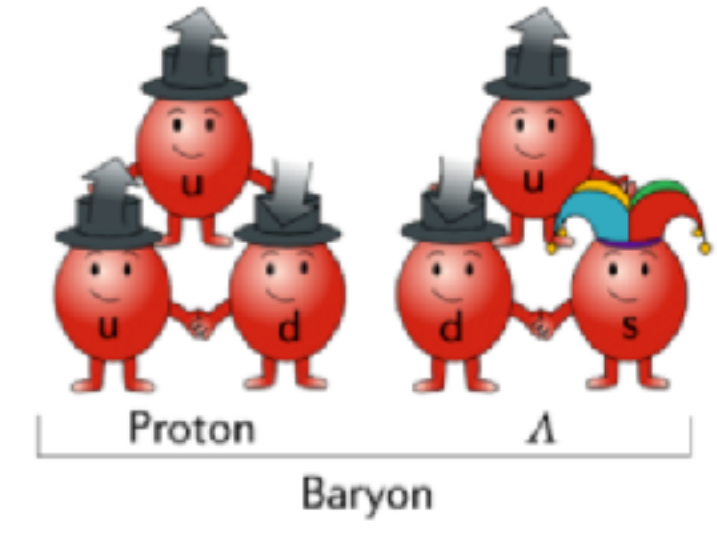
Quark model doesnot forbid hadrons with $N_{\text{quark}} \neq 2,3!$

Hadrons

mesons



baryons

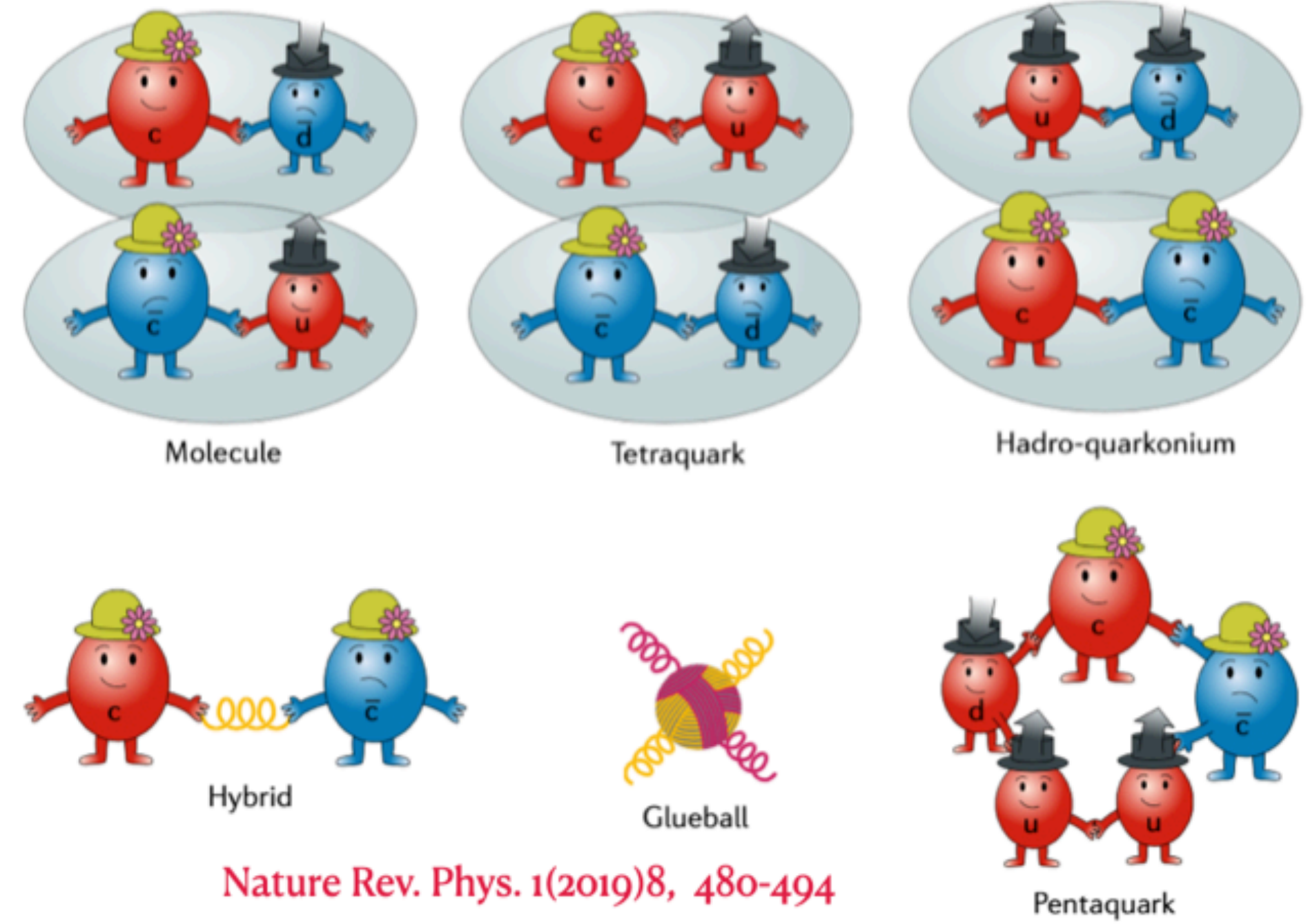


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M. Gell-Mann
Physics Letters 8, 214 (1964);

Normal & Exotic

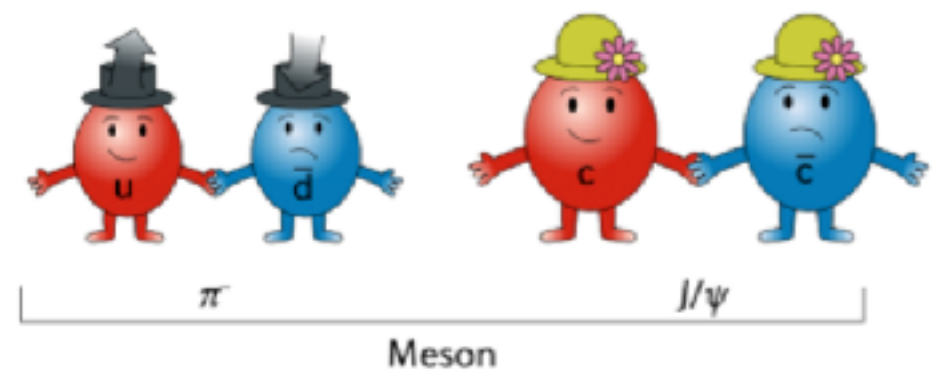


Quark model doesnot forbid hadrons with $N_{\text{quark}} \neq 2,3!$

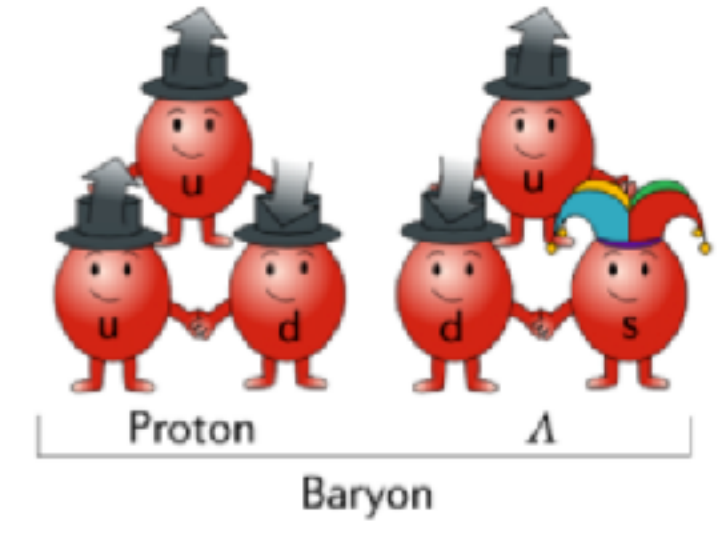
Nature Rev. Phys. 1(2019)8, 480-494

Hadrons

mesons



baryons

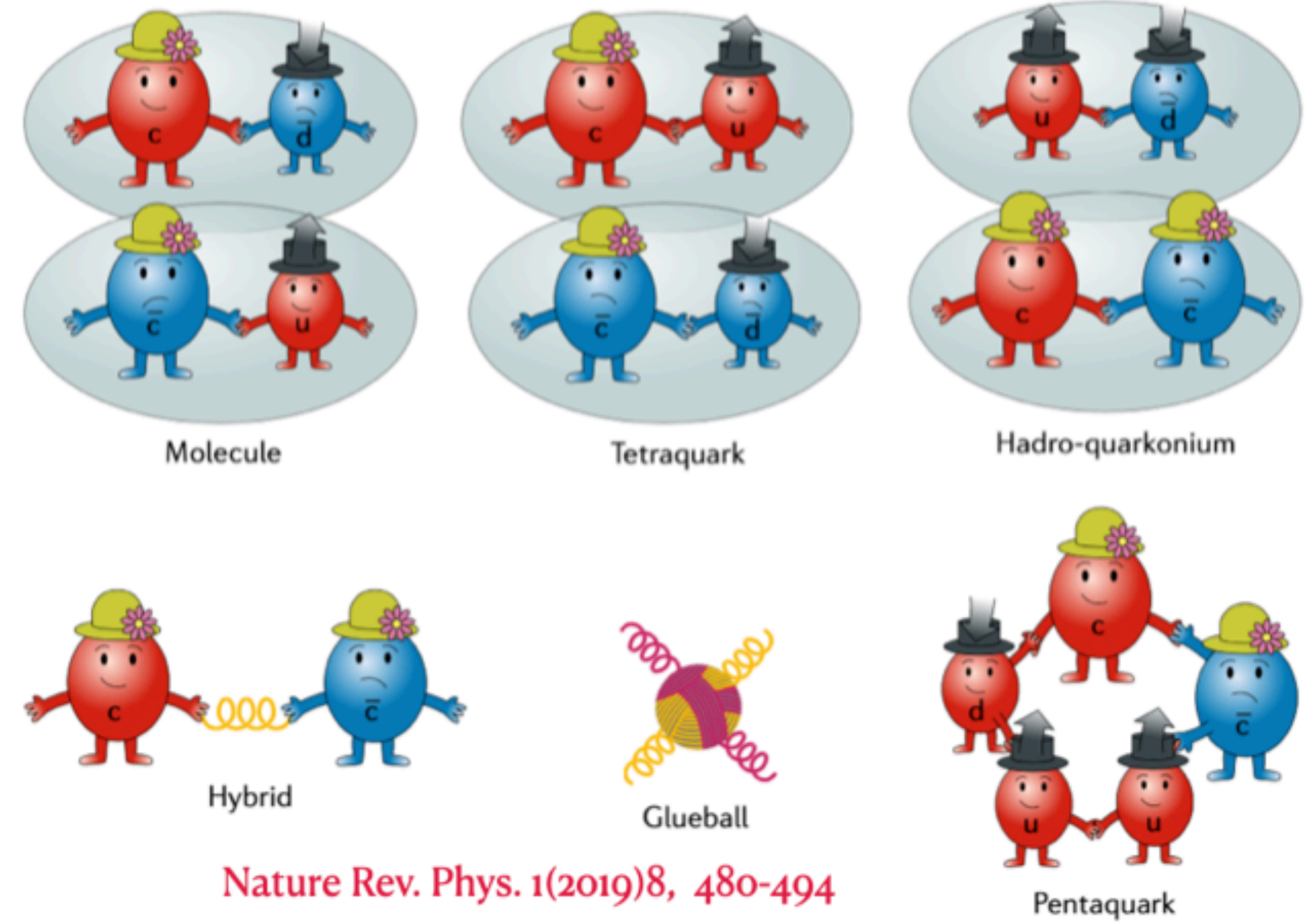


A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon \mathbf{b} if we assign to the triplet \mathbf{t} the following properties: spin $\frac{1}{2}$, $z = -\frac{1}{3}$, and baryon number $\frac{1}{3}$. We then refer to the members $u^{\frac{2}{3}}$, $d^{-\frac{1}{3}}$, and $s^{-\frac{1}{3}}$ of the triplet as "quarks" q and the members of the anti-triplet as anti-quarks \bar{q} . **Baryons** can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while **mesons** are made out of $(q\bar{q})$, $(qq\bar{q}\bar{q})$, etc. It is assuming that the lowest baryon configuration (qqq) gives just the representations **1**, **8**, and **10** that have been observed, while the lowest meson configuration $(q\bar{q})$ similarly gives just **1** and **8**.



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Physics Letters 8, 214 (1964);

Normal & Exotic



Nature Rev. Phys. 1(2019)8, 480-494

Quark model doesnot forbid hadrons with $N_{\text{quark}} \neq 2,3!$

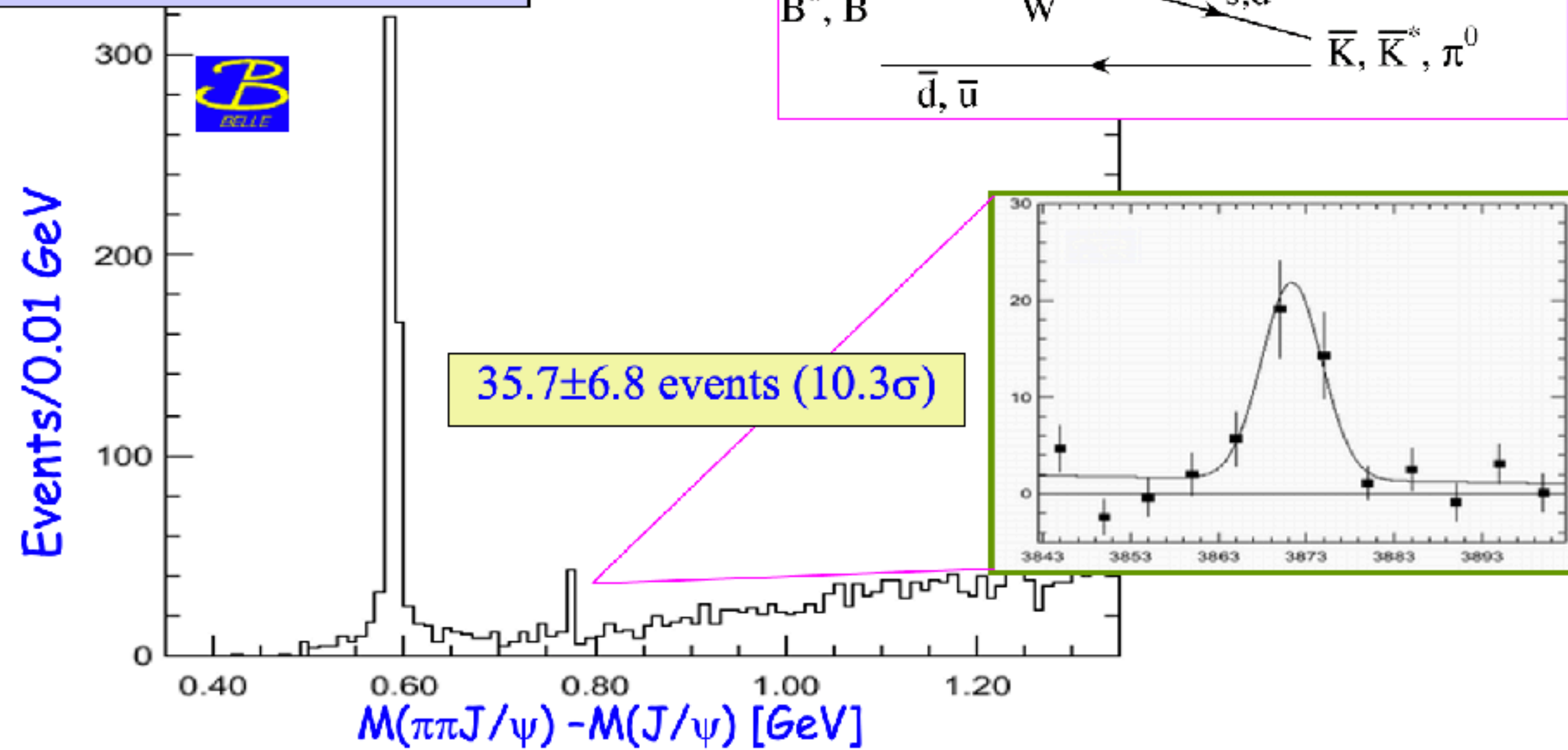
Where?

2003, Belle observed the first exotic state $X(3872)$

confirmed by BaBar, CDF, D0.

X(3872): first observation at Belle

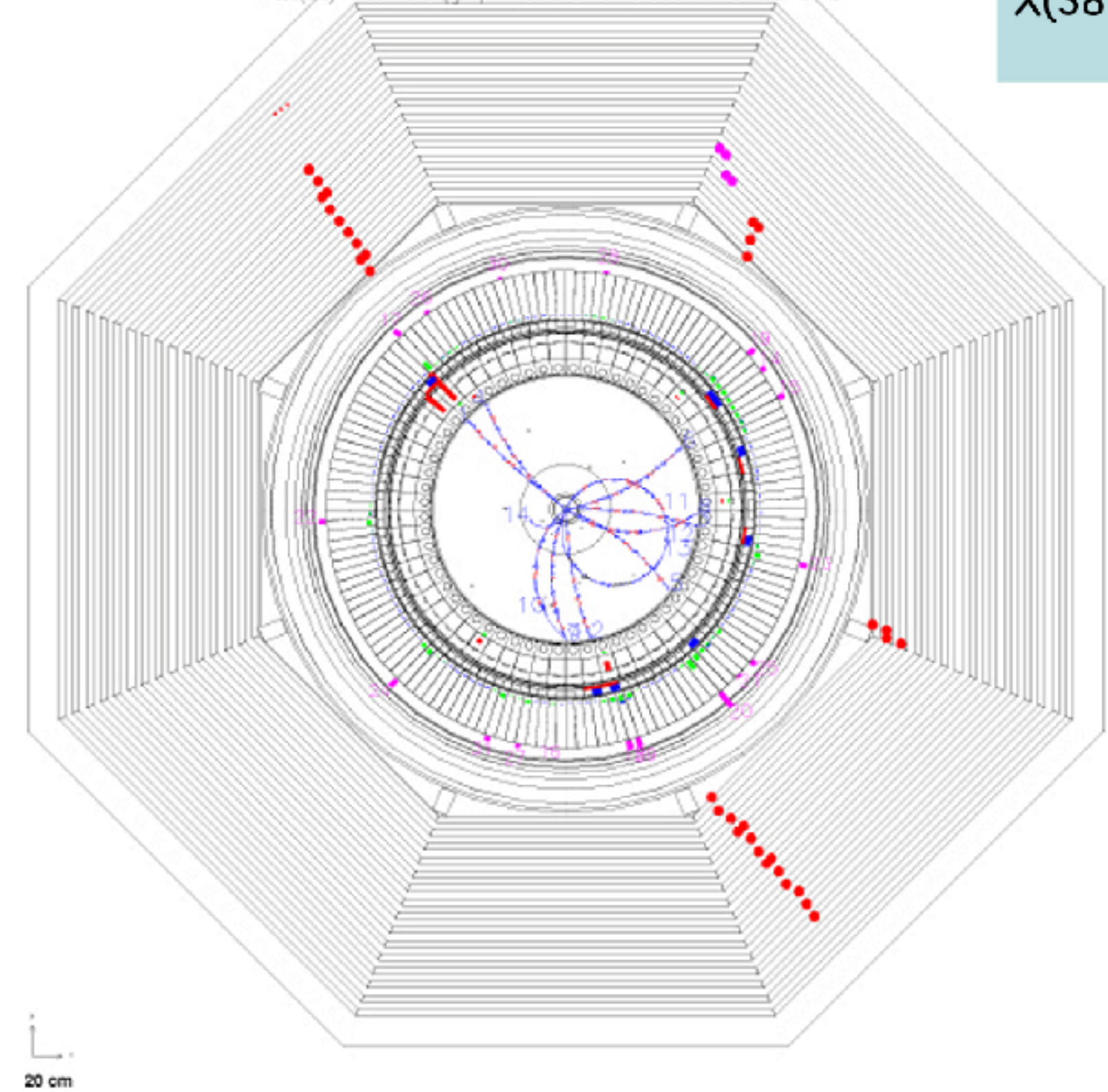
PRL91 262001 (2003)
140/fb data



BELLE

Exp 25 Run 1016 Form 0 Event 260737
Eher 0.00 Eler 0.00 Mon Feb 24 16z59z14 2003
TrgID 0 DetVer 0 MagID 0 BField 1.50 DspVer 7.50
Ptot(ch) 9.4 Etot(gm) 1.0 SVD-M 0 CDC-M 0 KLM-M 0

X(3872)事象の
1例



Why $X(3872)$ is exotic?

Mass:

- $3871.65 \pm 0.06 \text{ MeV}/c^2$

Not predicted in the potential model

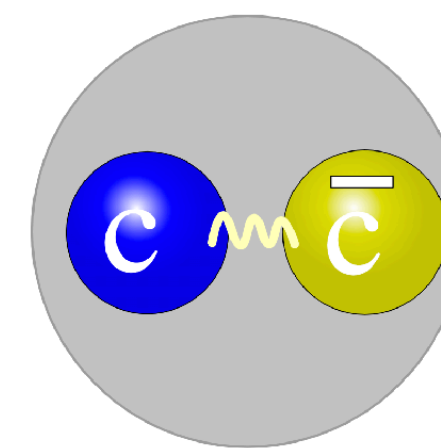
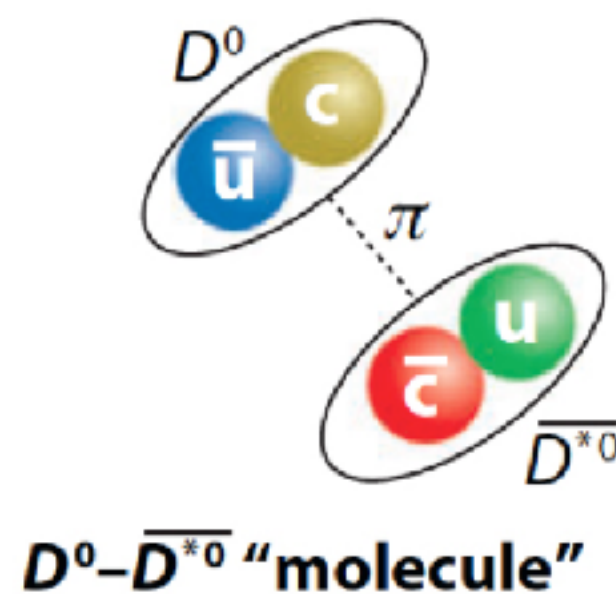
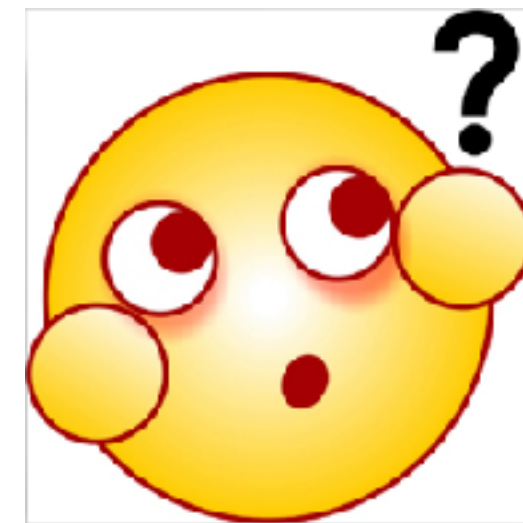
- Mass very close to $m(D^0 \bar{D}^{*0})$

Binding energy: $-10 \pm 90 \text{ keV}/c^2$

$$E_b = \hbar^2 / 2\mu a^2$$

$$\langle r \rangle = a / \sqrt{2} \geq 31.7_{-24.5}^{+\infty} \text{ fm}$$

$$E_b(\text{deuteron}) = -2.2 \text{ MeV}$$



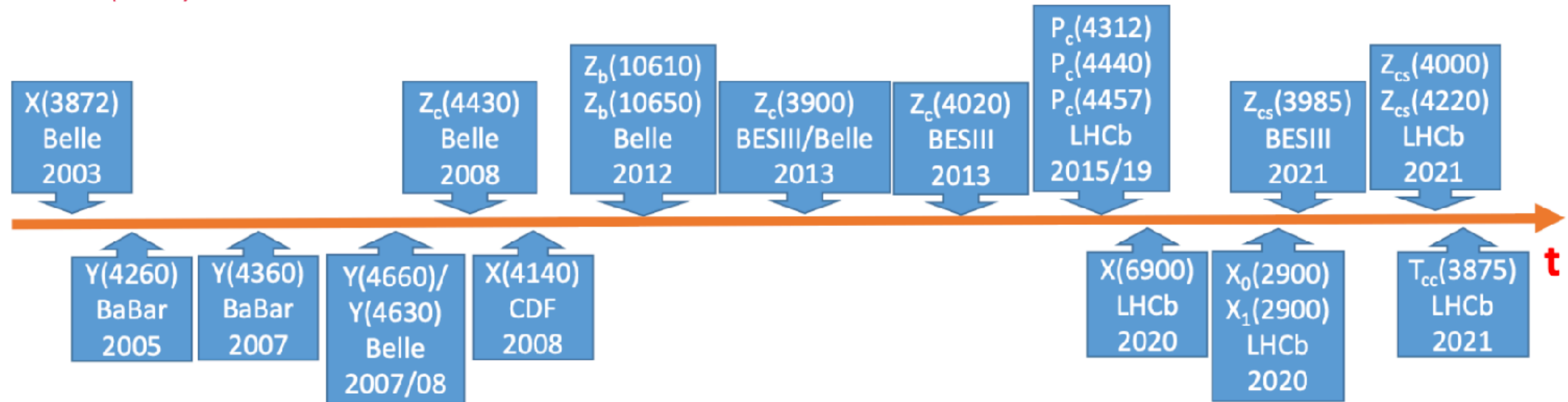
Width:

- $1.19 \pm 0.21 \text{ MeV}$
(The shorter the life, the larger the width)
- Lifetime is 20 times larger than the other charmonia nearby.

Production mechanism:

- B decay
- prompt production in pp collision
(should be small in case of charmonium)
- e^+e^- radiative decay

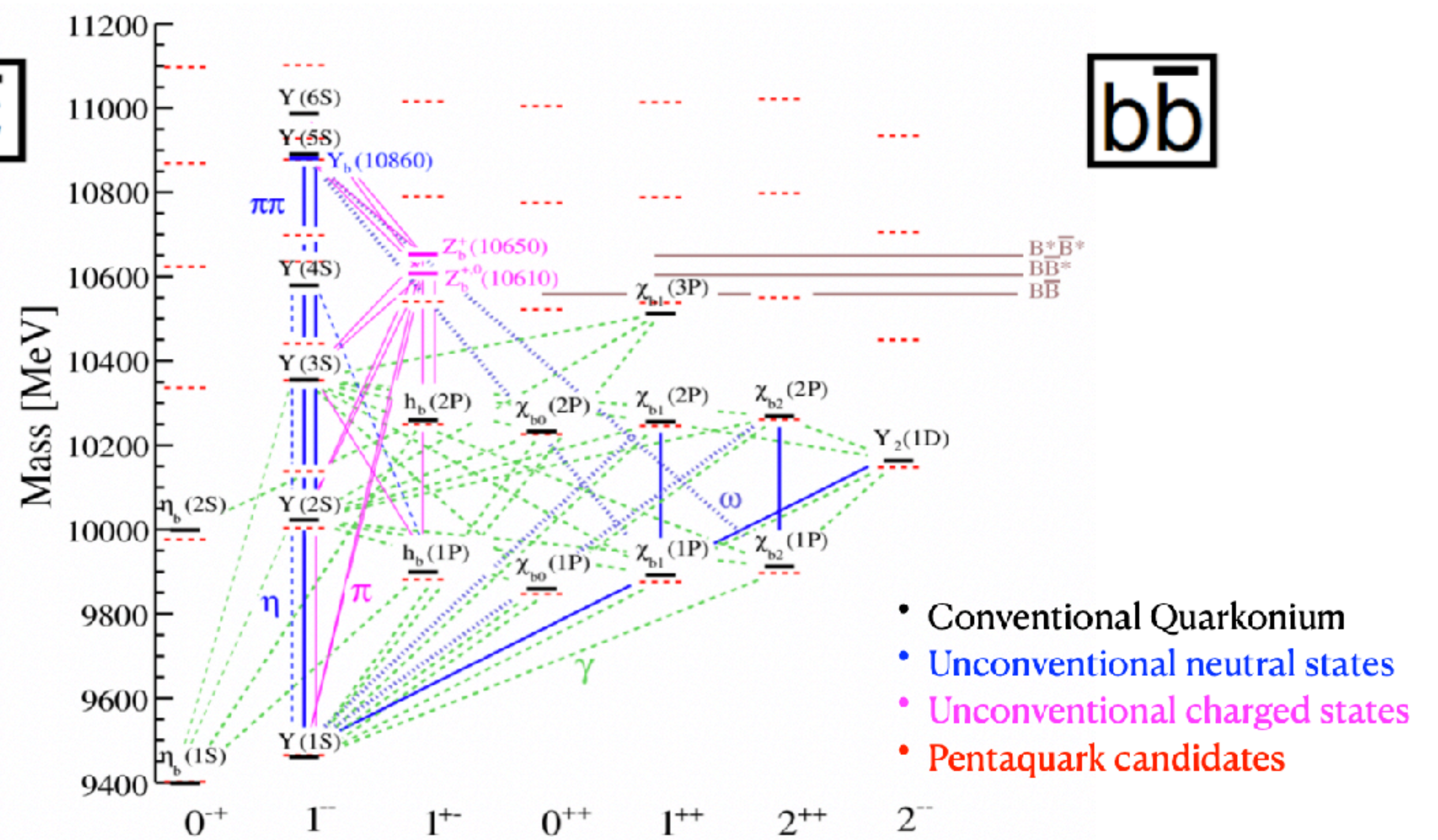
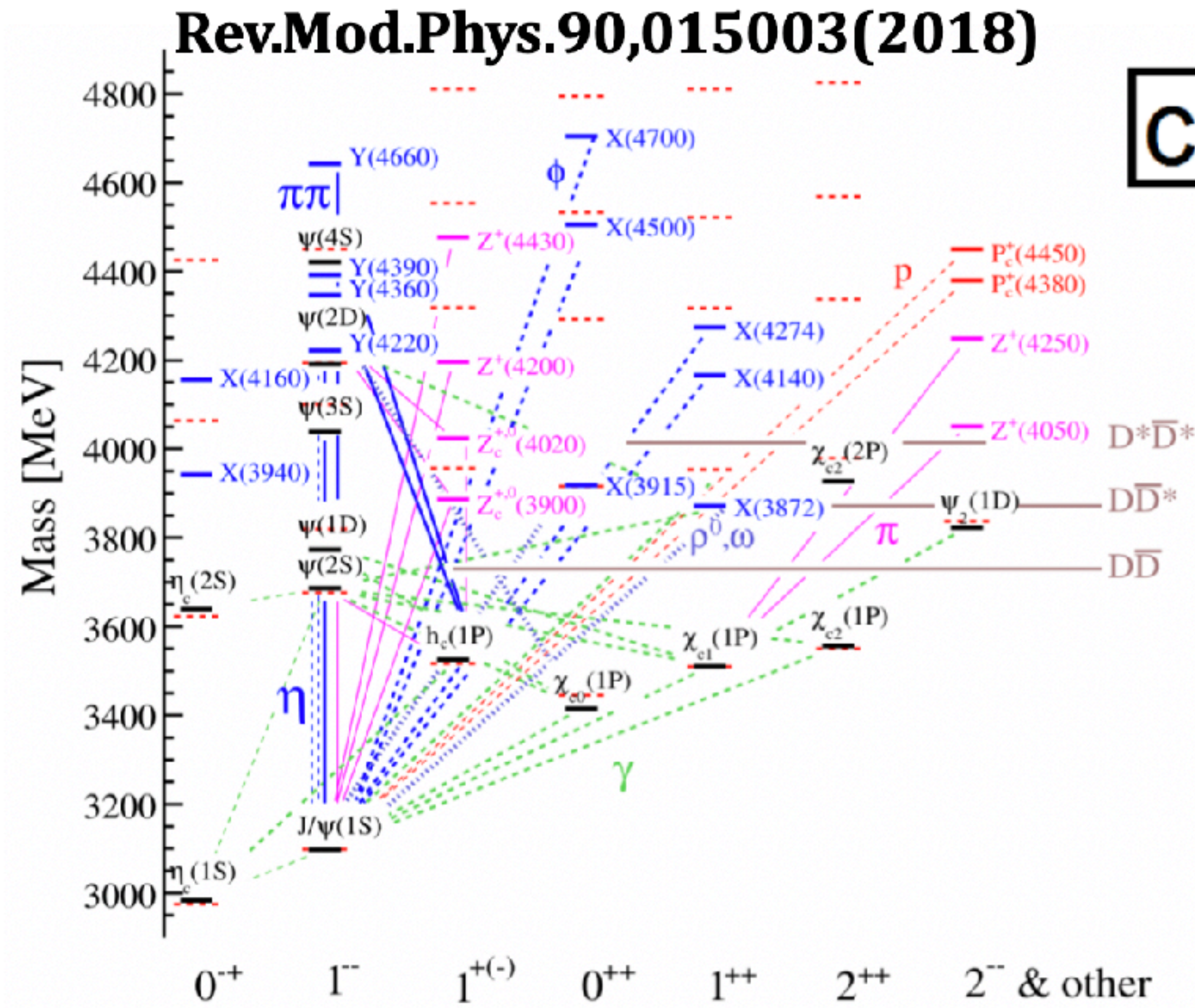
Phys.Rev.Lett. 91 (2003) 262001



- Ever since the discovery of $X(3872)$, we have a golden era in the discovery of the exotic states.

Quarkonium spectroscopy

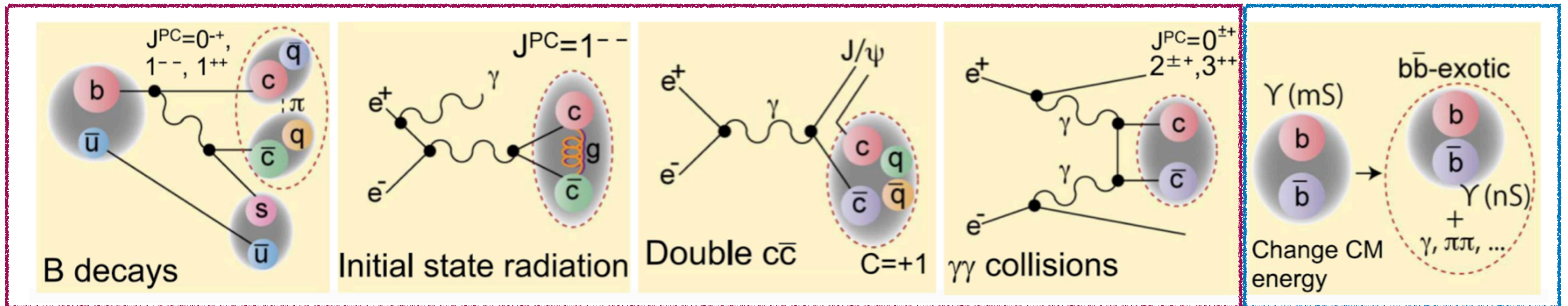
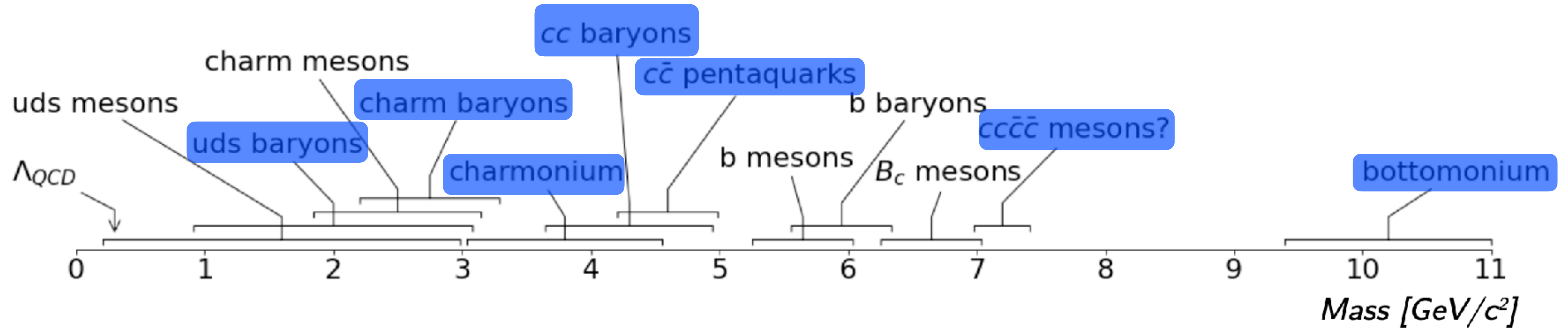
Parallel properties in $c\bar{c}$ and $b\bar{b}$.
Excellent experimental field!



Below $D\bar{D}/B\bar{B}$ threshold: Good agreement!

Above $D\bar{D}/B\bar{B}$ threshold: Unpredicted exotic states!!

Belle&Belle II capabilities on spectroscopy

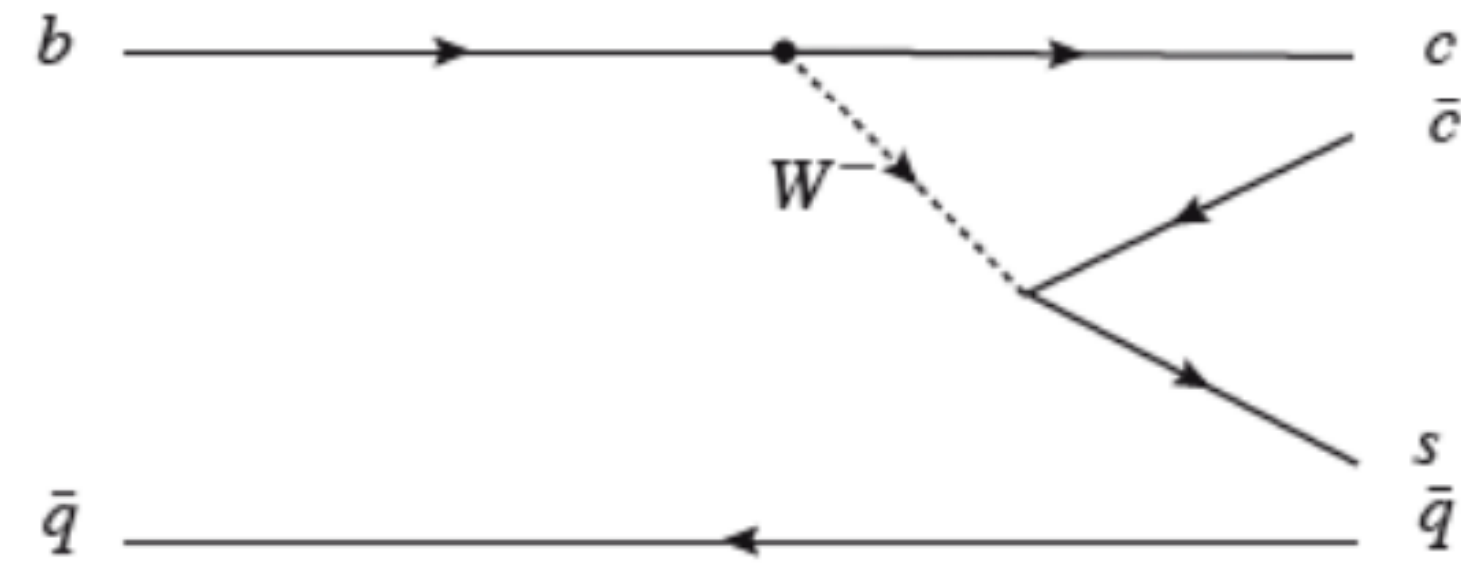


charmonium

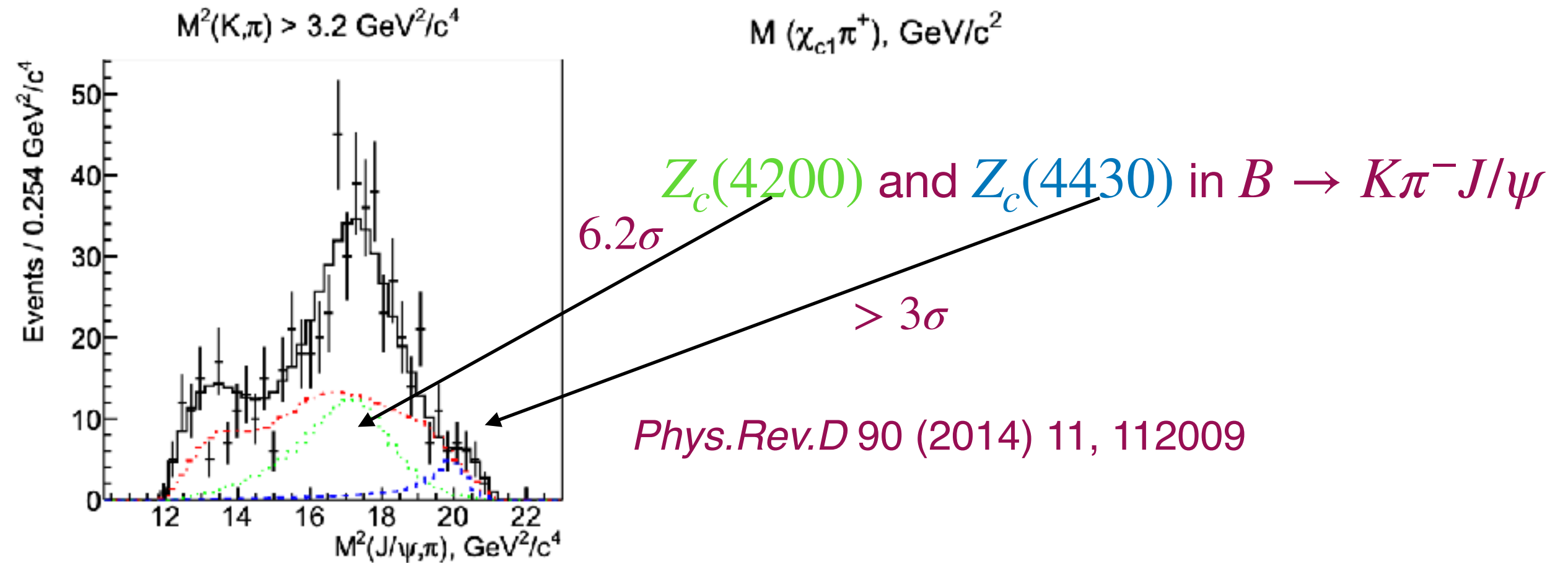
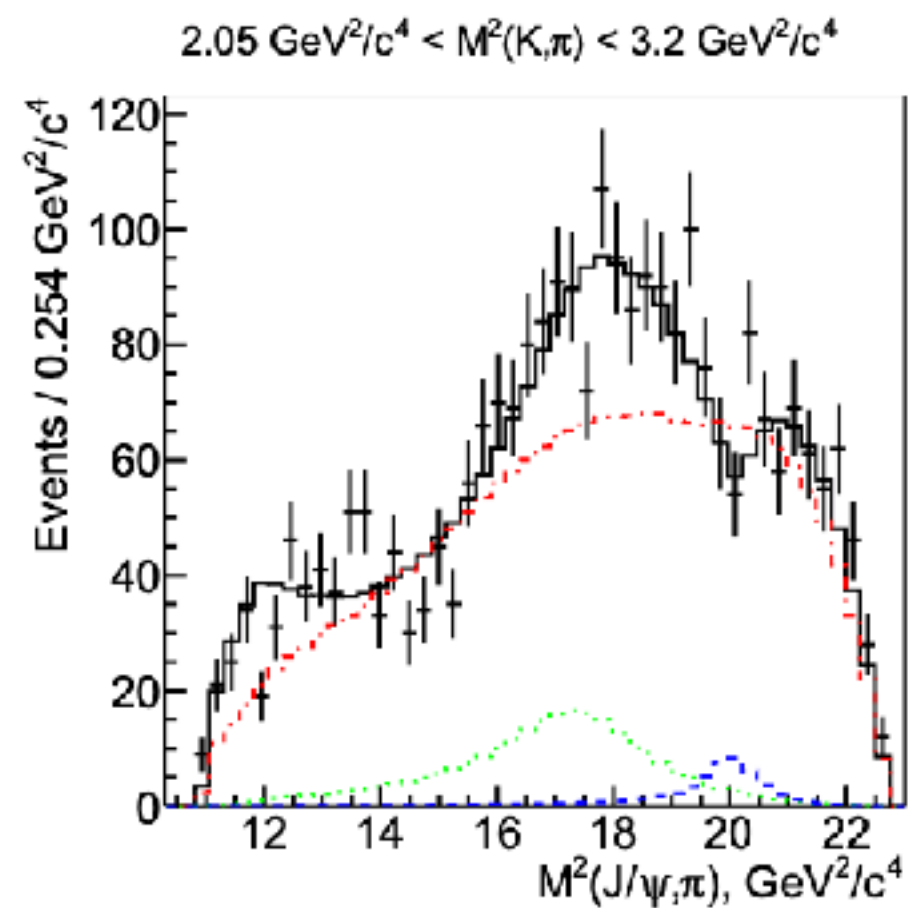
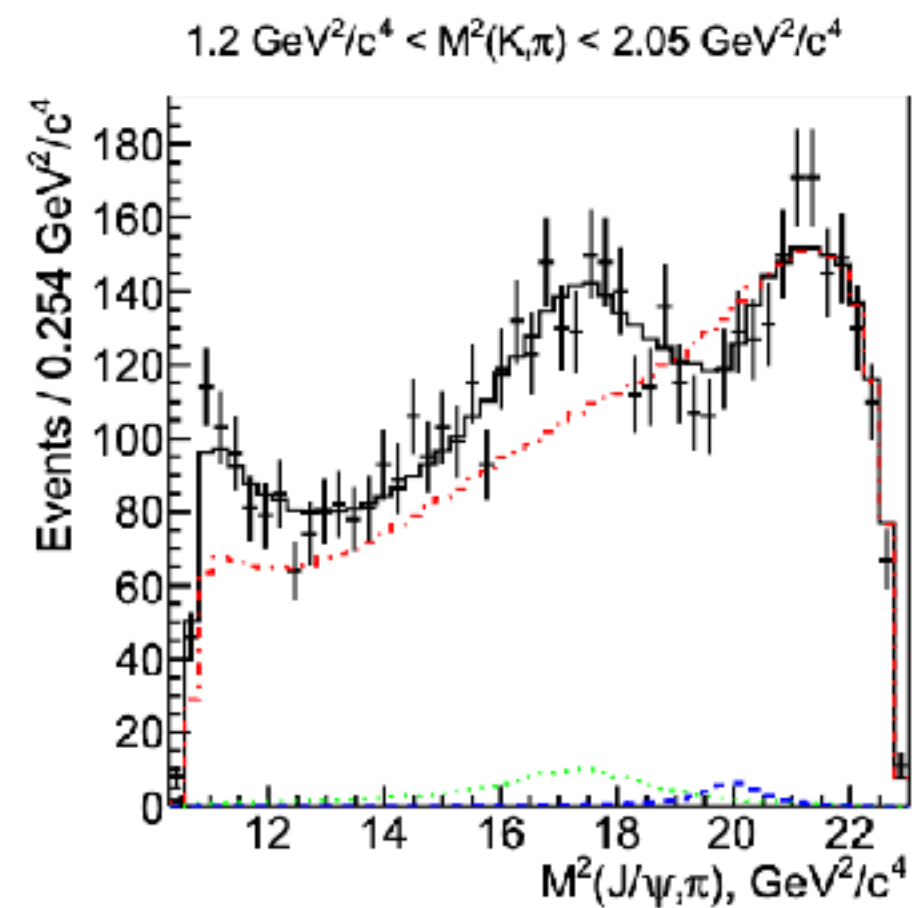
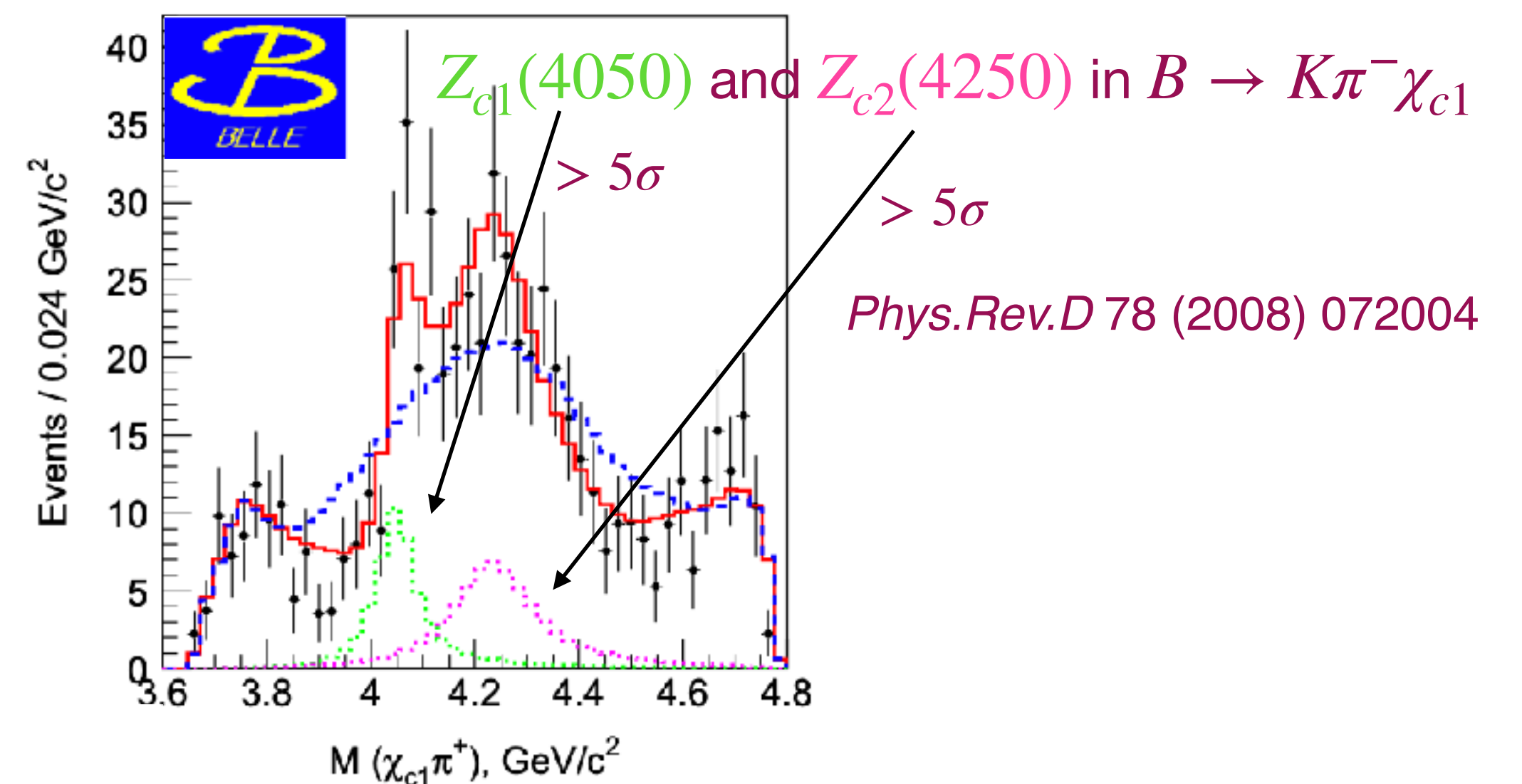
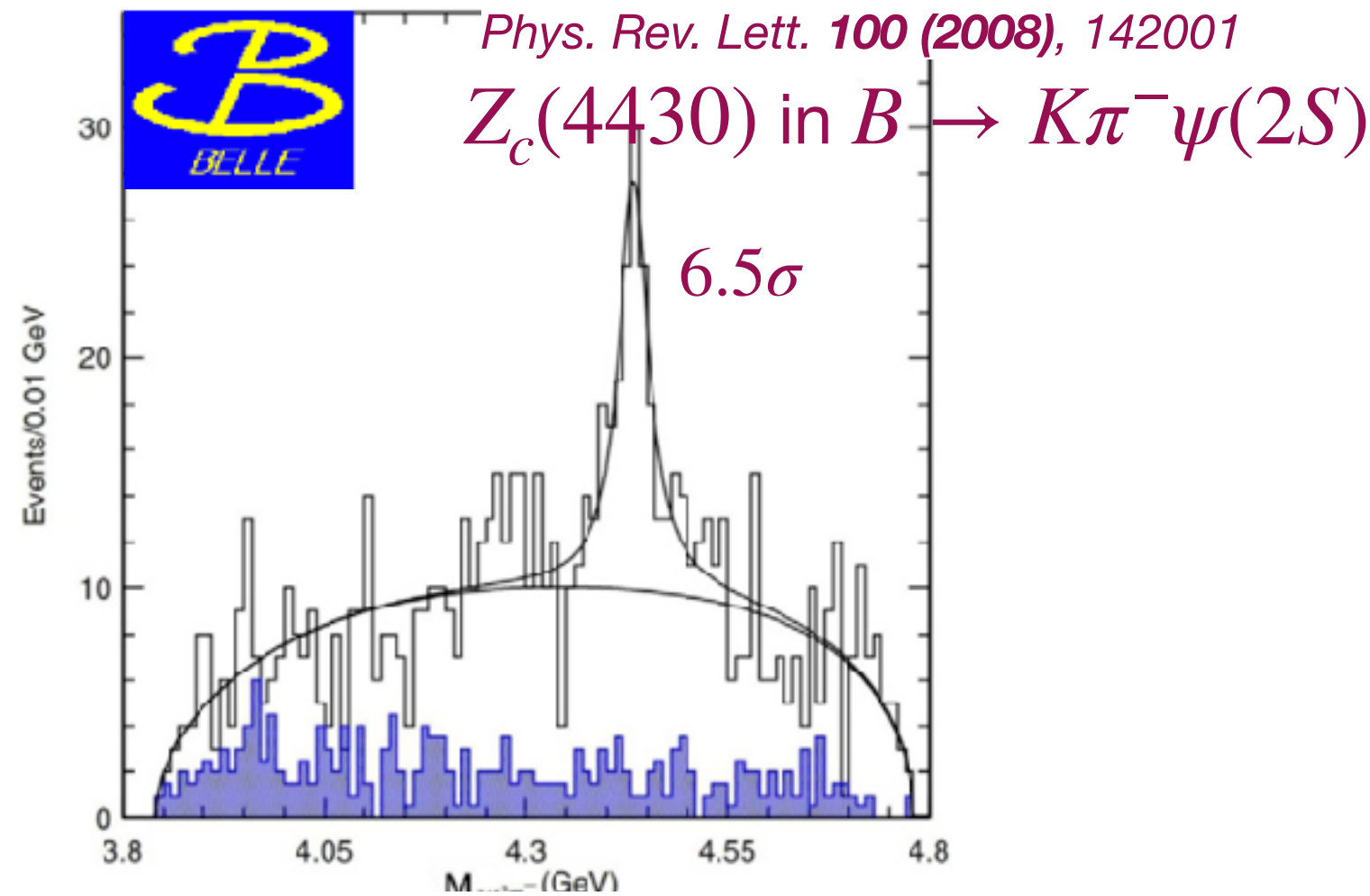
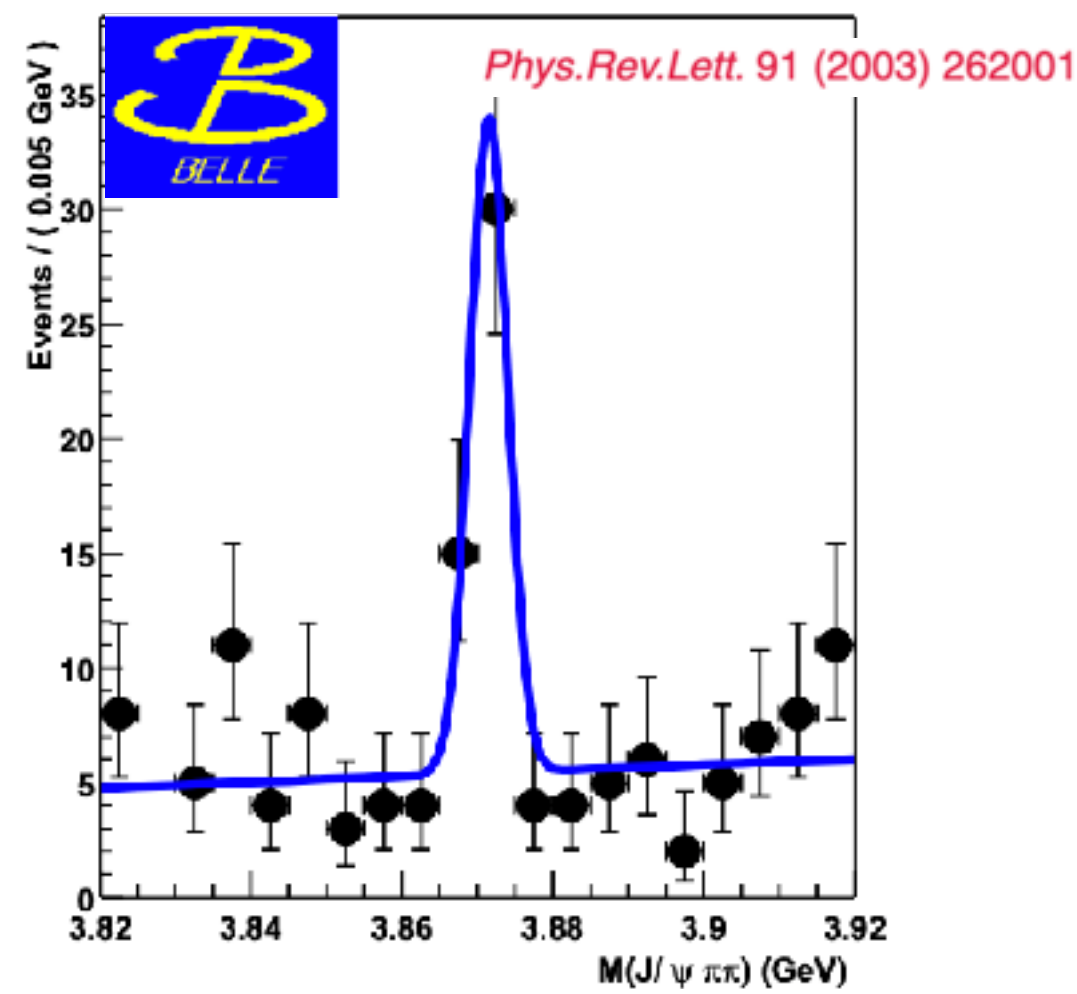
bottomonium

B decays

Large production rate provide a solid ground to search for exotics

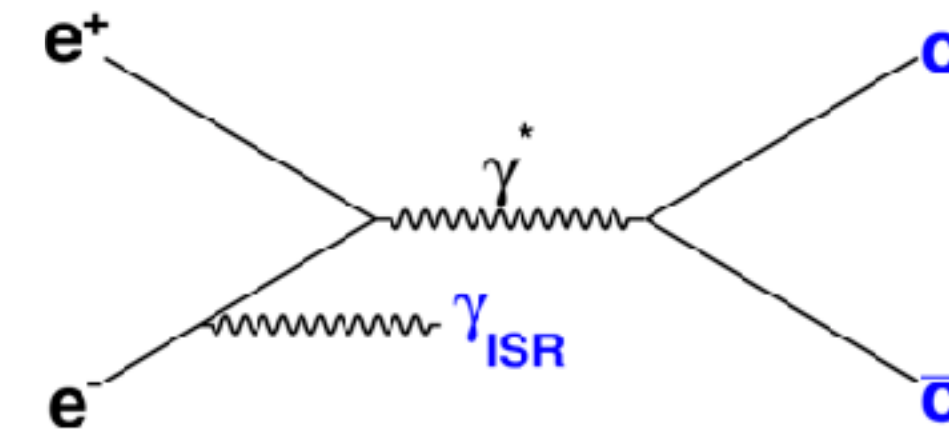


In history: observation of $X(3872)$ & establishment of various Z_c^+ states



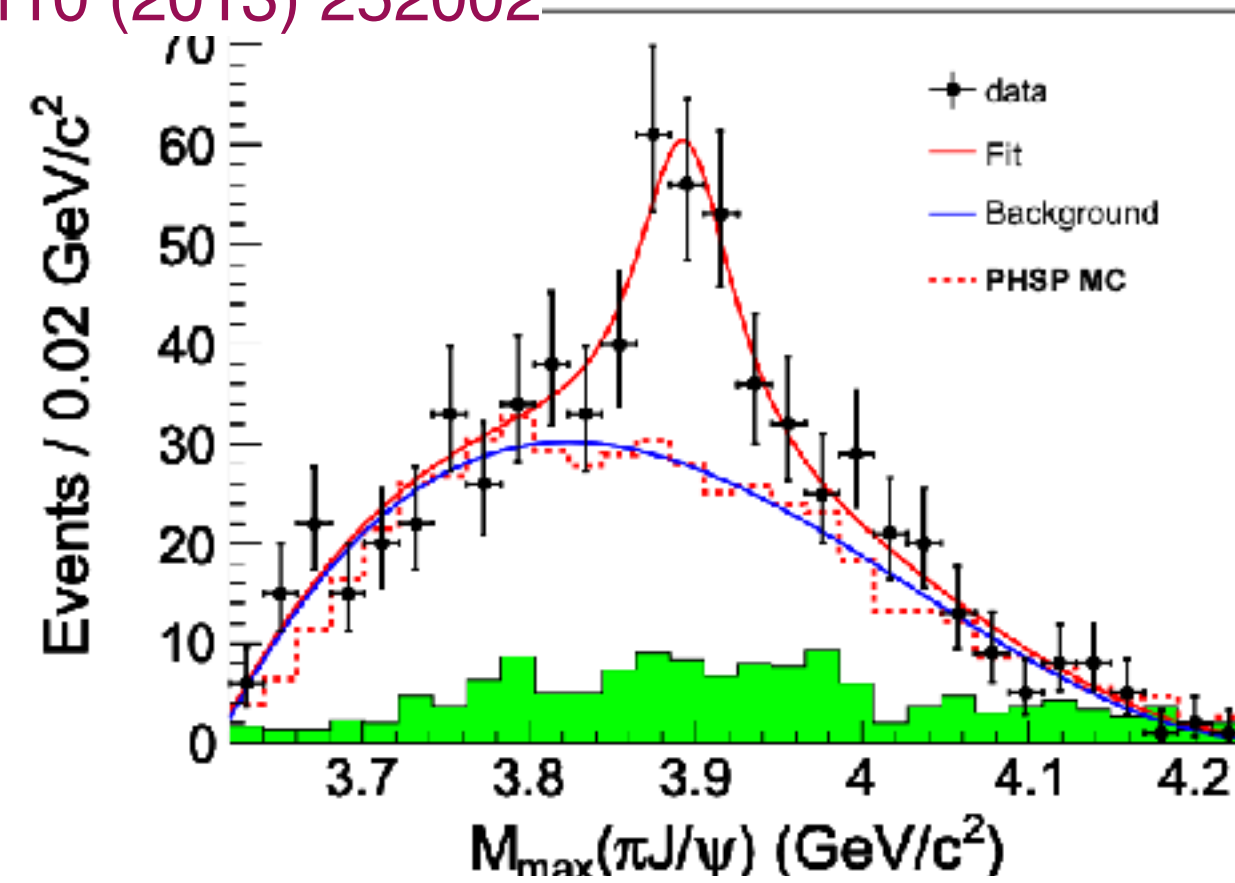
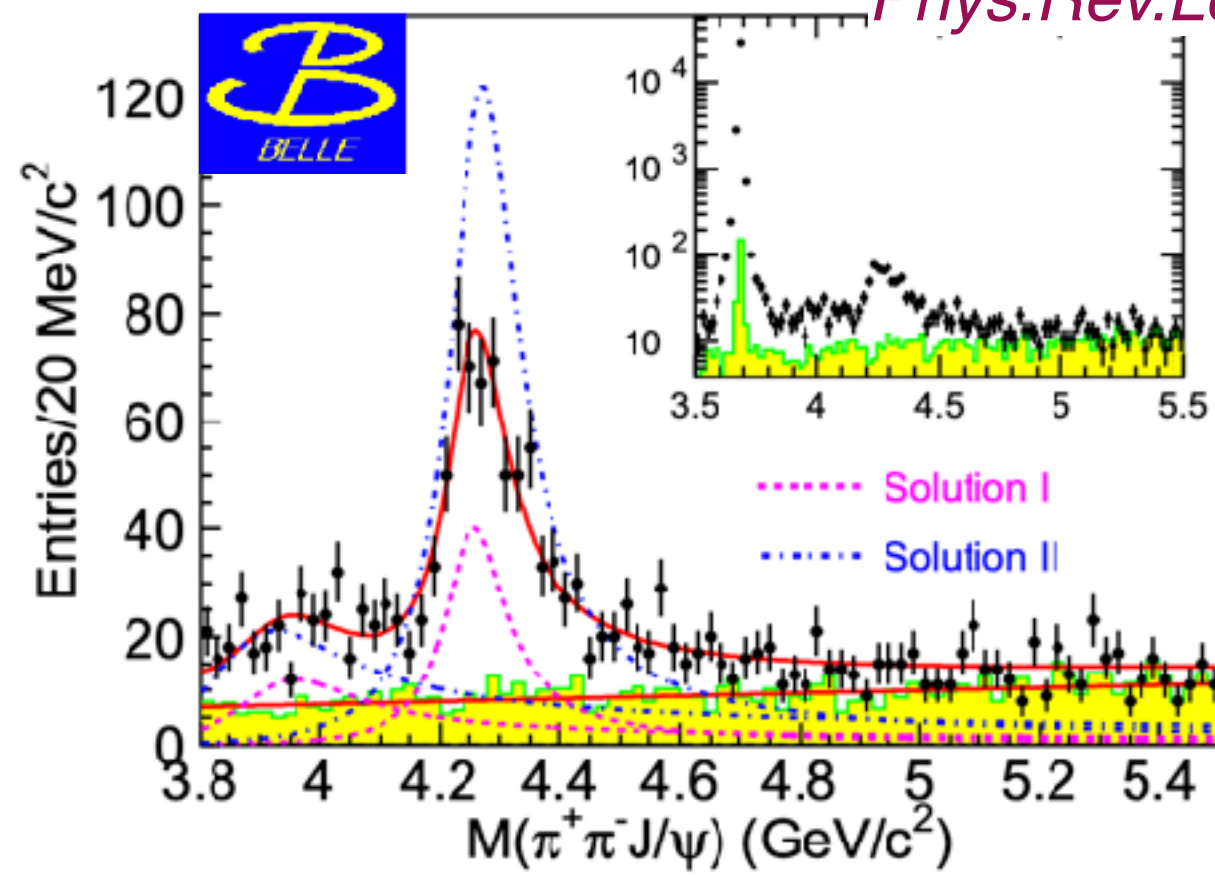
Initial state radiation

Allow us to reach lower c.m. energy "for free"

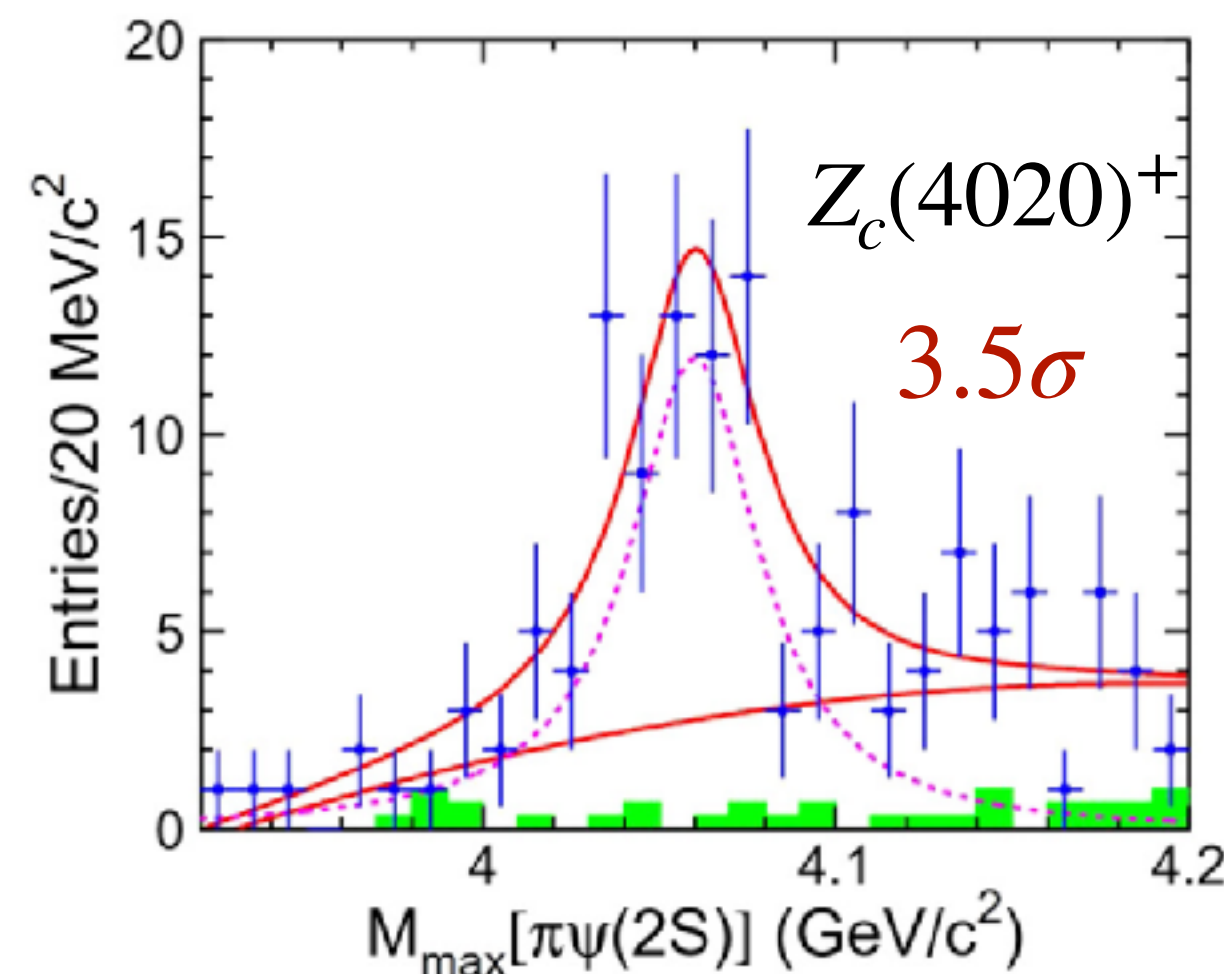
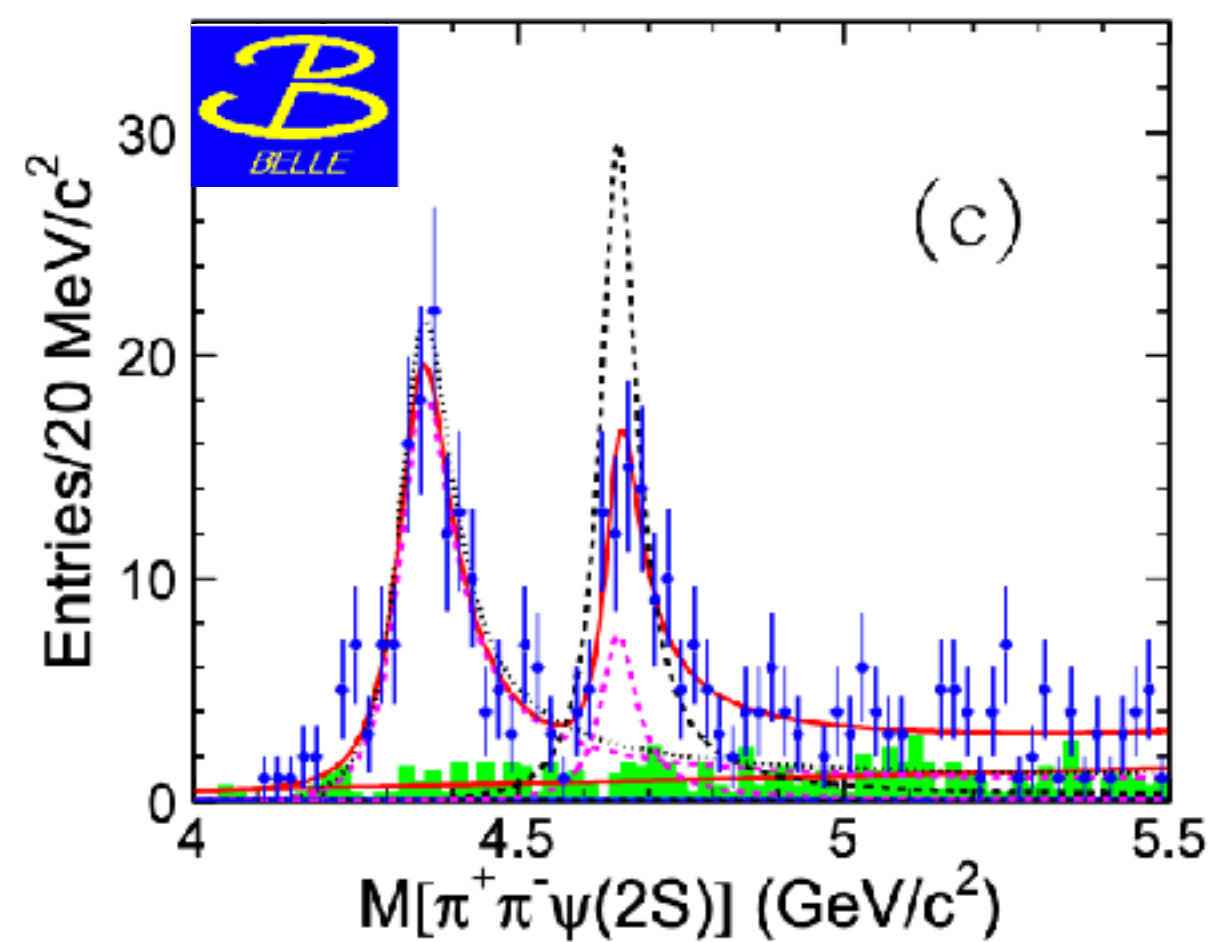
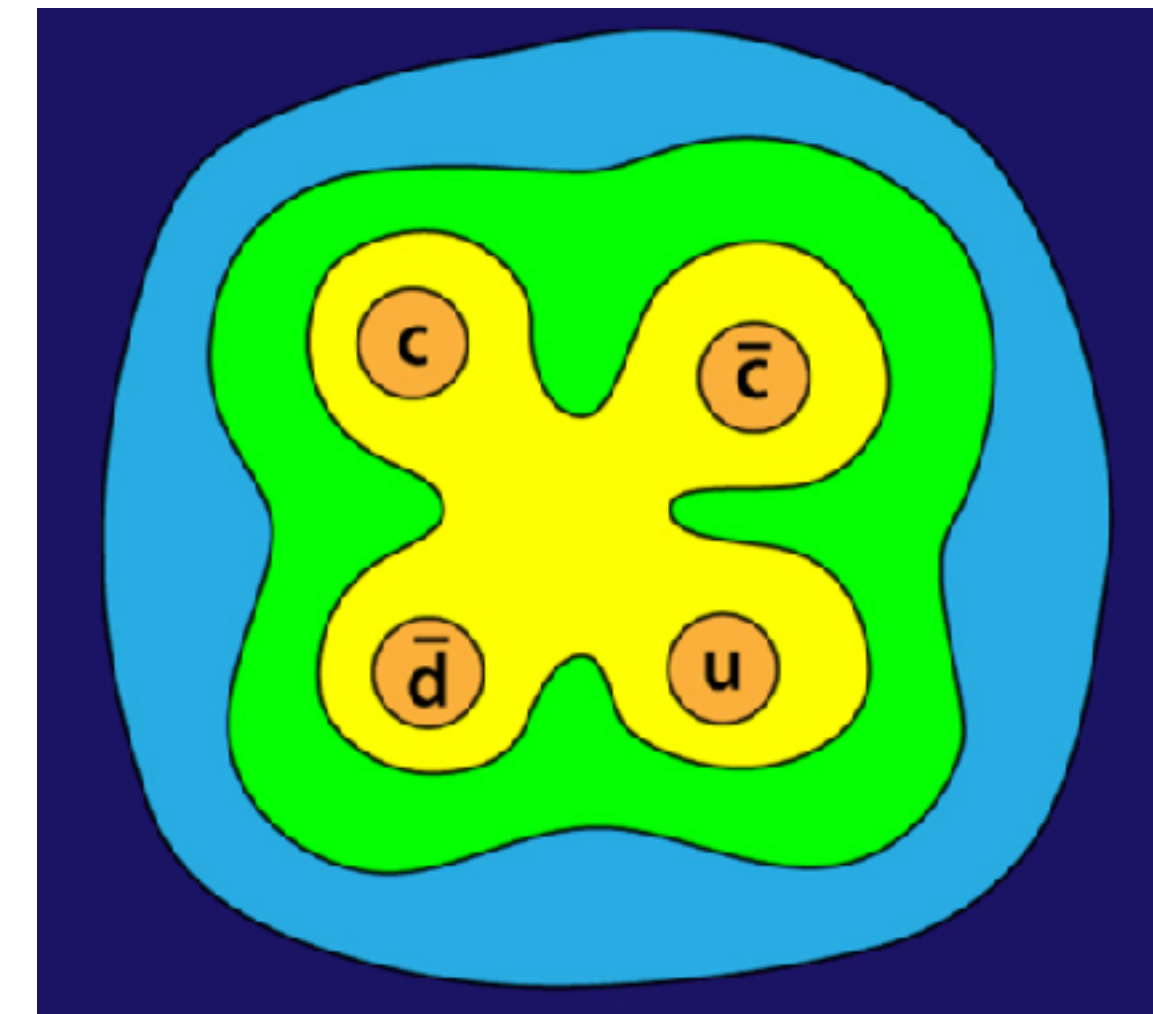


Great achievement in history: Observation of $Z_c(3900)^+$

Phys.Rev.Lett. 110 (2013) 252002



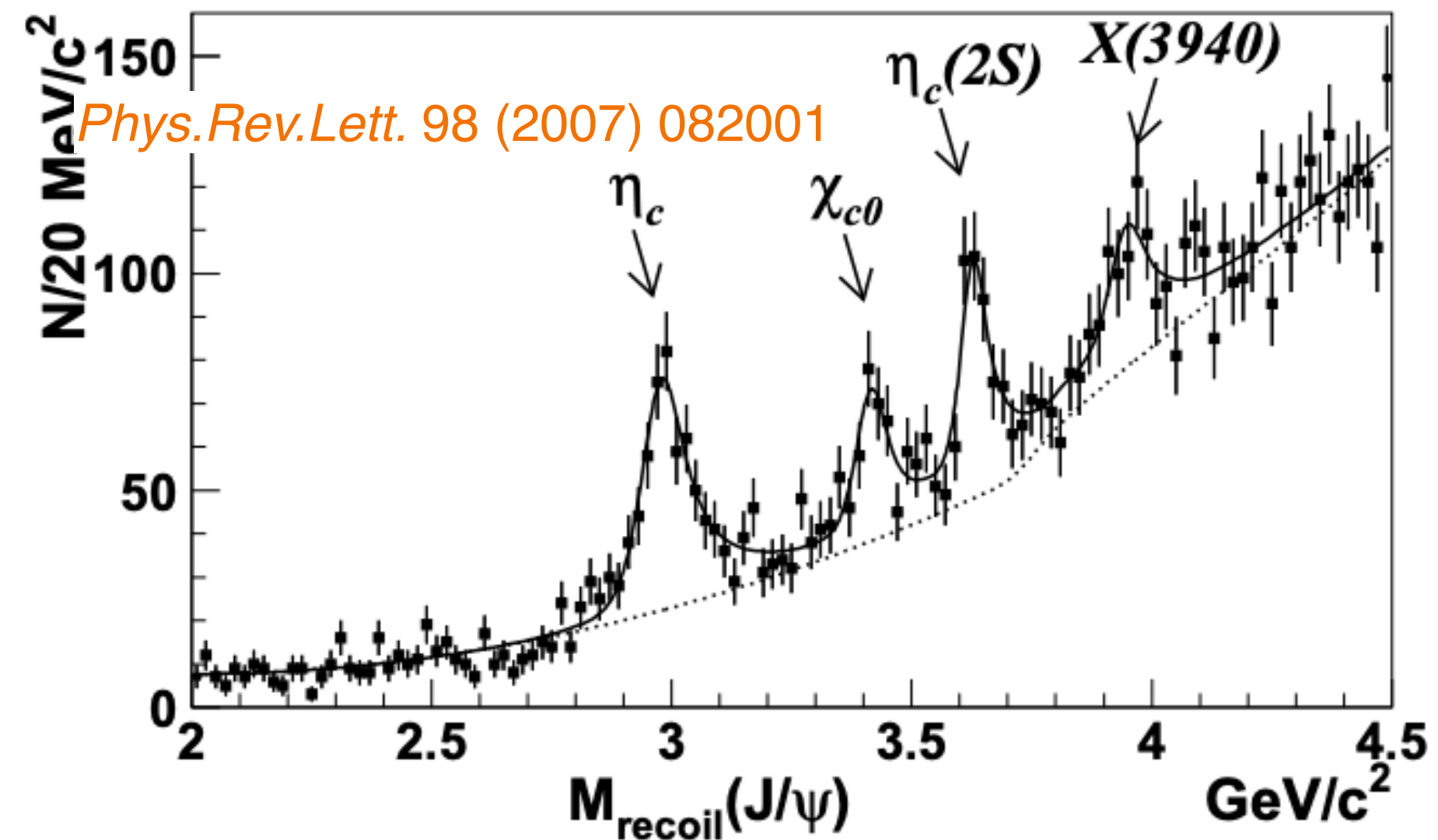
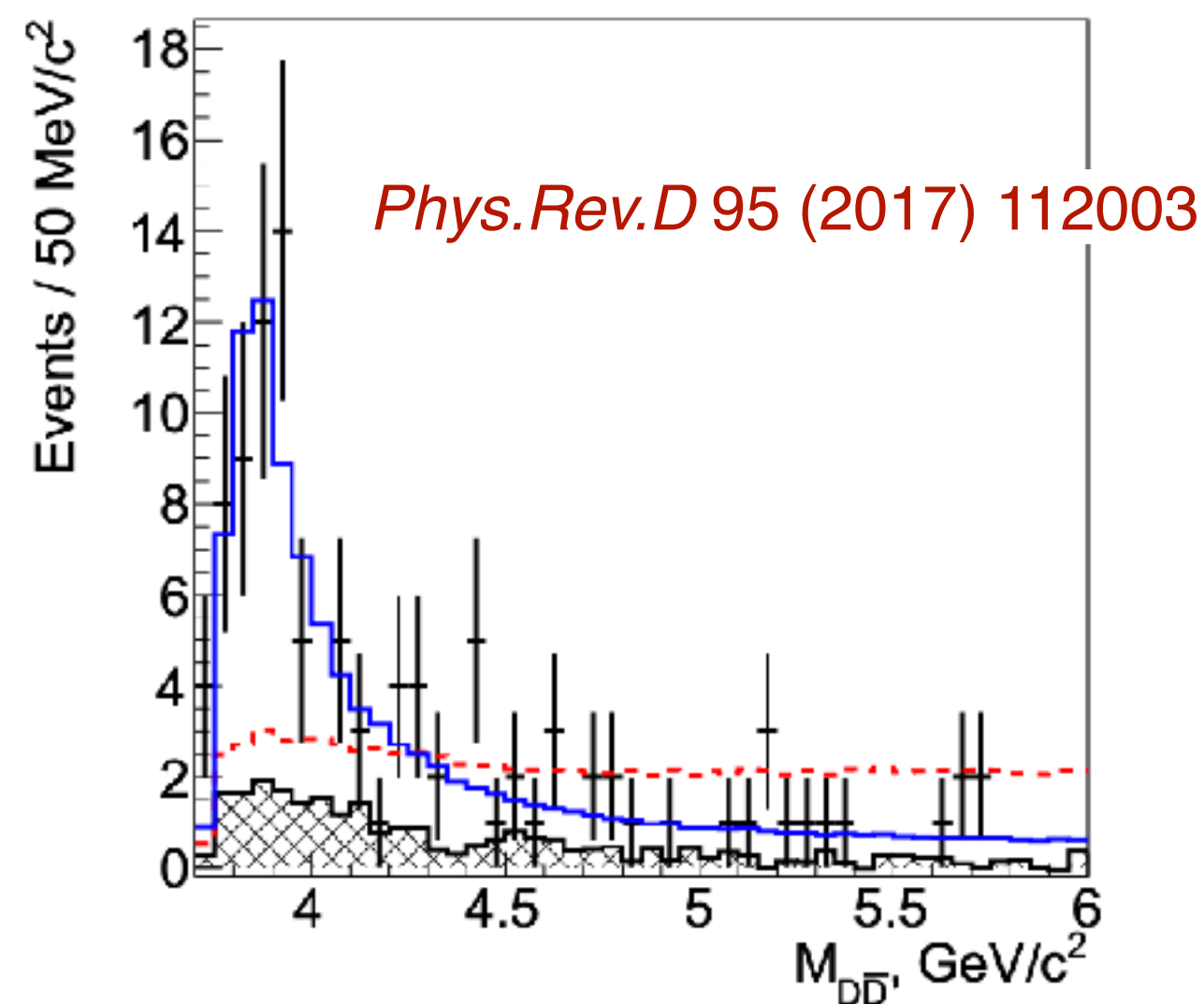
First solid four quark state!



Double charmonium production

Unique field to produce charmonium(-like) particles

Rich resonances produced against J/ψ

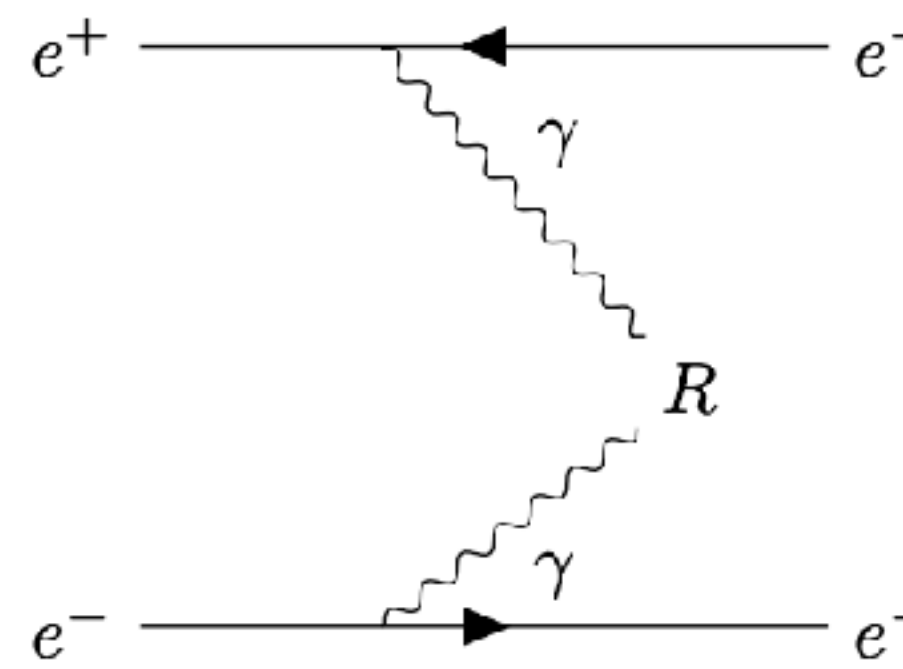


$\chi_{c0}(3860)$ was observed in $e^+e^- \rightarrow J/\psi D\bar{D}$.

$X(3940)$ was discovered in J/ψ recoiling mass while dominantly decays into $D\bar{D}^*$.

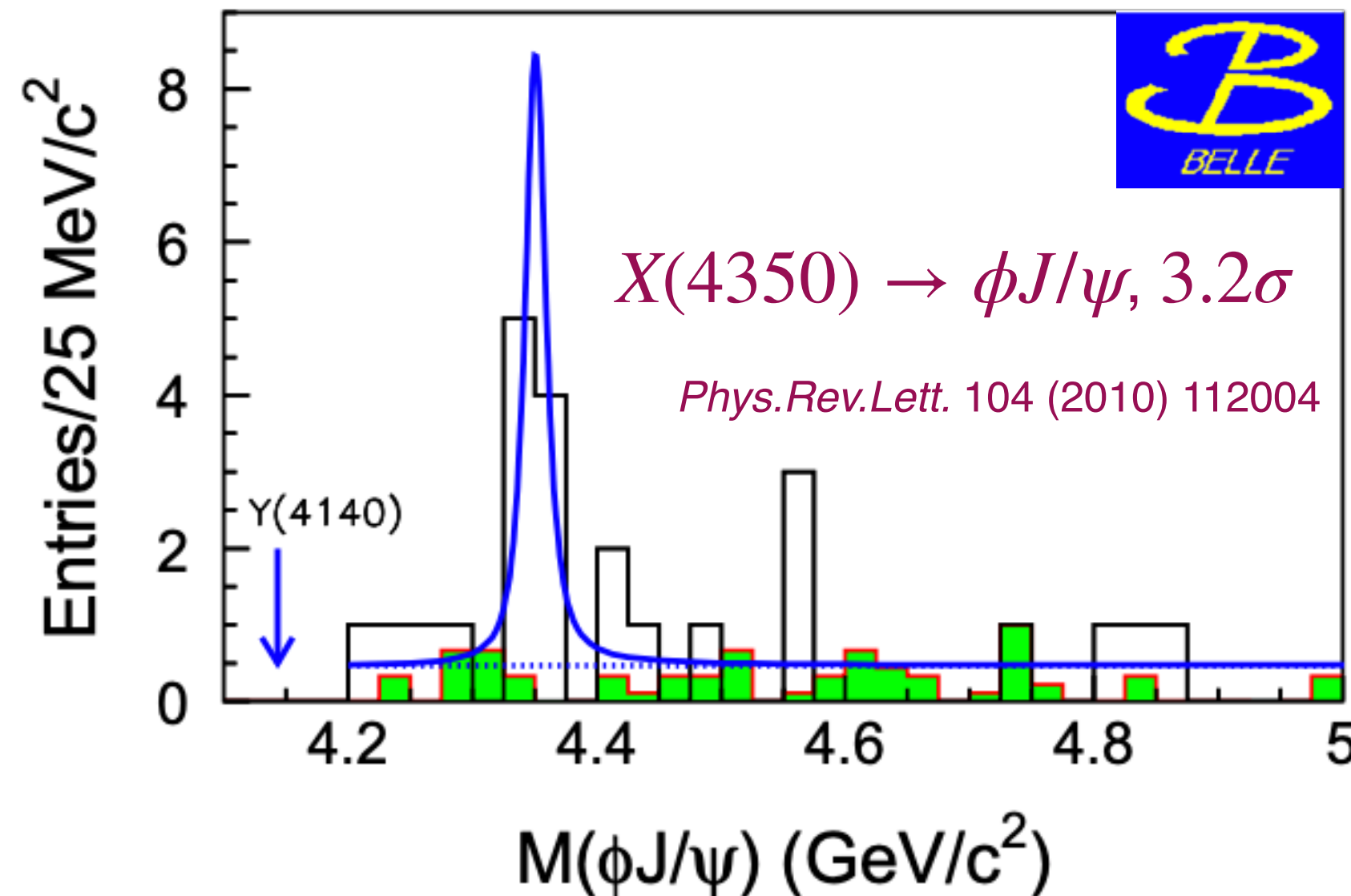
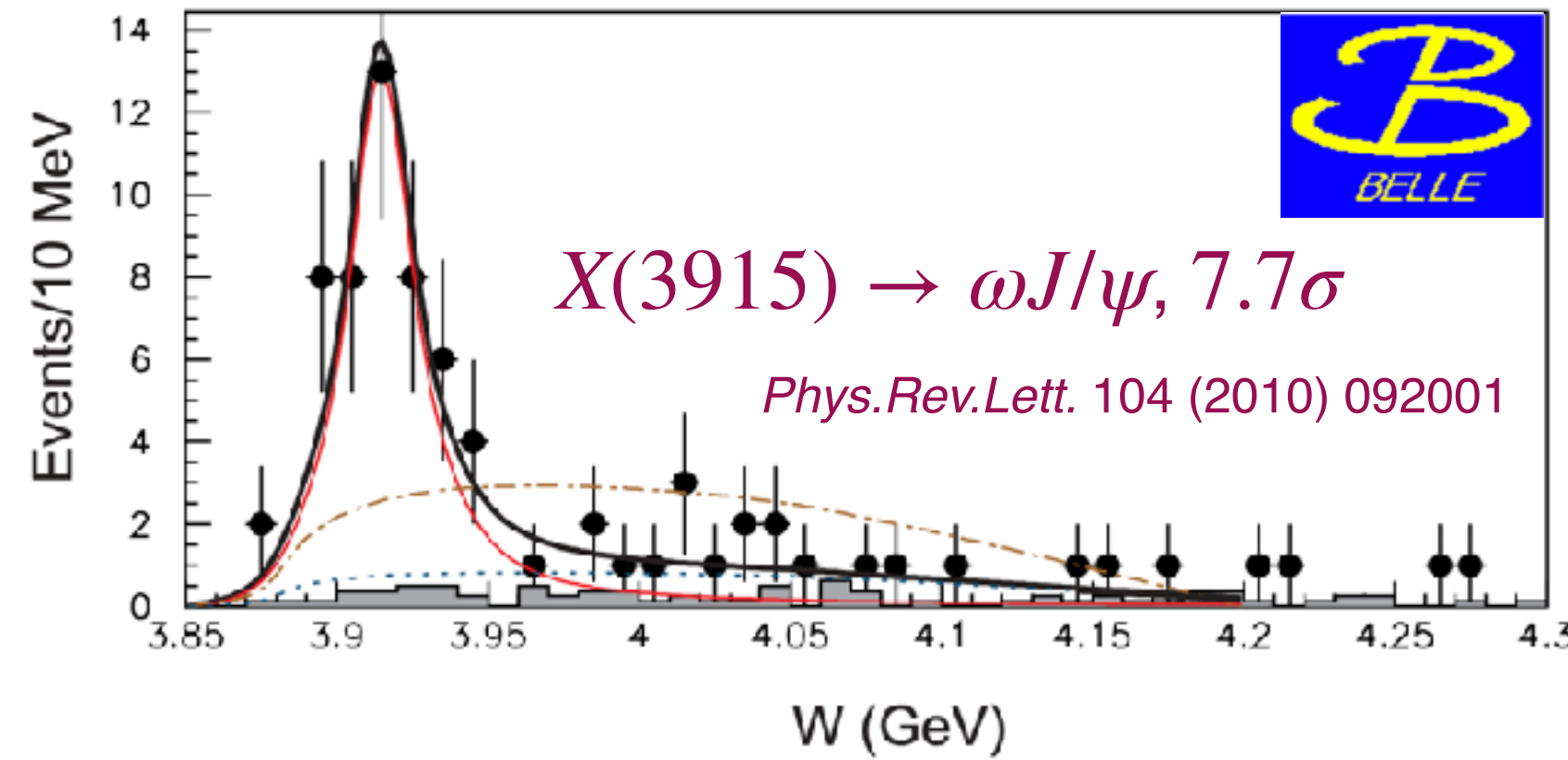
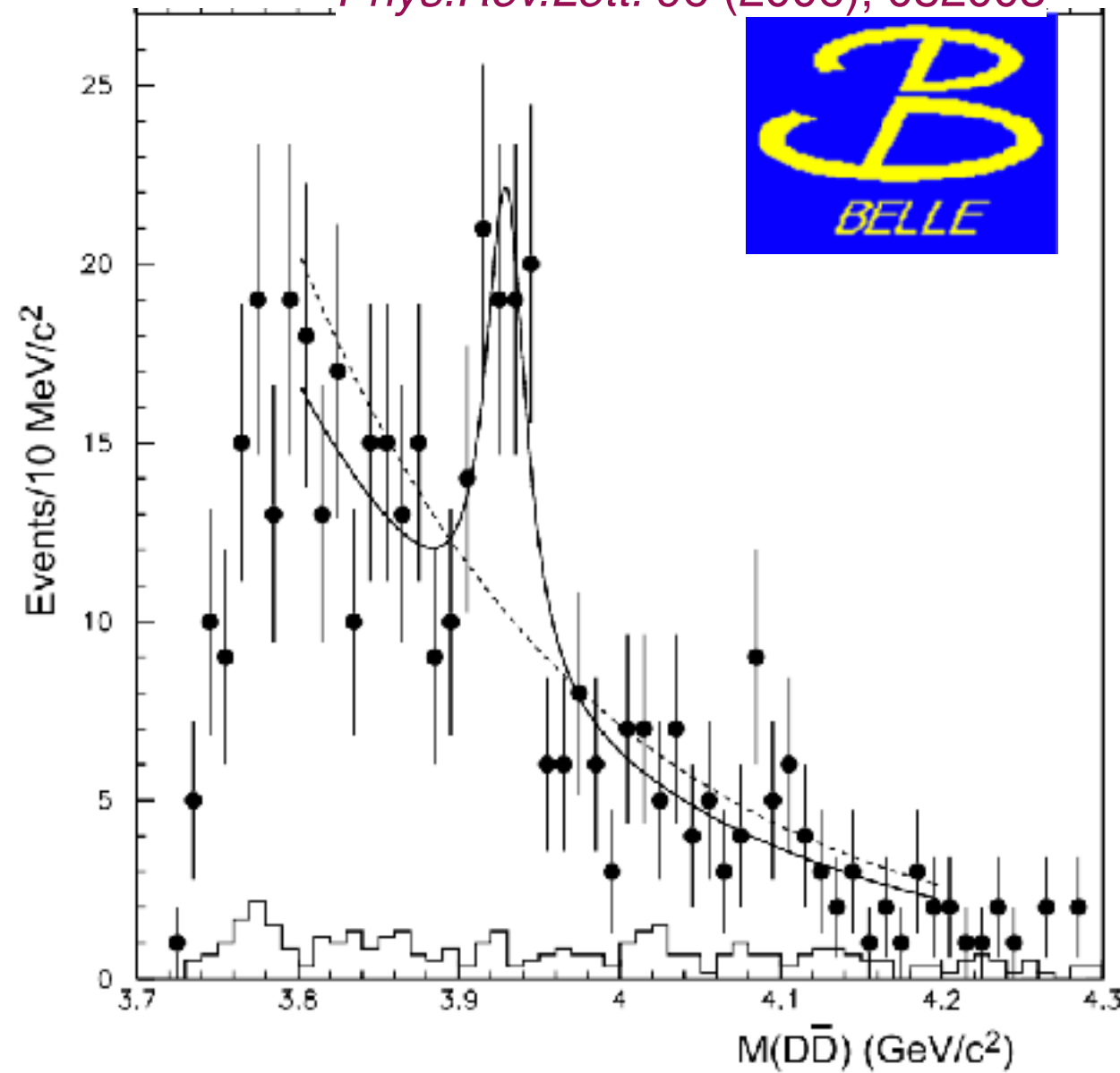
Two photon process

Unique field to produce charmonium(-like) particles

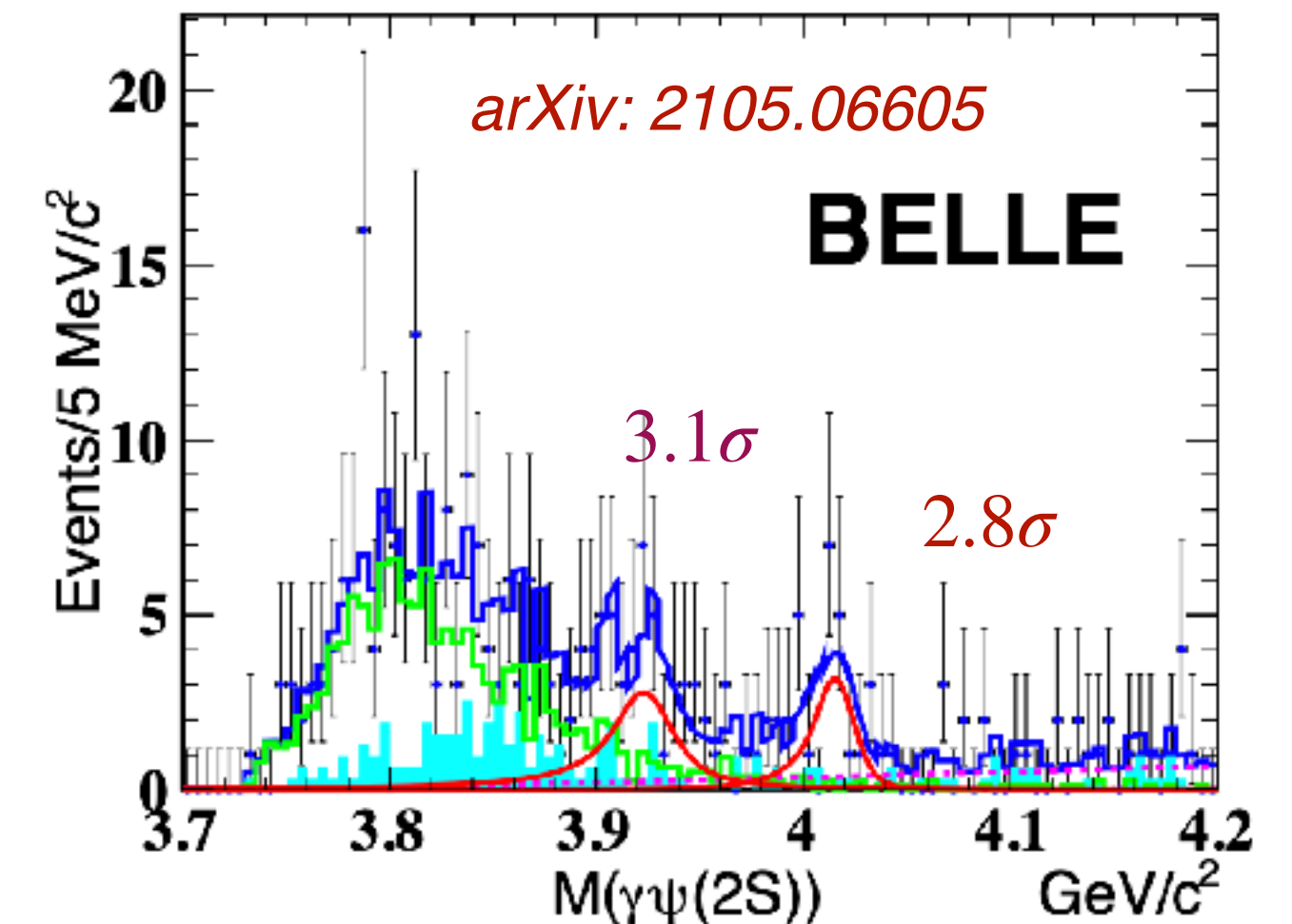
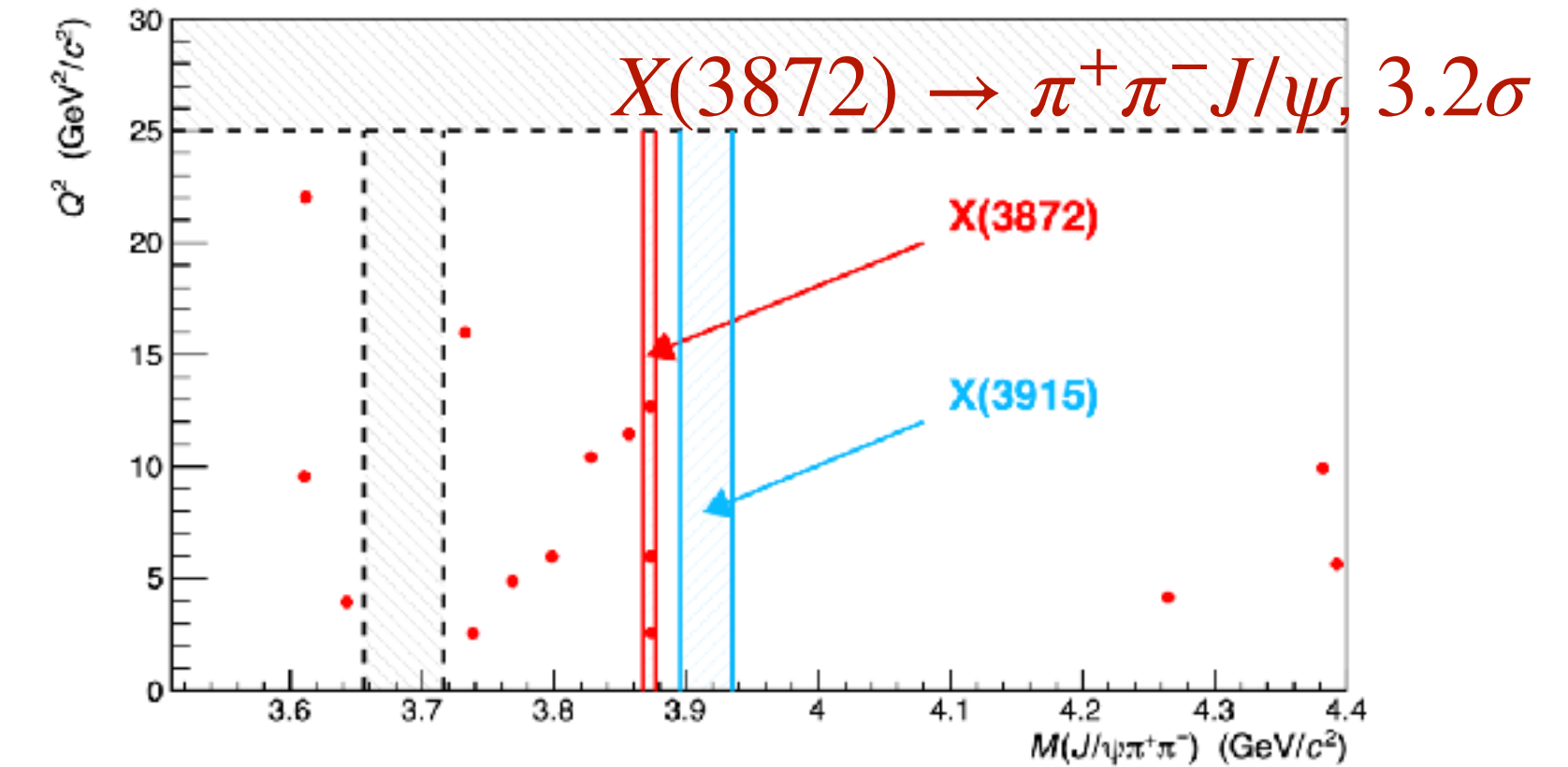


Establishment of exotic states in various final states

$X(3930) \rightarrow D\bar{D}, 5.3\sigma$
Phys.Rev.Lett. 96 (2006), 082003

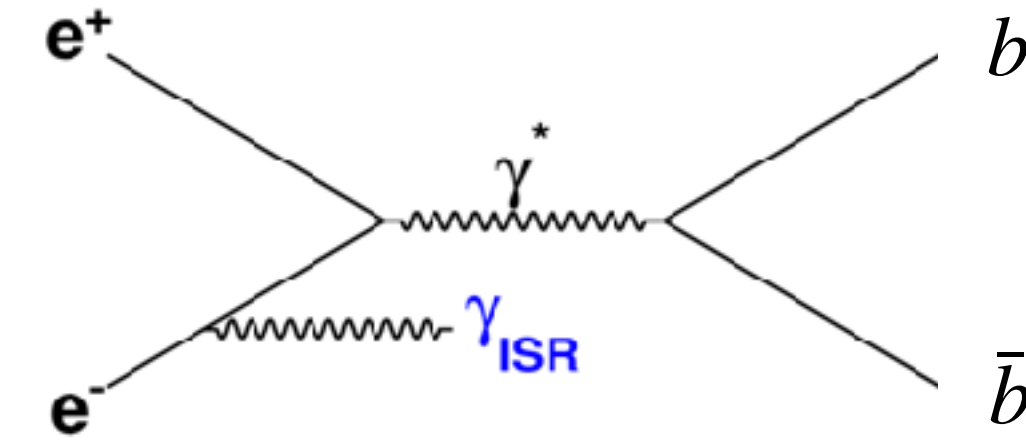


Phys.Rev.Lett. 126 (2021) 12, 122001



Bottomonium production

Excellent play ground of NRQCD & unique properties



Questions raised: $\Upsilon(5S)$ mass & abnormal transition rate

Phys.Rev.Lett. 100 (2008) 112001

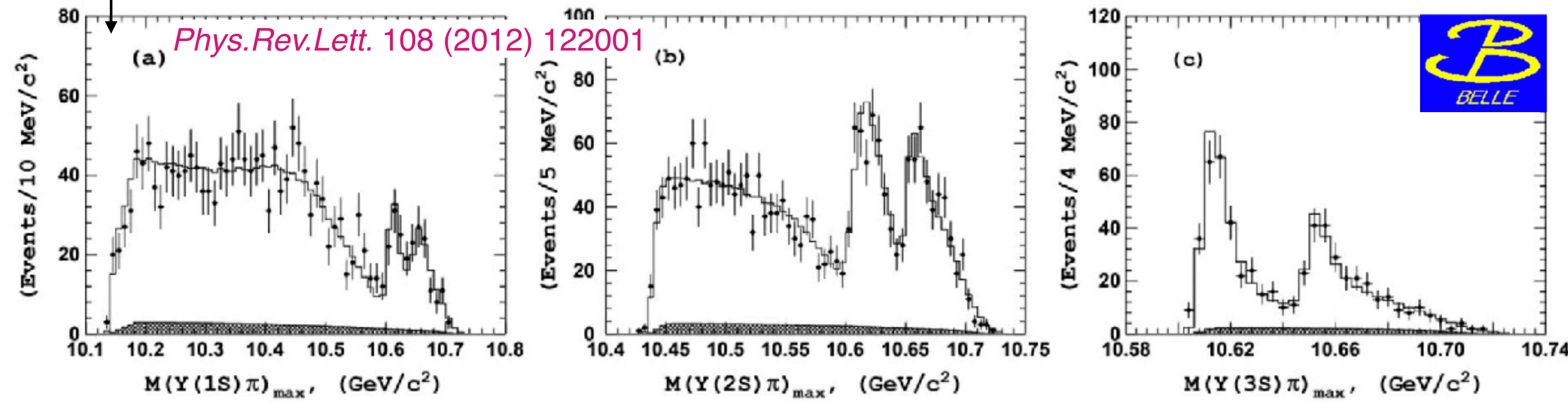
Process	Γ_{total}	$\Gamma_{e^+e^-}$	$\Gamma_{\Upsilon(1S)\pi^+\pi^-}$
$\Upsilon(2S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	0.032 MeV	0.612 keV	0.0060 MeV
$\Upsilon(3S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	0.020 MeV	0.443 keV	0.0009 MeV
$\Upsilon(4S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	20.5 MeV	0.272 keV	0.0019 MeV
$\Upsilon(10860) \rightarrow \Upsilon(1S)\pi^+\pi^-$	110 MeV	0.31 keV	0.59 MeV

$> 10^2 \times$

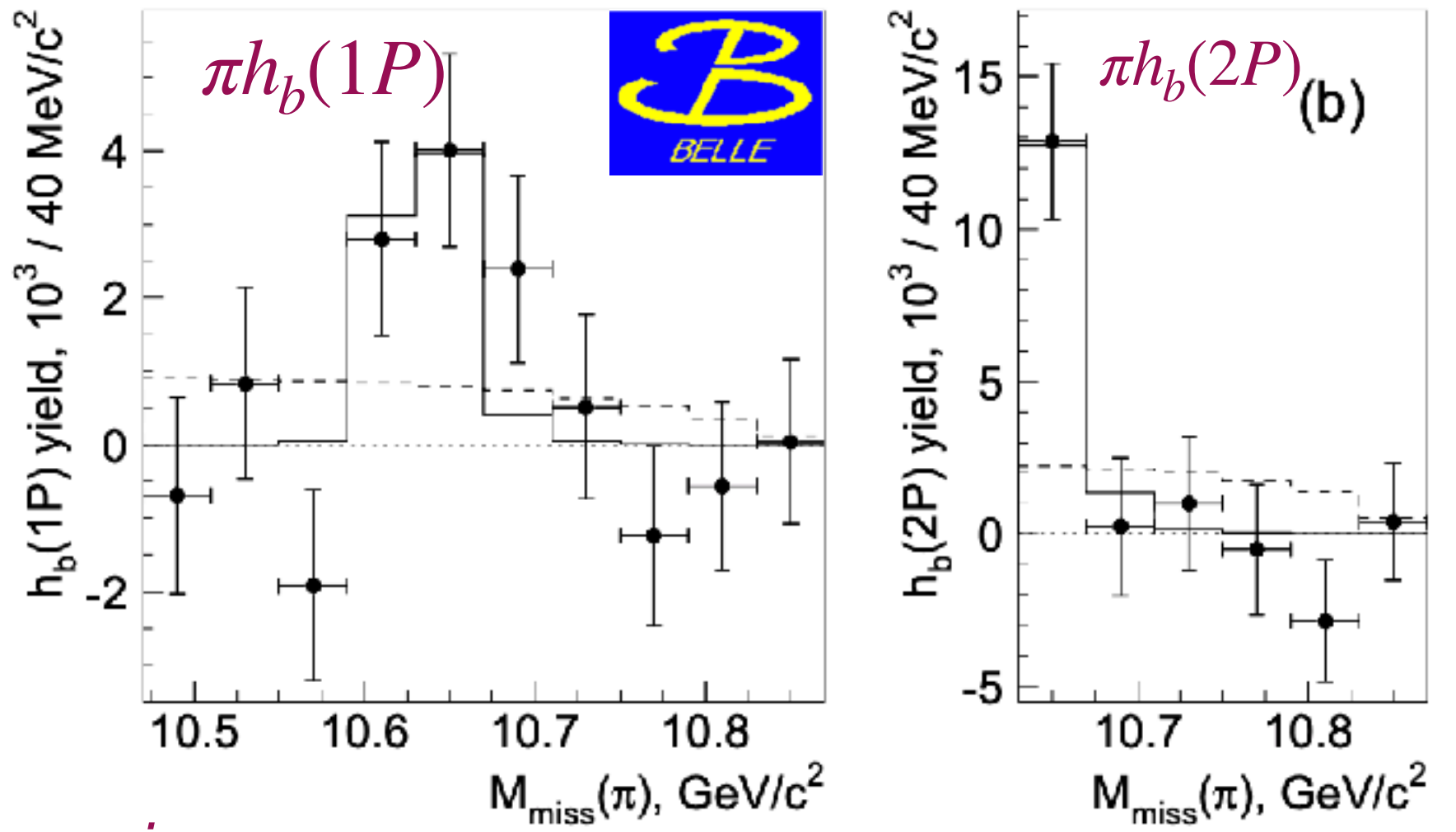
$$\frac{\Gamma(\Upsilon(5S) \rightarrow h_b(1P)\pi^+\pi^-)}{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(1S)\pi^+\pi^-)} = 0.46 \pm 0.08^{+0.07}_{-0.12}$$

$$\frac{\Gamma(\Upsilon(5S) \rightarrow h_b(2P)\pi^+\pi^-)}{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(2S)\pi^+\pi^-)} = 0.77 \pm 0.08^{+0.22}_{-0.17}$$

Observation of $Z_b(10610)^+$ and $Z_b(10650)^+$



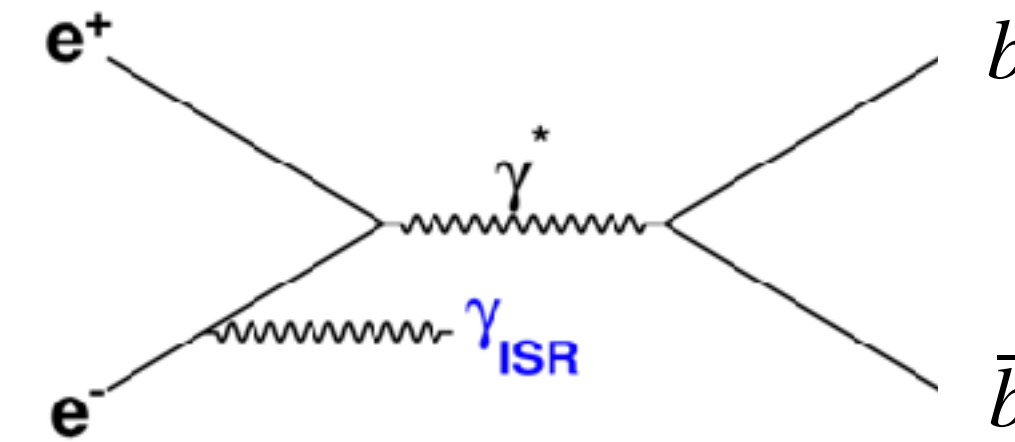
Phys.Rev.Lett. 117 (2016) 14, 142001



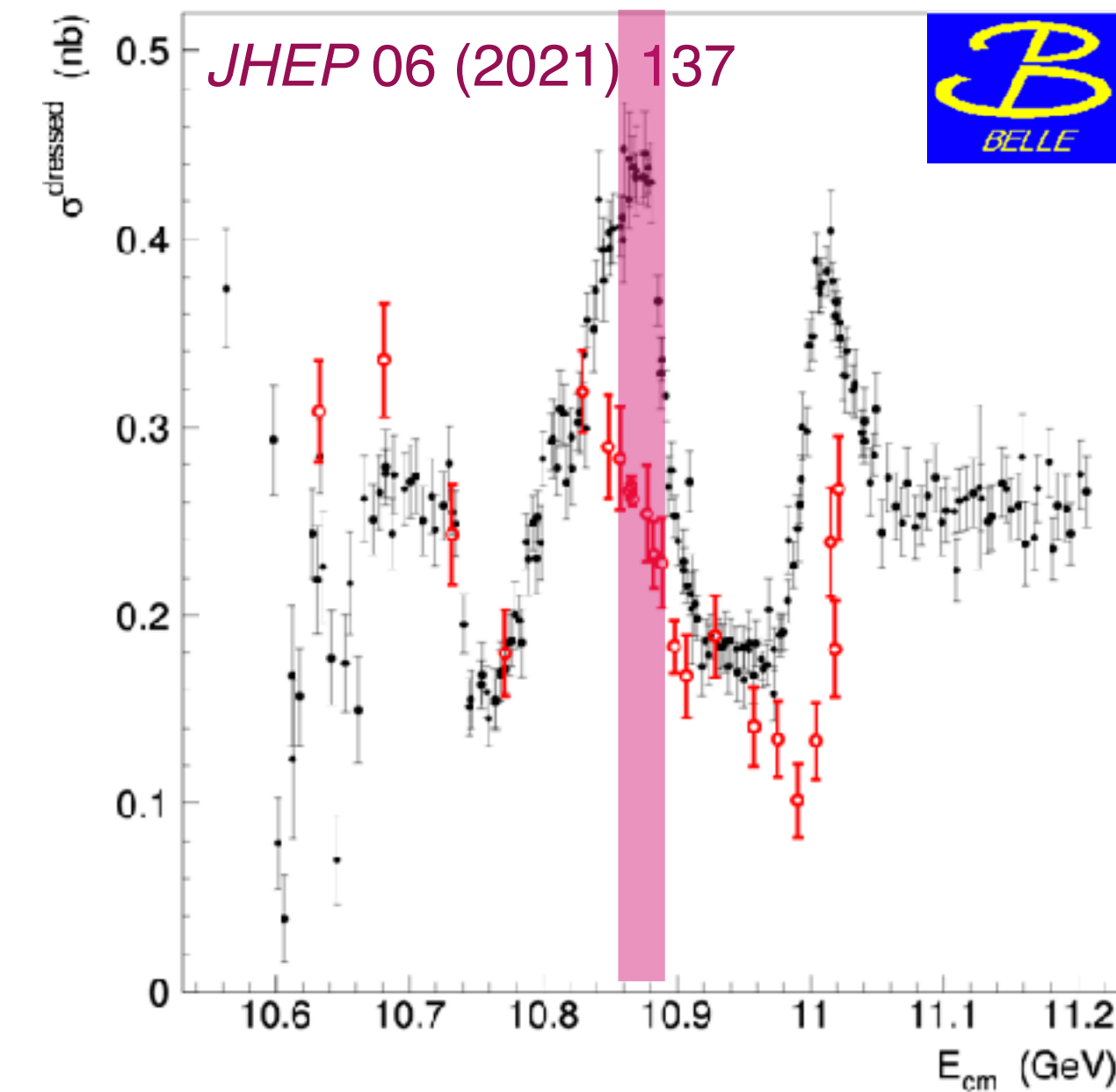
****Similar feature is also found in higher charmonia**

Bottomonium production

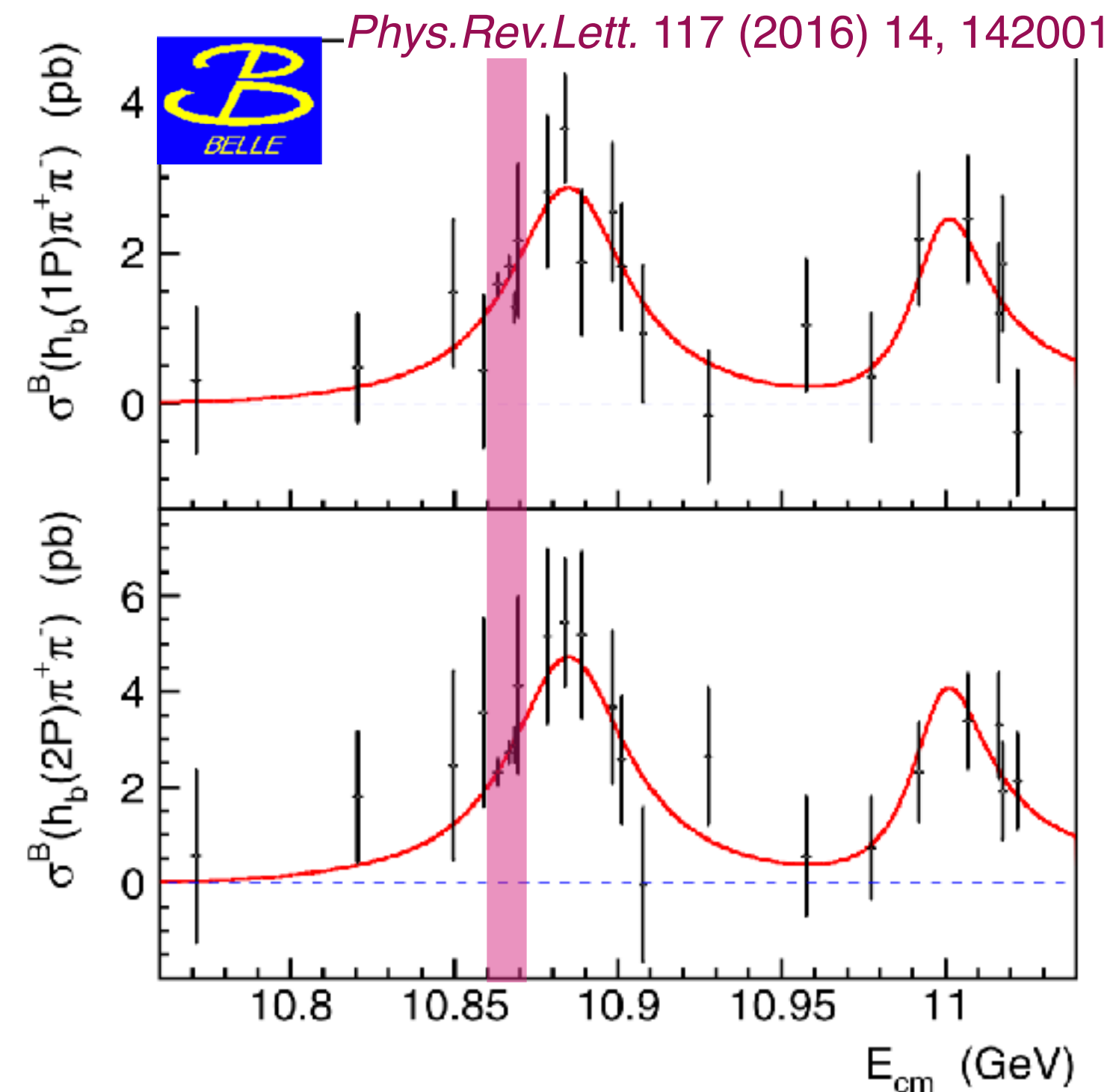
Excellent play ground of NRQCD & unique properties



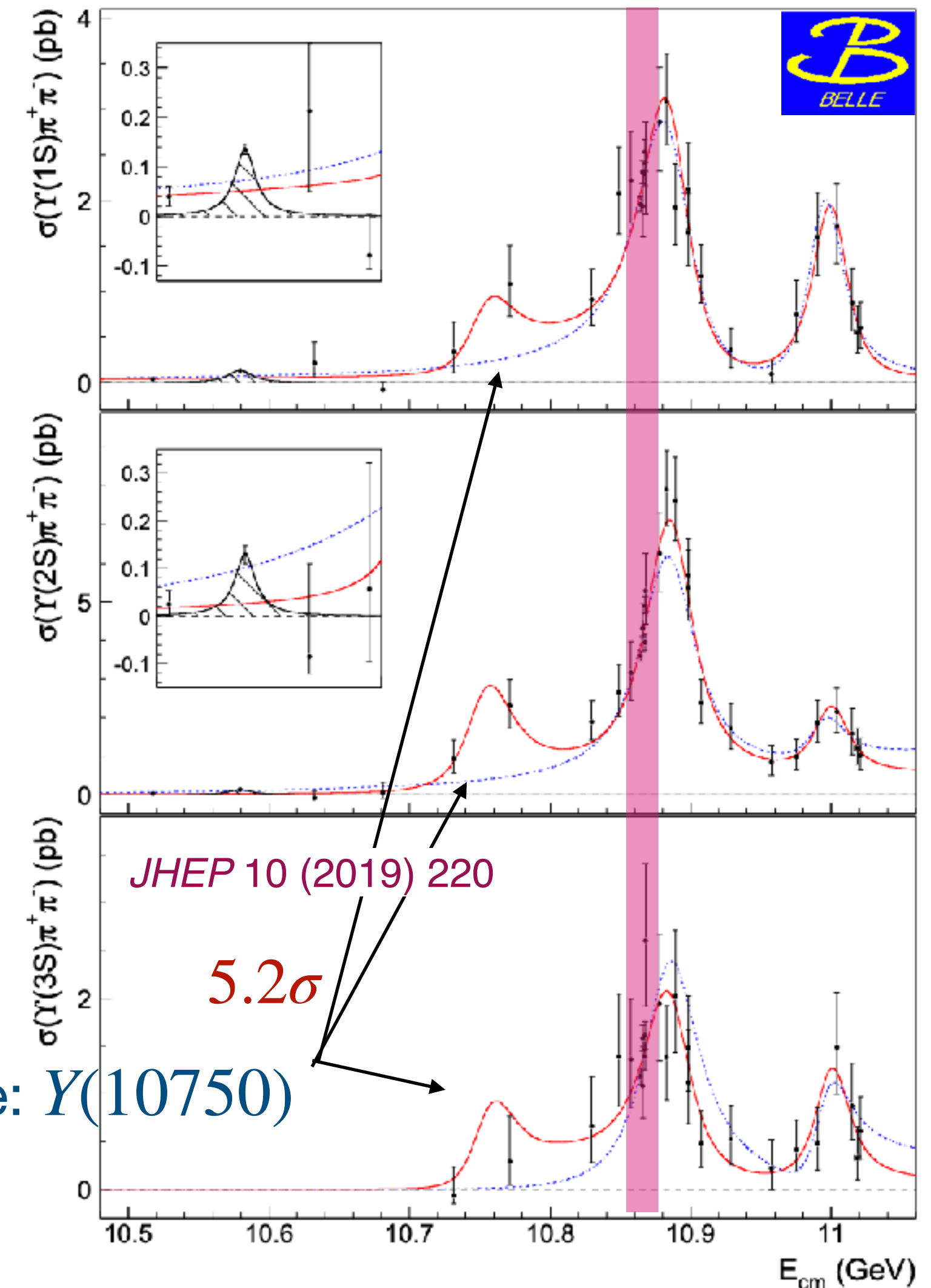
Questions raised: $\Upsilon(5S)$ mass & abnormal transition rate



Red dots: Measured $\sigma[e^+e^- \rightarrow B^{(*)}\bar{B}^{(*)}]$
 Black dots: total $b\bar{b}$ dressed cross section
Chin. Phys. C 44, 083001 (2020)



Phys.Rev.Lett. 117 (2016) 14, 142001



JHEP 10 (2019) 220

A new state: $Y(10750)$

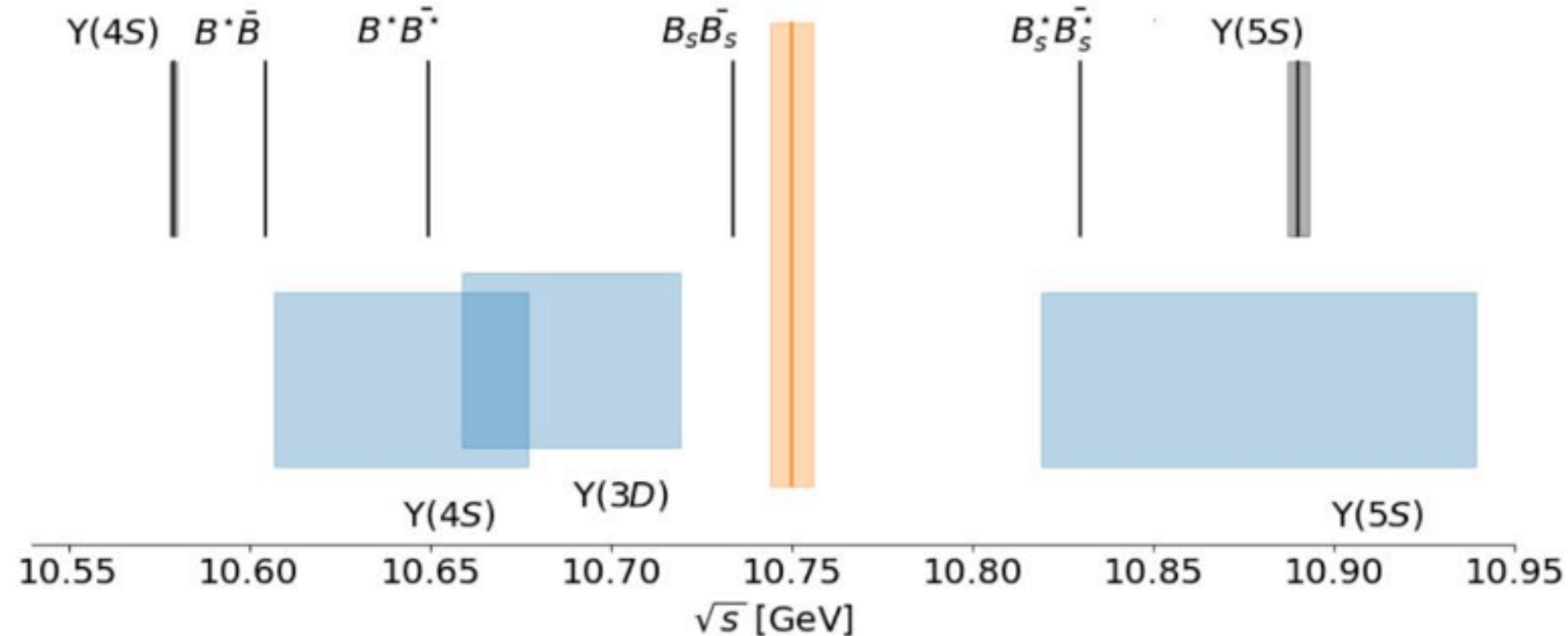
- ◆ Excess between $e^+e^- \rightarrow b\bar{b}$ and $B^{(*)}\bar{B}^{(*)}$, others?
- ◆ Cross sections do not peak at $\Upsilon(5S)$ mass.

Y(10753): why it's important



Uncertain nature:

- No clear conventional $b\bar{b}$ candidate
- Molecule? 10.75 GeV isn't a threshold...
- Tetraquark?



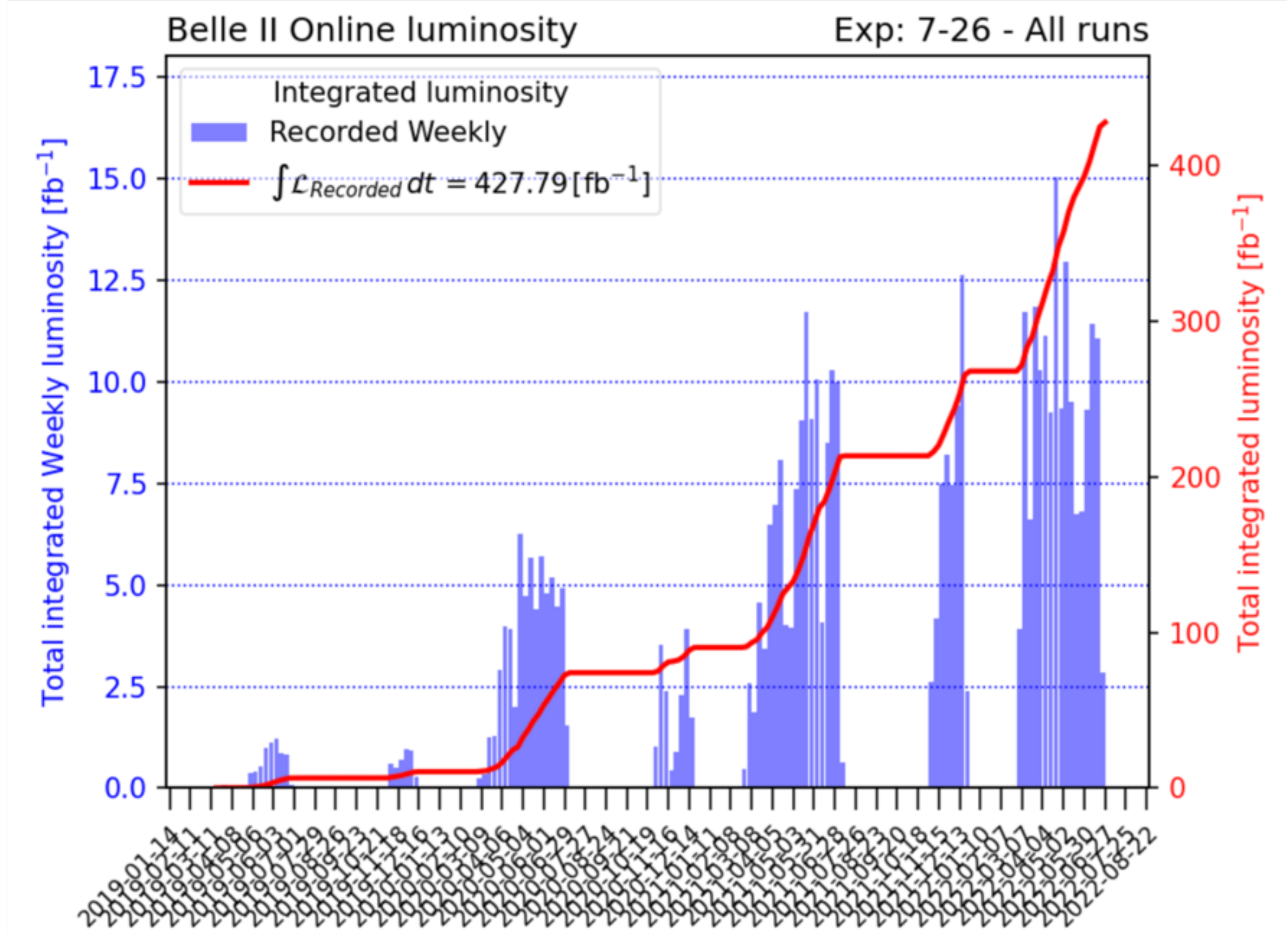
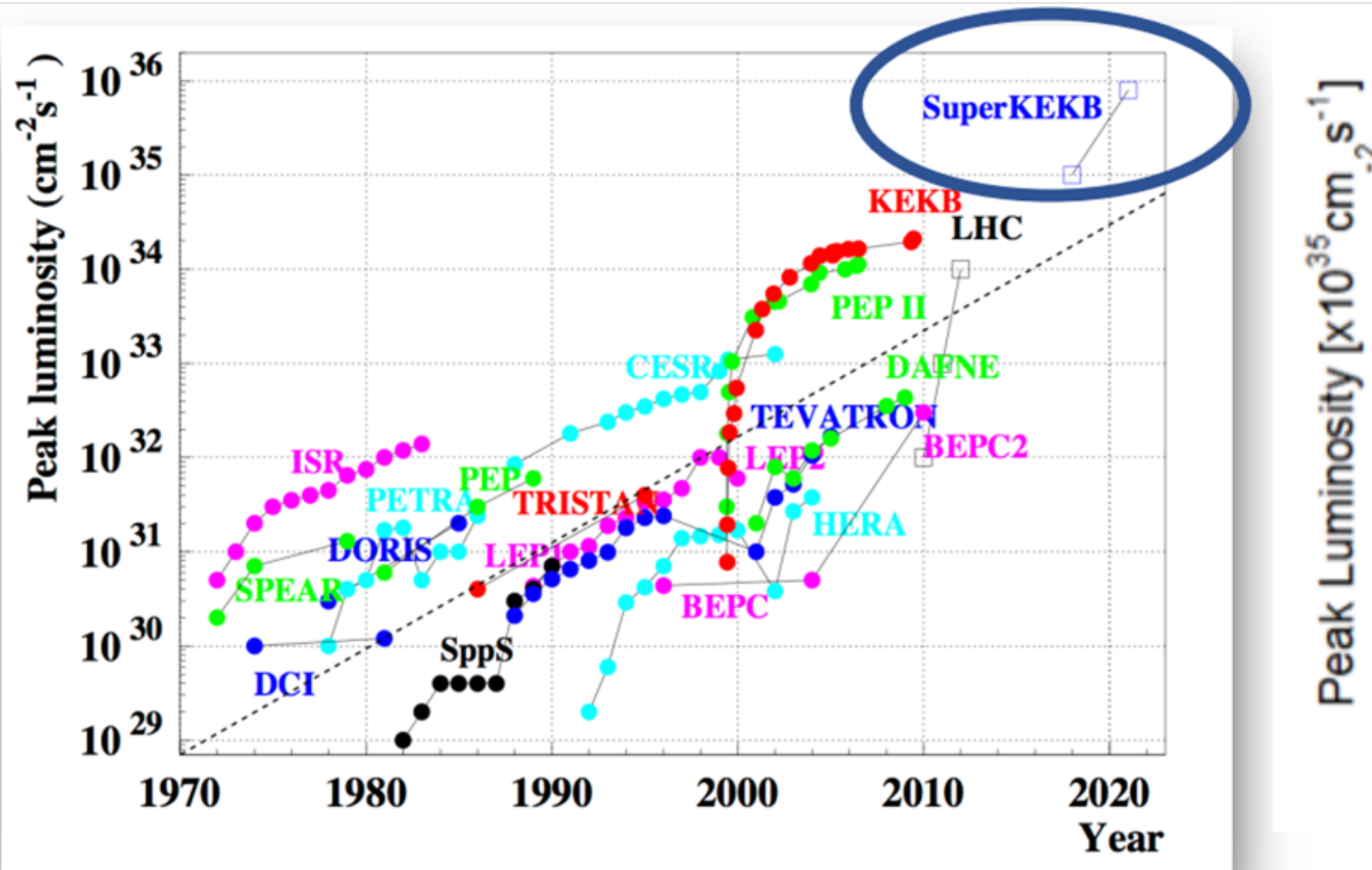
Conventional interpretations:

- Chen, Zhang & He, PRD 101, 014020 (2020)
- Giron & Lebed, PRD 102, 014036 (2020)
- Li et al., EPJC 80, 59, (2020)
- Li et al., PRD 104, 034036 (2021)
- van Beveren & Oset, PPNP 117, 103845 (2021)
- Bai et al., PRD 105, 074007 (2022)
- Husken, Mitchell & Swanson, arXiv:2204.11915 (2022)
- Kher et al., EPJ+ 137, 357 (2022)
- Li, Bai & Liu, arXiv:2205.04049 (2022)
- Liang, Ikeno & Oset, PLB 803, 135340 (2020)
- ...

Exotic interpretations:

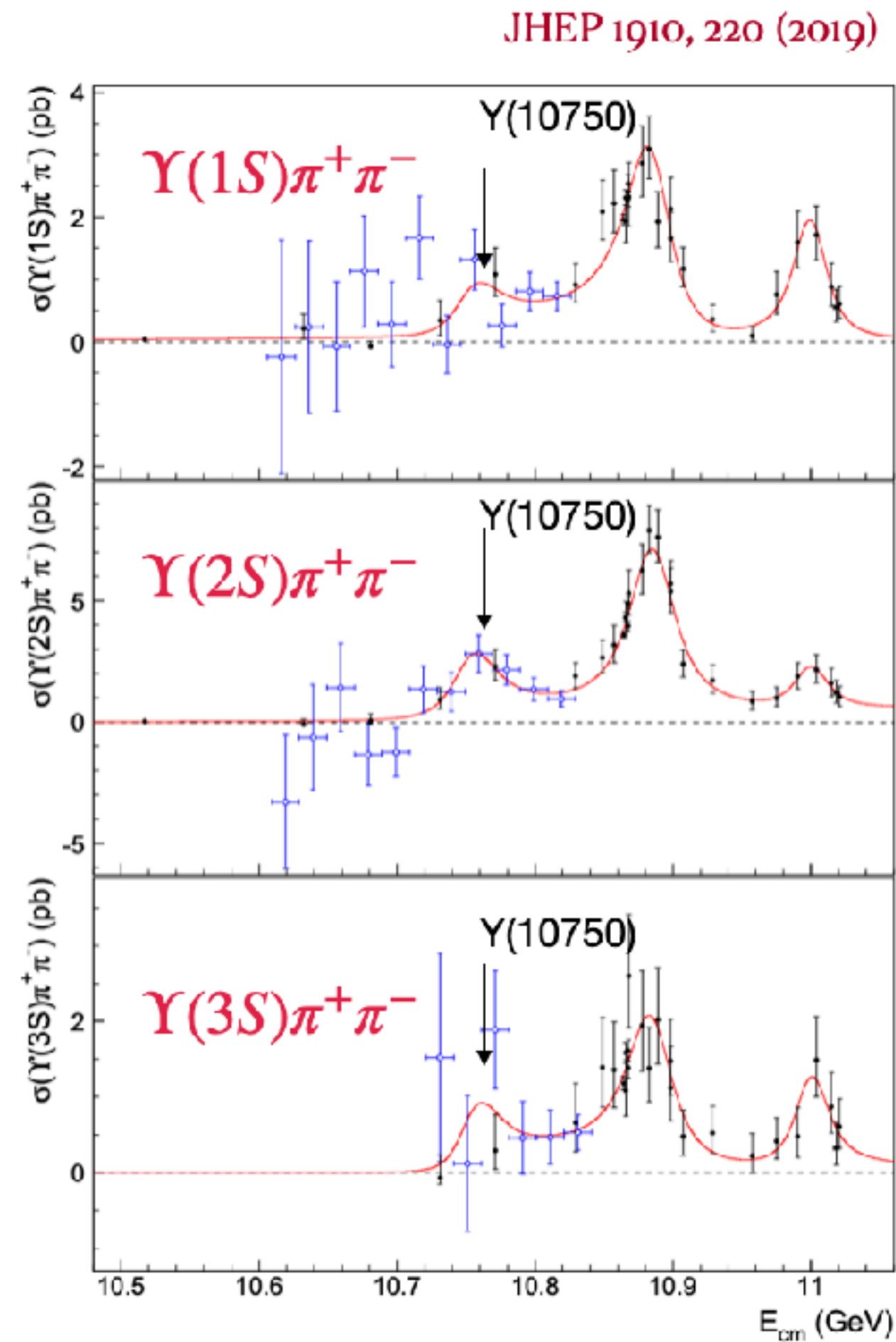
- Wang, CPC 43, 123102 (2019)
- Ali, Maiani, Parkhomenko & Wang, PLB 802, 135217 (2020)
- Bicudo, Cardoso & Wagner, PRD 103, 074507 (2020)
- Castella & Passemar, PRD 104, 034019 (2021)
- ...

Belle II luminosity High precision frontier

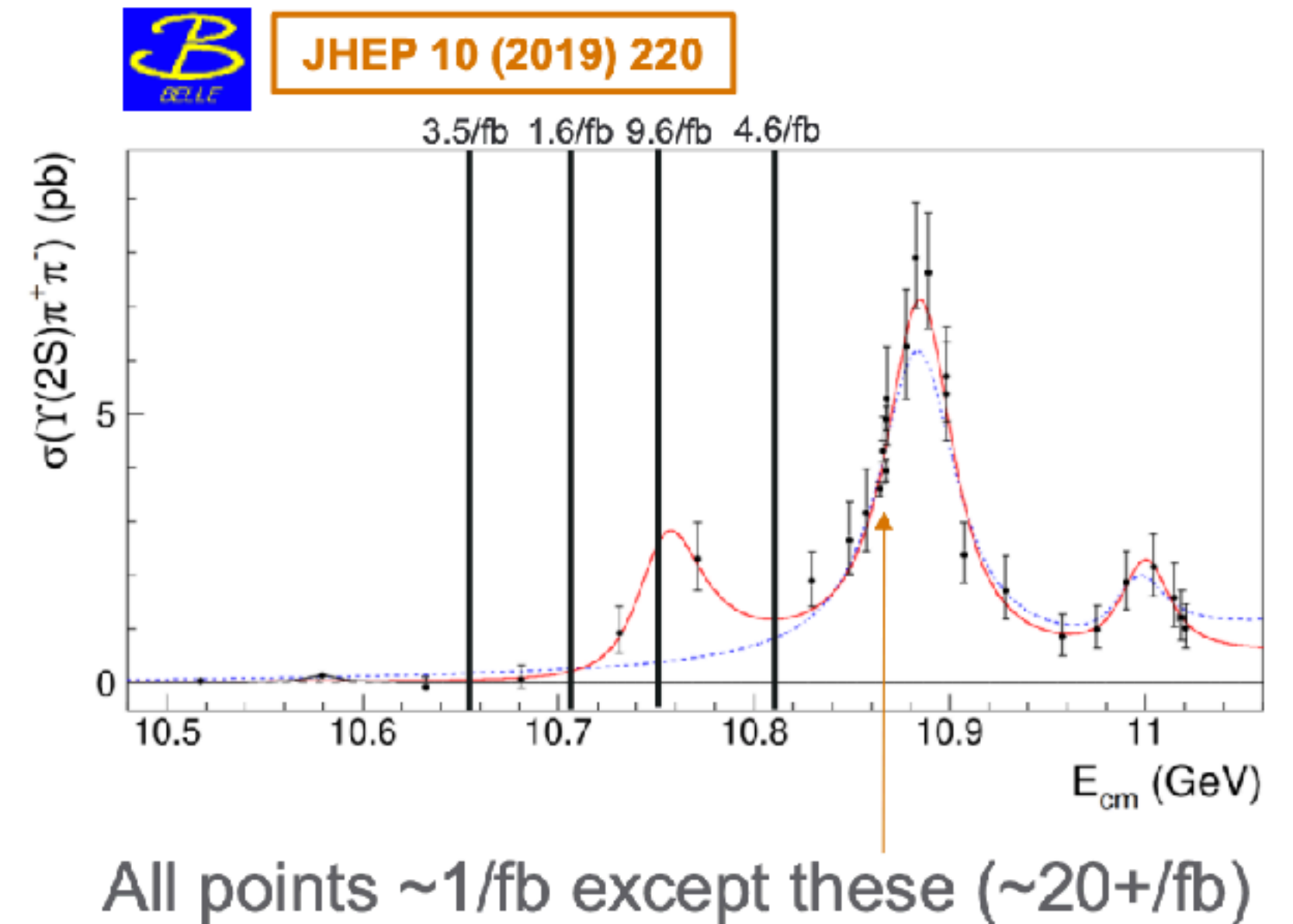
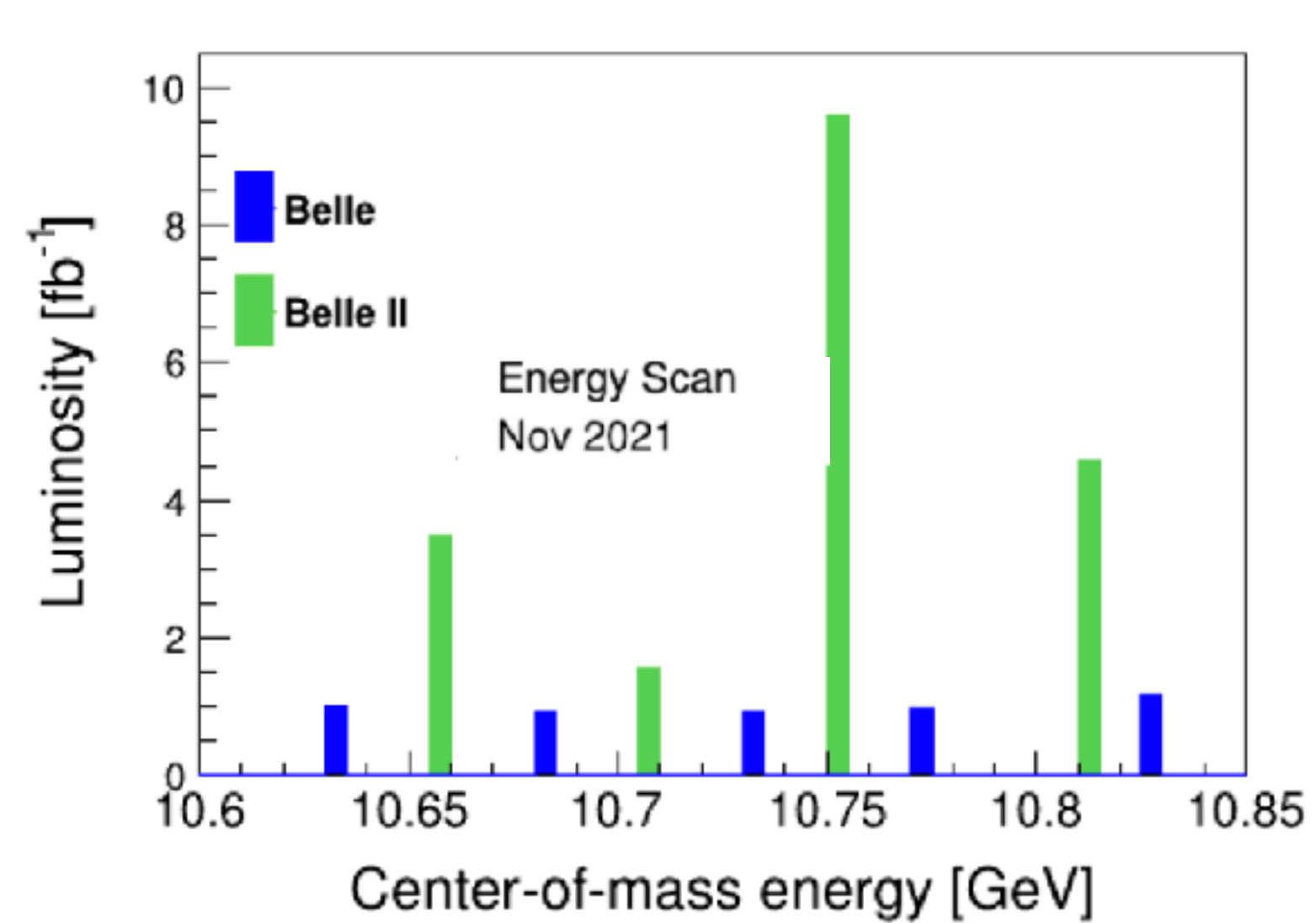


Belle II already achieve the world record instantaneous luminosity: $4.7 \times 10^{34} / \text{cm}^2 / \text{s}$

A little data may tell a big story



Scan data samples with 1 fb^{-1} at each point on Belle



All points $\sim 1/\text{fb}$ except these ($\sim 20+/\text{fb}$)

- Data collected at 4 energy points around 10.75 GeV
- Physics goal: understand the nature of $Y(10750)$ energy region

Belle II energy scan: new result



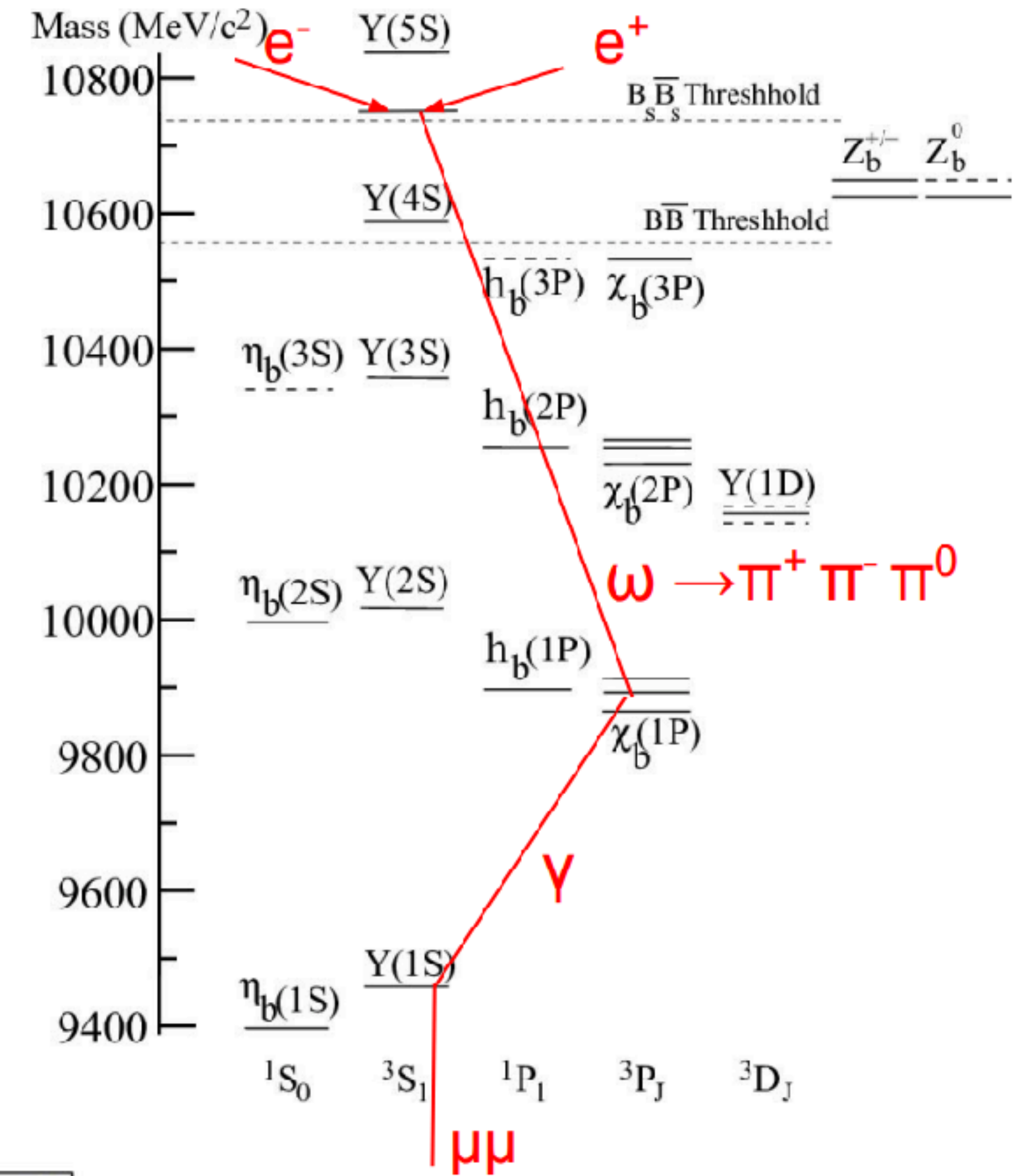
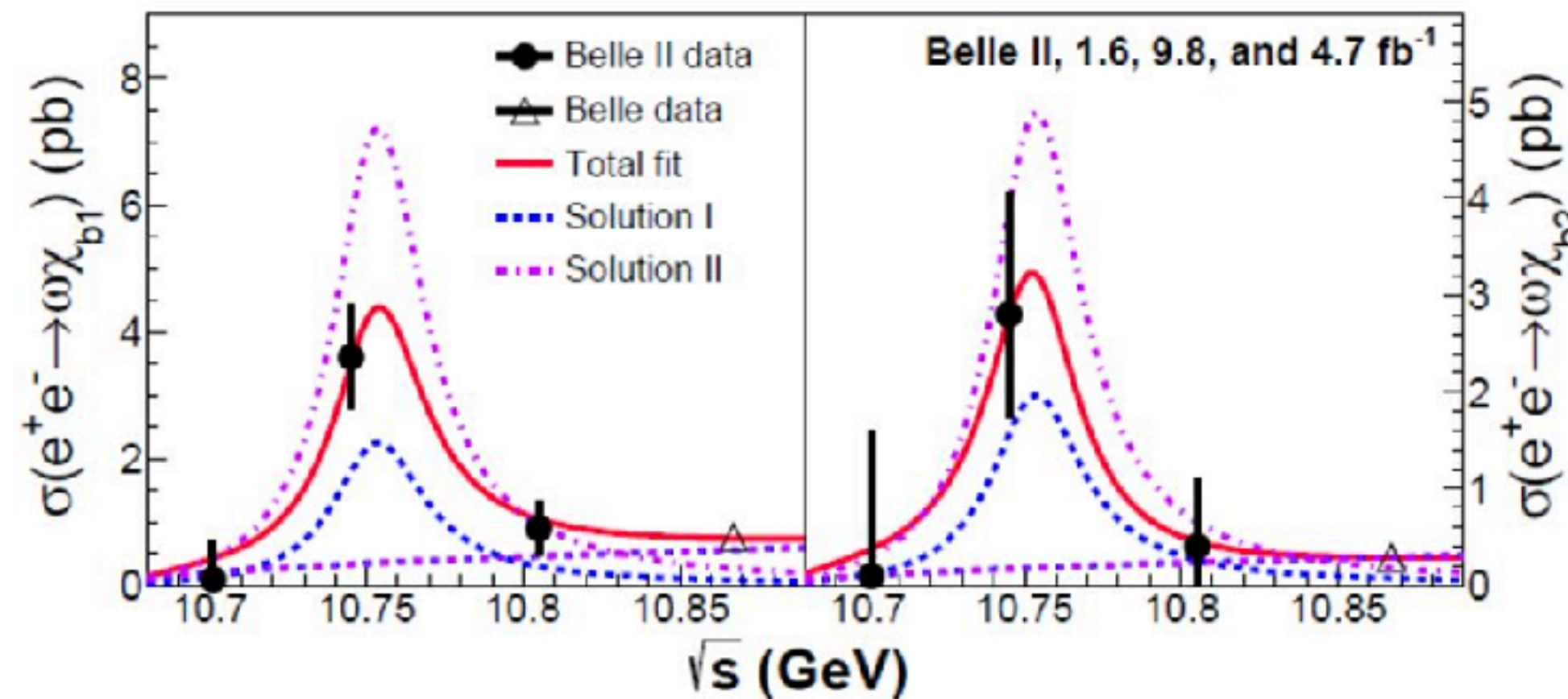
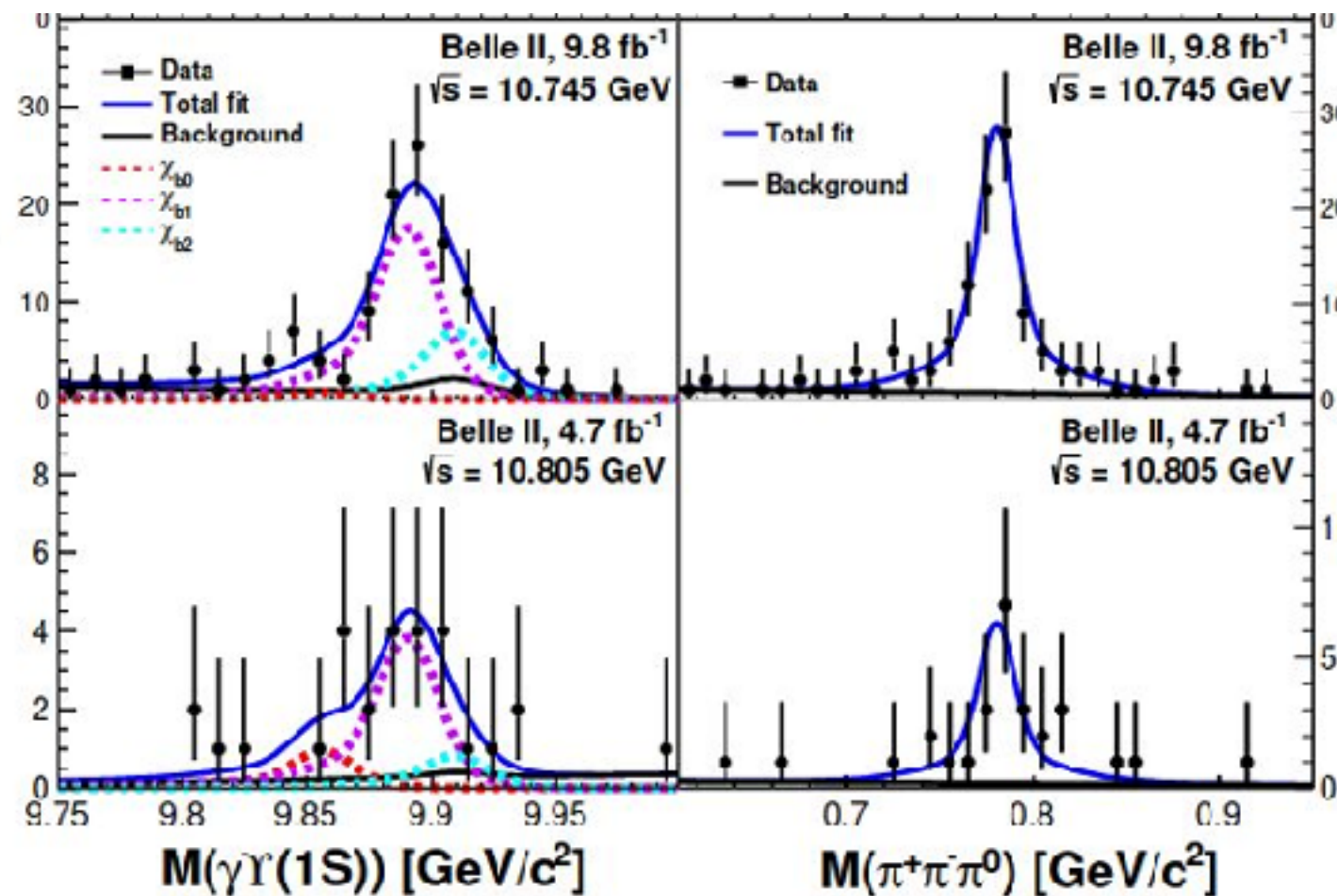
Observation of $e^+e^- \rightarrow \omega\chi_{bJ}(1P)$ and search for $X_b \rightarrow \omega\Upsilon(1S)$ at \sqrt{s} near 10.75 GeV

I. Adachi, L. Aggarwal, H. Ahmed, H. Aihara, N. Akopov, A. Aloisio, N. Anh Ky, T. Aushev, V. Aushev, H. Bae, P. Bambade, Sw. Banerjee, J. Baudot, M. Bauer, A. Beaubien, J. Becker, P. K. Behera, J. V. Bennett, E. Bernieri, F. U. Bernlochner, V. Bertacchi, M. Bertemes, E. Bertholet, M. Bessner, S. Bettarini, B. Bhuyan, F. Bianchi, T. Bilka, D. Biswas, D. Bodrov, A. Bolz, A. Bondar, J. Borah, A. Bozek, M. Bračko, P. Branchini, T. E. Browder, A. Budano, S. Bussino, M. Campajola, L. Cao, G. Casarosa, M.-C. Chang, P. Cheema, V. Chekelian, Y. Q. Chen, K. Chilikin, K. Chirapatpimol, H.-E. Cho, K. Cho, S.-J. Cho, S.-K. Choi, S. Choudhury, D. Cinabro, L. Corona, S. Cunliffe, S. Das, F. Dattola, E. De La Cruz-Burclo, S. A. De La Motte, G. De Nardo, M. De Nuccio, G. De Pietro, R. de Sangro, M. Destefanis, S. Dey, A. De Yta-Hernandez, R. Dhamija, A. Di Canto, F. Di Capua, Z. Doležal, I. Domínguez Jiménez, T. V. Dong, M. Dorigo, K. Dort, S. Dreyer, S. Dubey, G. Dujany, M. Eliachevitch, D. Epifanov, P. Feichtinger, T. Ferber, D. Ferlewicz, T. Fillinger, G. Finocchiaro, A. Fodor, F. Forti, B. G. Fulsom, ...

(The Belle II Collaboration)

We study the processes $e^+e^- \rightarrow \omega\chi_{bJ}(1P)$ ($J = 0, 1, \text{ or } 2$) using samples at center-of-mass energies $\sqrt{s} = 10.701, 10.745, \text{ and } 10.805$ GeV, corresponding to 1.6, 9.8, and 4.7 fb^{-1} of integrated luminosity, respectively. These data were collected with the Belle II detector during a special run of the SuperKEKB collider above the $\Upsilon(4S)$ resonance. We report the first observation of $\omega\chi_{bJ}(1P)$ signals at $\sqrt{s} = 10.745$ GeV. By combining Belle II data with Belle results at $\sqrt{s} = 10.867$ GeV, we find energy dependencies of the Born cross sections for $e^+e^- \rightarrow \omega\chi_{b1,b2}(1P)$ to be consistent with the shape of the $\Upsilon(10753)$ state. Including data at $\sqrt{s} = 10.653$ GeV, we also search for the bottomonium equivalent of the $X(3872)$ state decaying into $\omega\Upsilon(1S)$. No significant signal is observed for masses between 10.45 and 10.65 GeV/c^2 .

[arXiv:2208.13189](https://arxiv.org/abs/2208.13189) [hep-ex]



Confirm the existence of $\Upsilon(10753)$.
Find a new decay mode.

How to identify exotics

- Exotic quantum numbers
 - $J^{PC}=0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, \dots$
 - $I=2, 5/2$
- Minimum quark content >3
 - $s s \bar{u} \bar{d}, c s \bar{u} \bar{d}, c c \bar{u} \bar{d}, u u d d \bar{s}, u u d d \bar{c}, \dots$
- Too many states, beyond expectation of quark model
 - Vector charmonia: $N_{\Upsilon} + N_{\Psi} > N_{\Psi(3S1)} + N_{\Psi(3D1)}$
- State with exotic properties
 - Large mass but very narrow
 - Very close to threshold
 - Strong coupling to heavy quarkonium but charged
 - ...



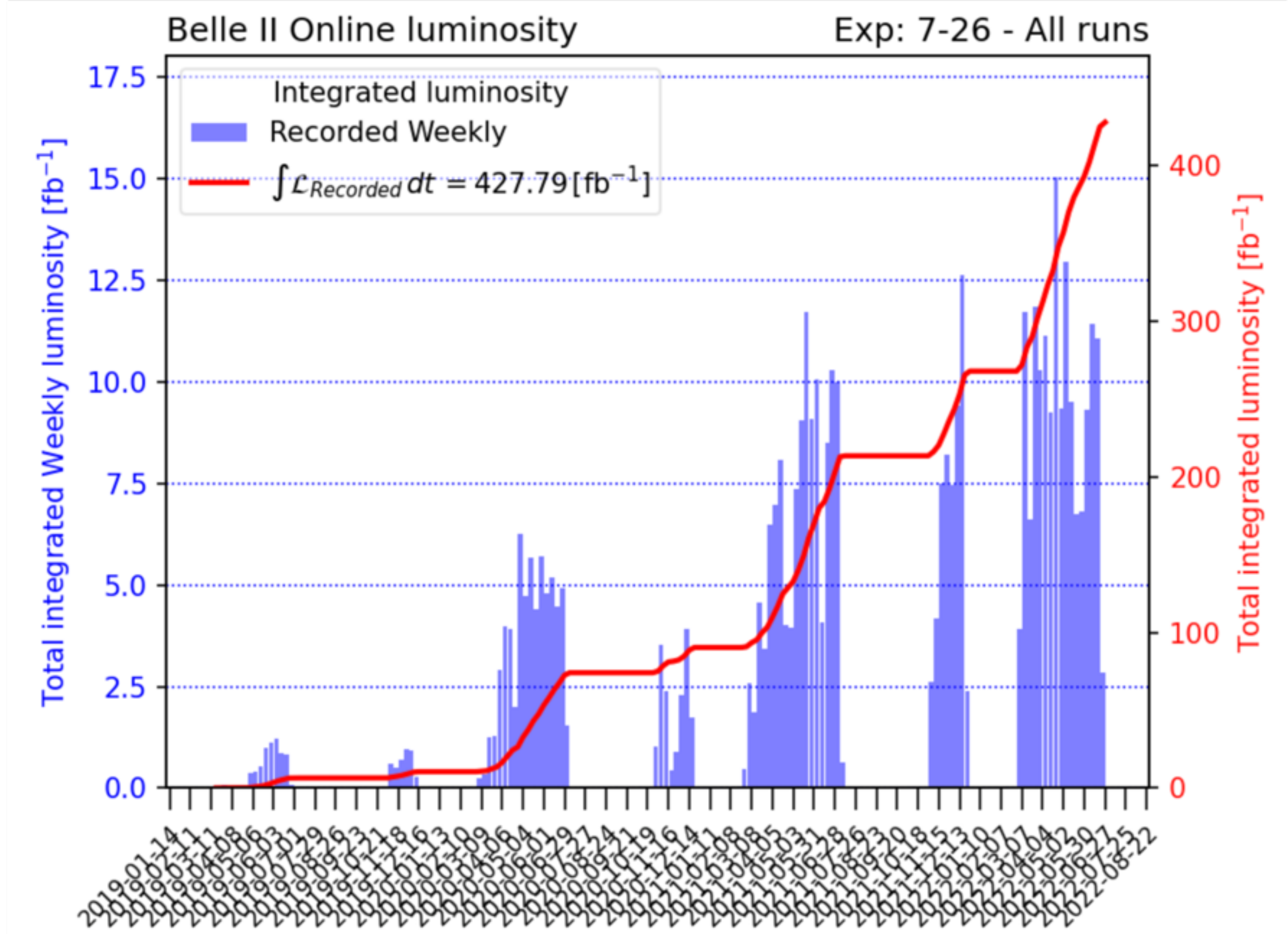
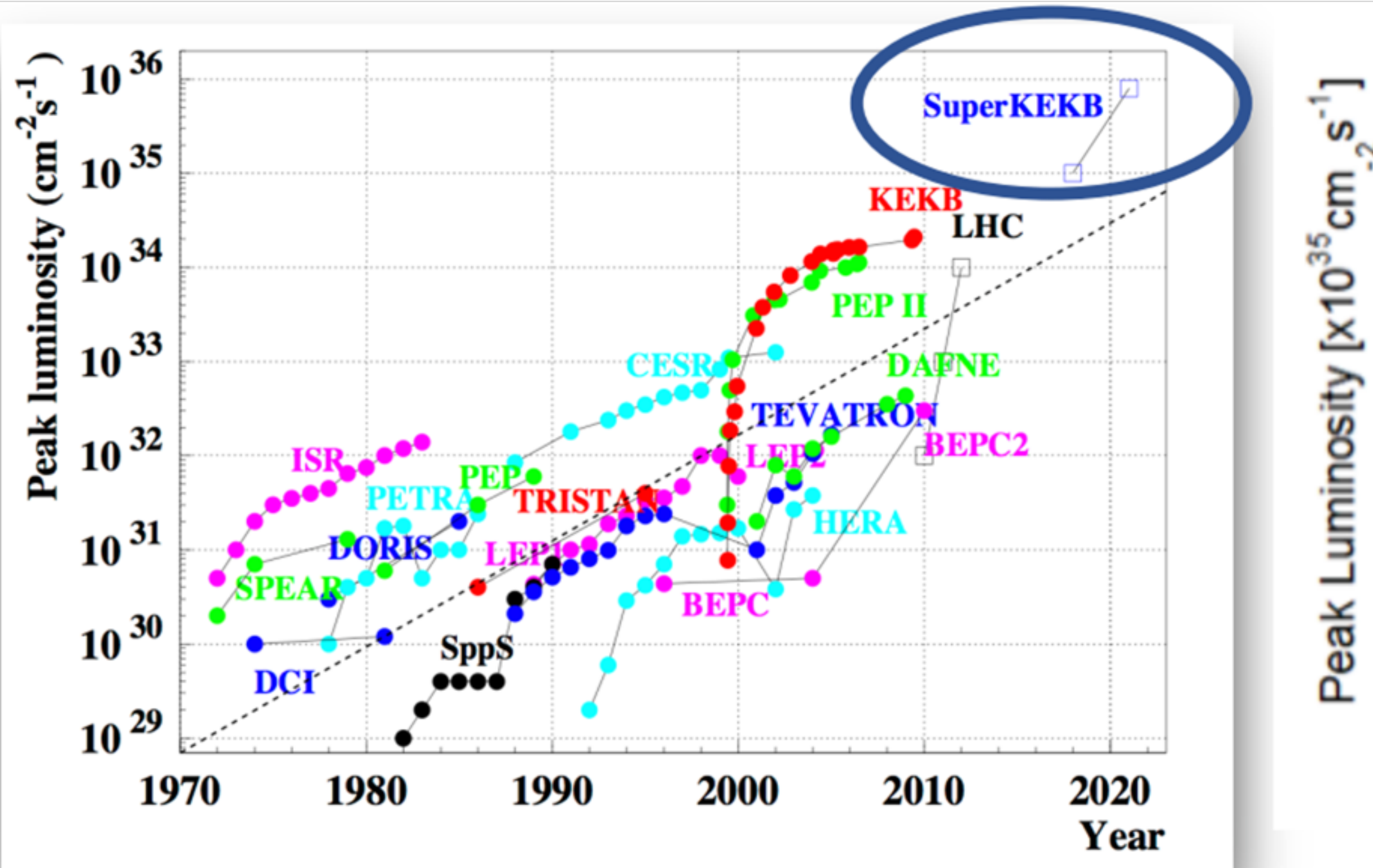
Not so easy...

Summary

- Quarkonium spectroscopy is an extraordinary place to test the QCD theory, especially the observation of exotic states, which gives us new insight and shed lights to the understanding of QCD theory.
- Belle II, the next generation B-factory, can make significant impacts in quarkonia study.
 - ◆ Precisely measure lineshapes.
 - ◆ Determine spin-parities, transitions, and quantum numbers.
 - ◆ Search for new decay channels.
 - ◆ Test predictions for unobserved states.
 - ◆ Unique datasets.
 - ◆ ...
- The effort goes on with the upgraded facility, SuperKEKB collider, and Belle II detector.

BACK UP

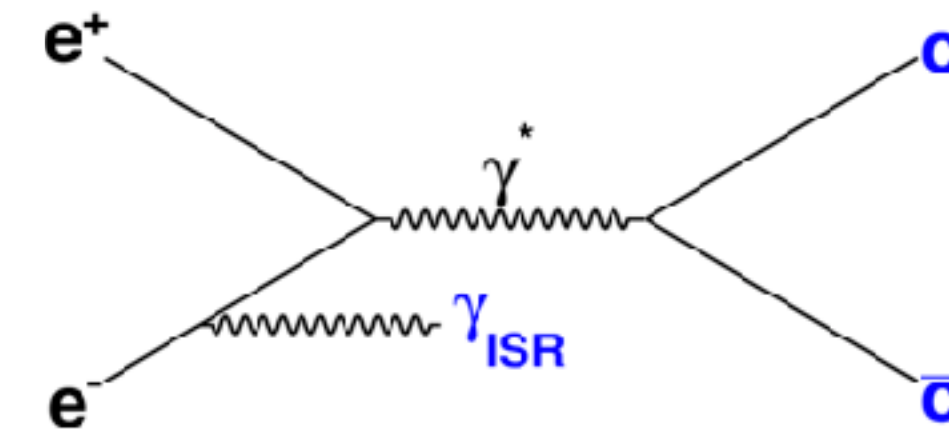
Belle II luminosity High precision frontier



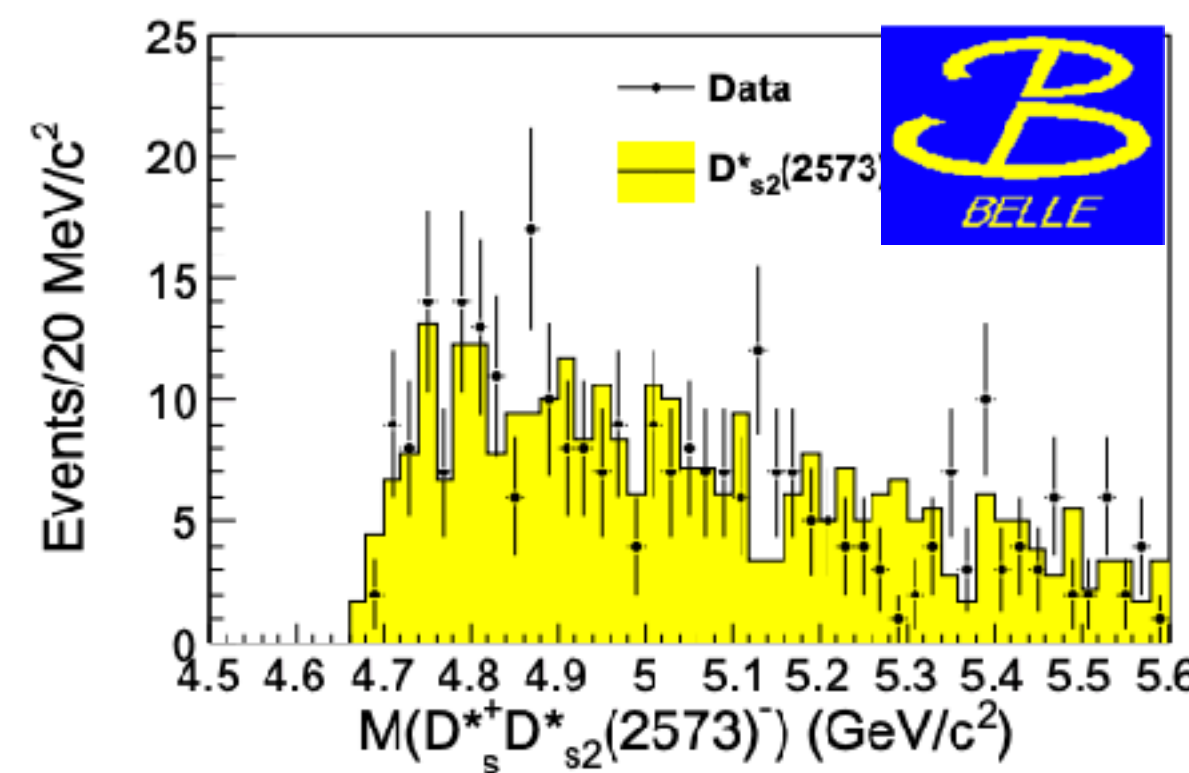
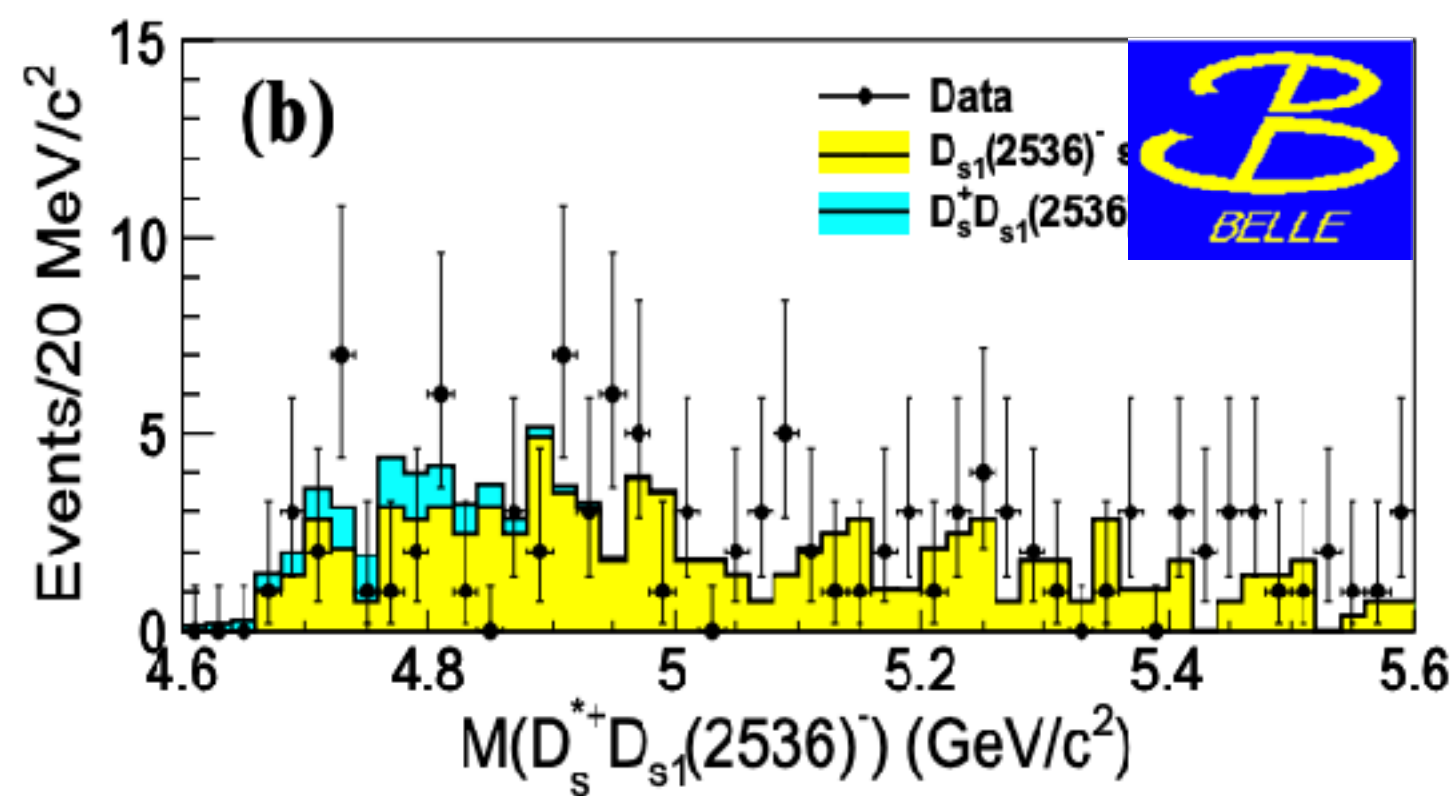
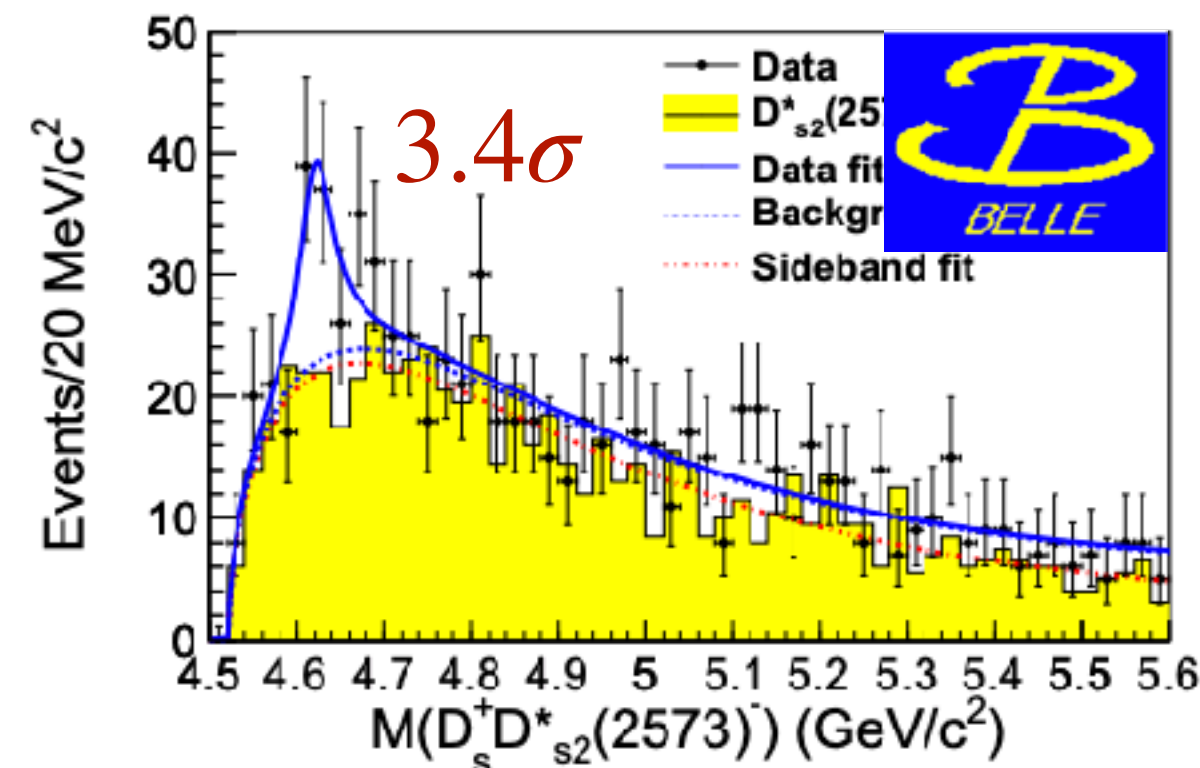
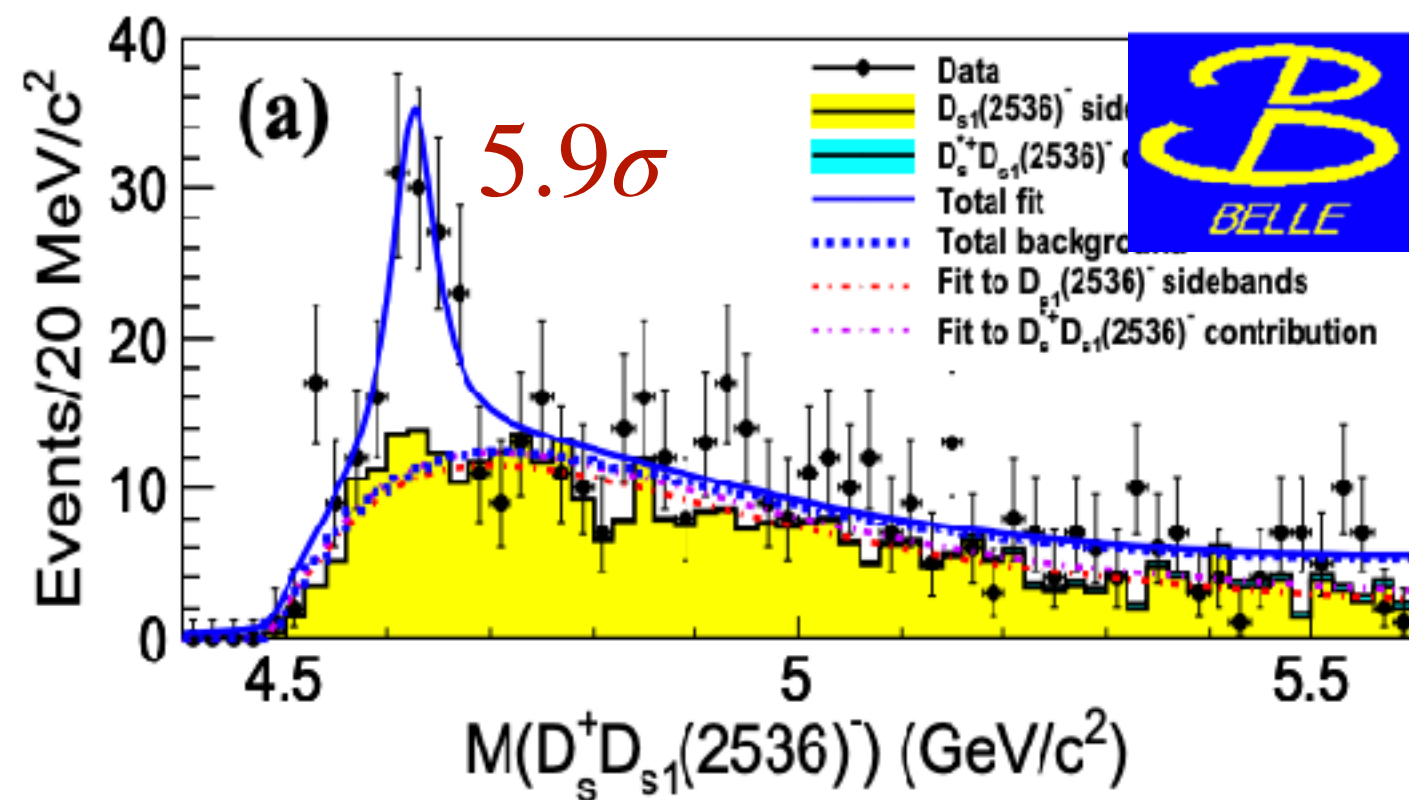
Belle II already achieve the world record instantaneous luminosity: $4.7 \times 10^{34} / \text{cm}^2 / \text{s}$

Initial state radiation

Allow us to reach lower c.m. energy “for free”



Recent: A new vector state in $e^+e^- \rightarrow D_s^+ D_{s1}(2536)^-, D_s^+ D_{s2}^*(2573)^-$



$Y(4620) = Y(4660)?$

Experiments	Mass (MeV/c ²)	Width (MeV)
Belle, $\Lambda_c^+ \Lambda_c^-$	4634_{-7-8}^{+8+5}	92_{-24-21}^{+40+10}
Belle, $\pi^+ \pi^- J/\psi$	$4652 \pm 10 \pm 8$	$68 \pm 11 \pm 1$
BaBar, $\pi^+ \pi^- J/\psi$	$4669 \pm 21 \pm 3$	$104 \pm 48 \pm 10$
Belle, $D_s^+ D_{s1}(2536)^-$	$4625.9_{-6.0}^{+6.2} \pm 0.4$	$49.8_{-11.5}^{+13.9} \pm 4.0$
Belle, $D_s^+ D_{s1}^*(2573)^-$	$4619.8_{-8.0}^{+8.9} \pm 2.3$	$47.0_{-14.8}^{+31.3} \pm 4.6$

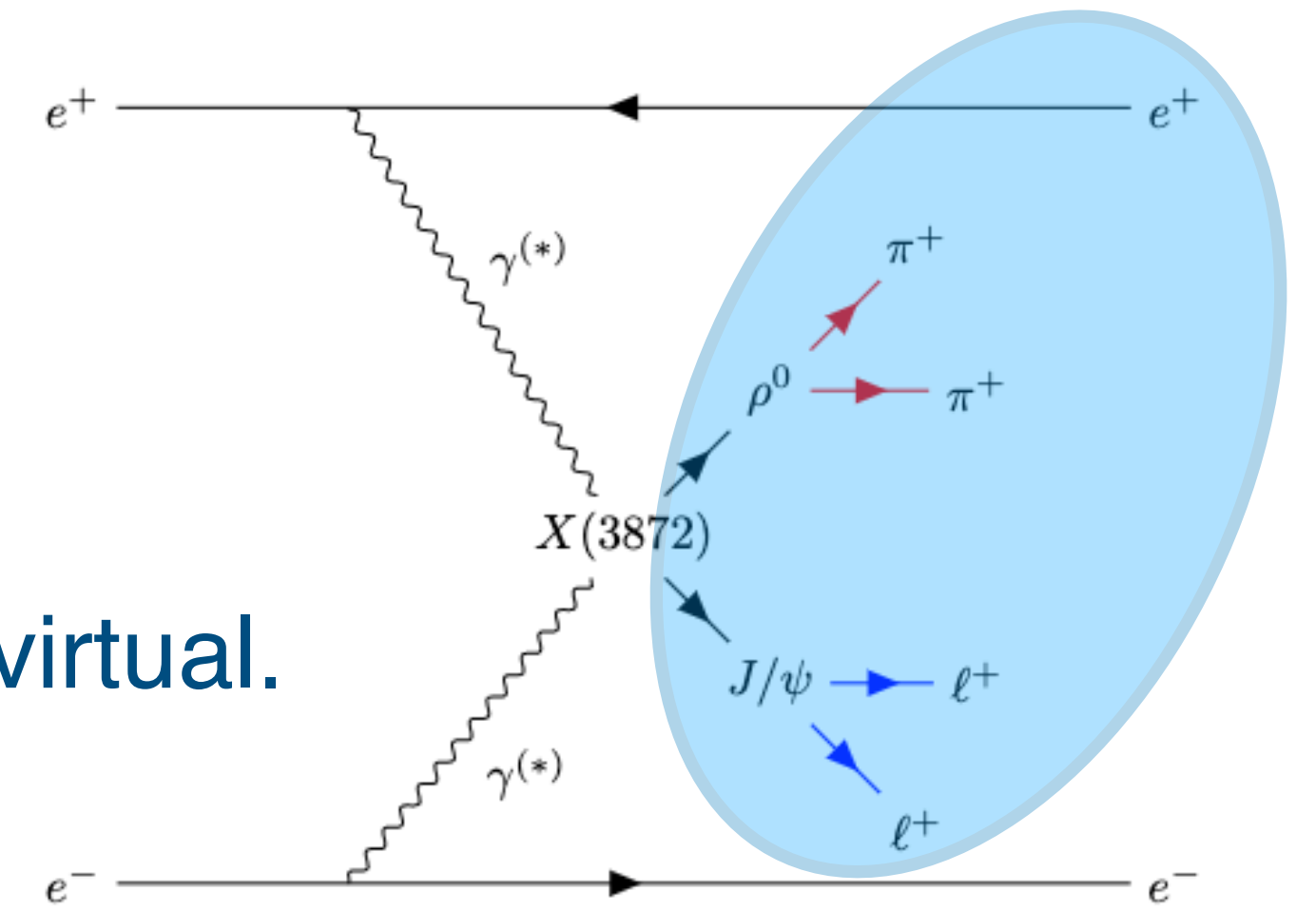
Phys.Rev.D 100 (2019) 11, 111103

Phys.Rev.D 101 (2020) 9, 091101

Evidence of $\gamma\gamma^* \rightarrow X(3872)$

Axial-vector particles are forbidden to decay to two real photons.

Mesons with $J^{PC} = 1^{++}$ could be produced if one or both photons are virtual.

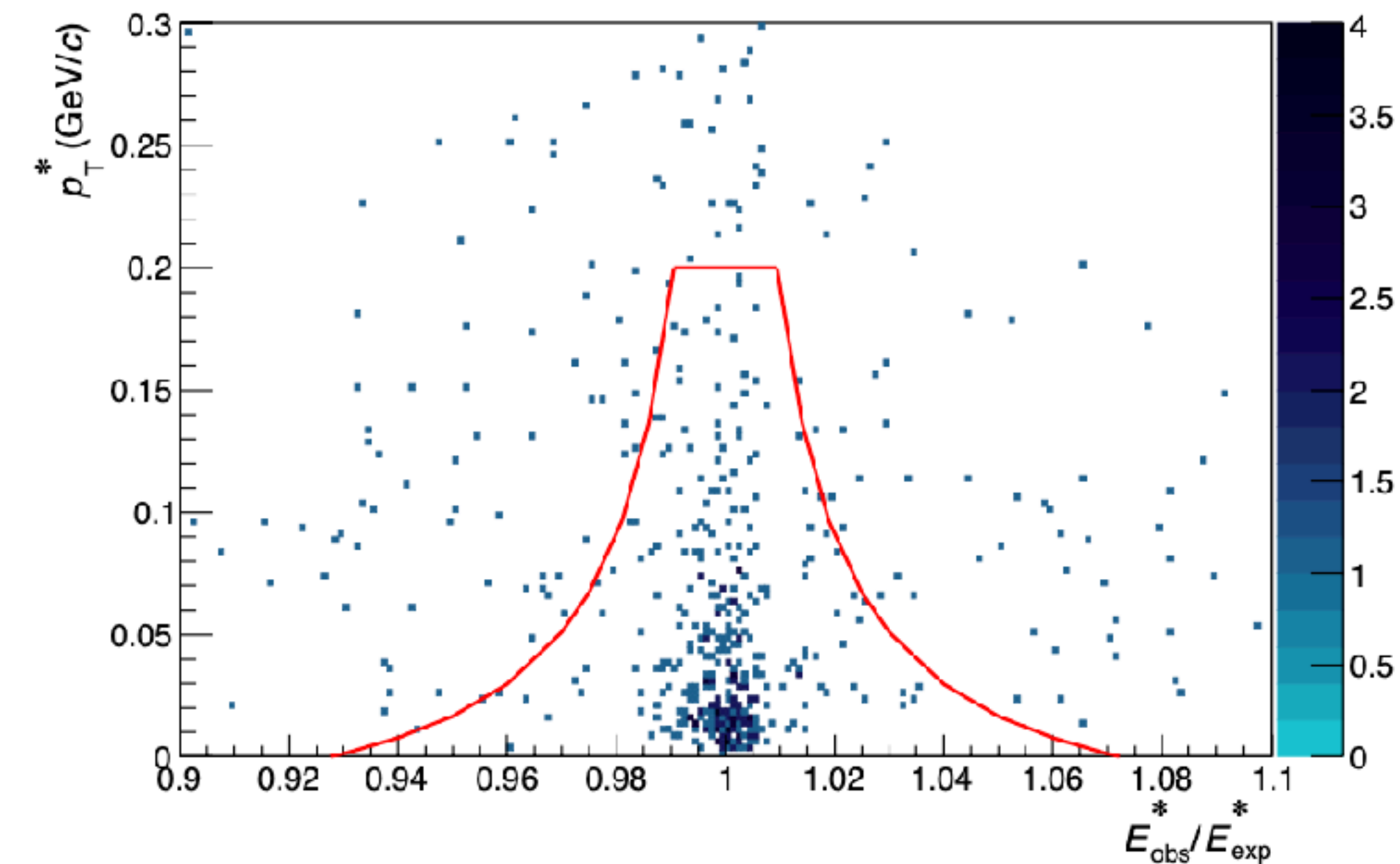


reconstruction:

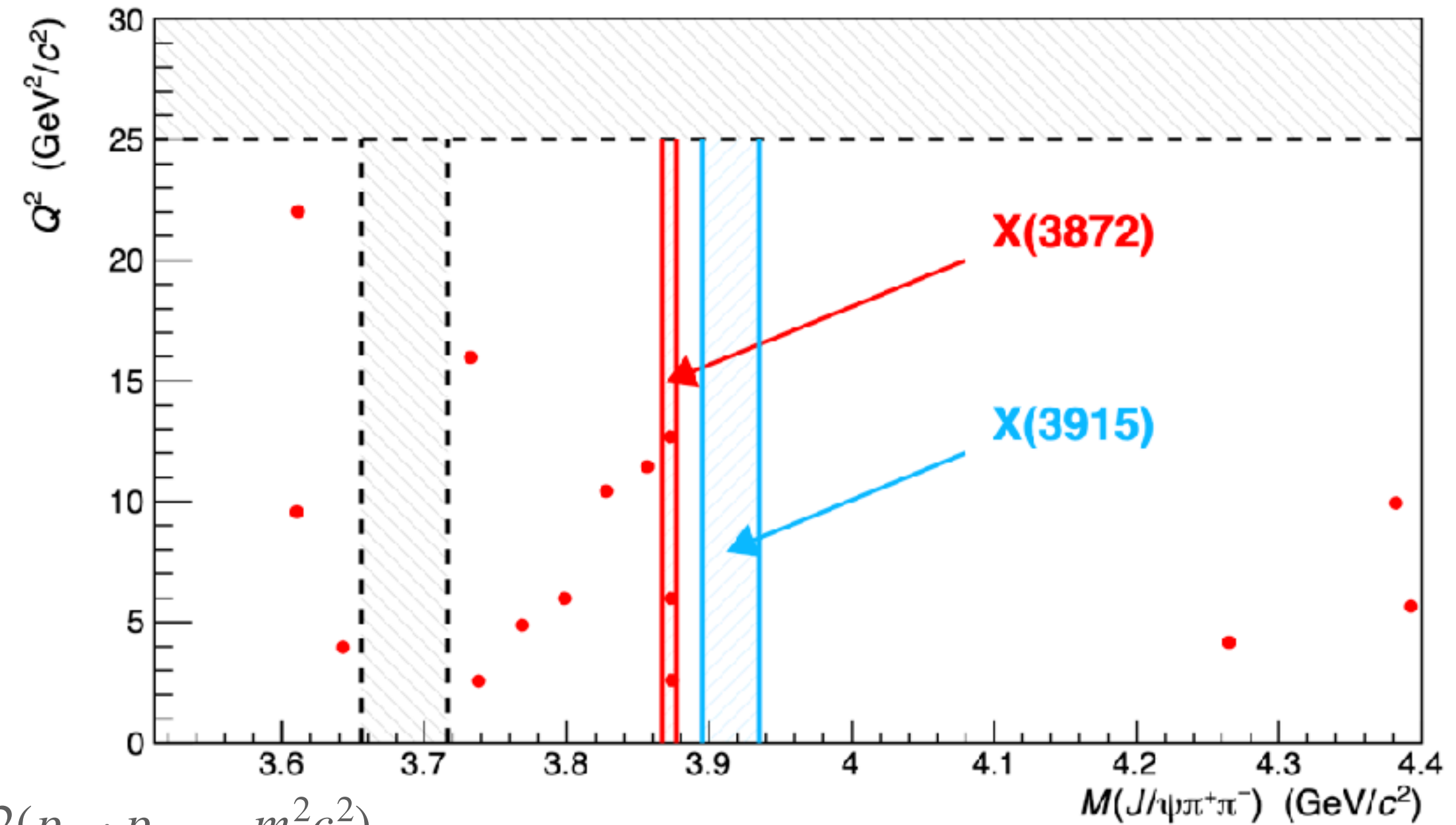
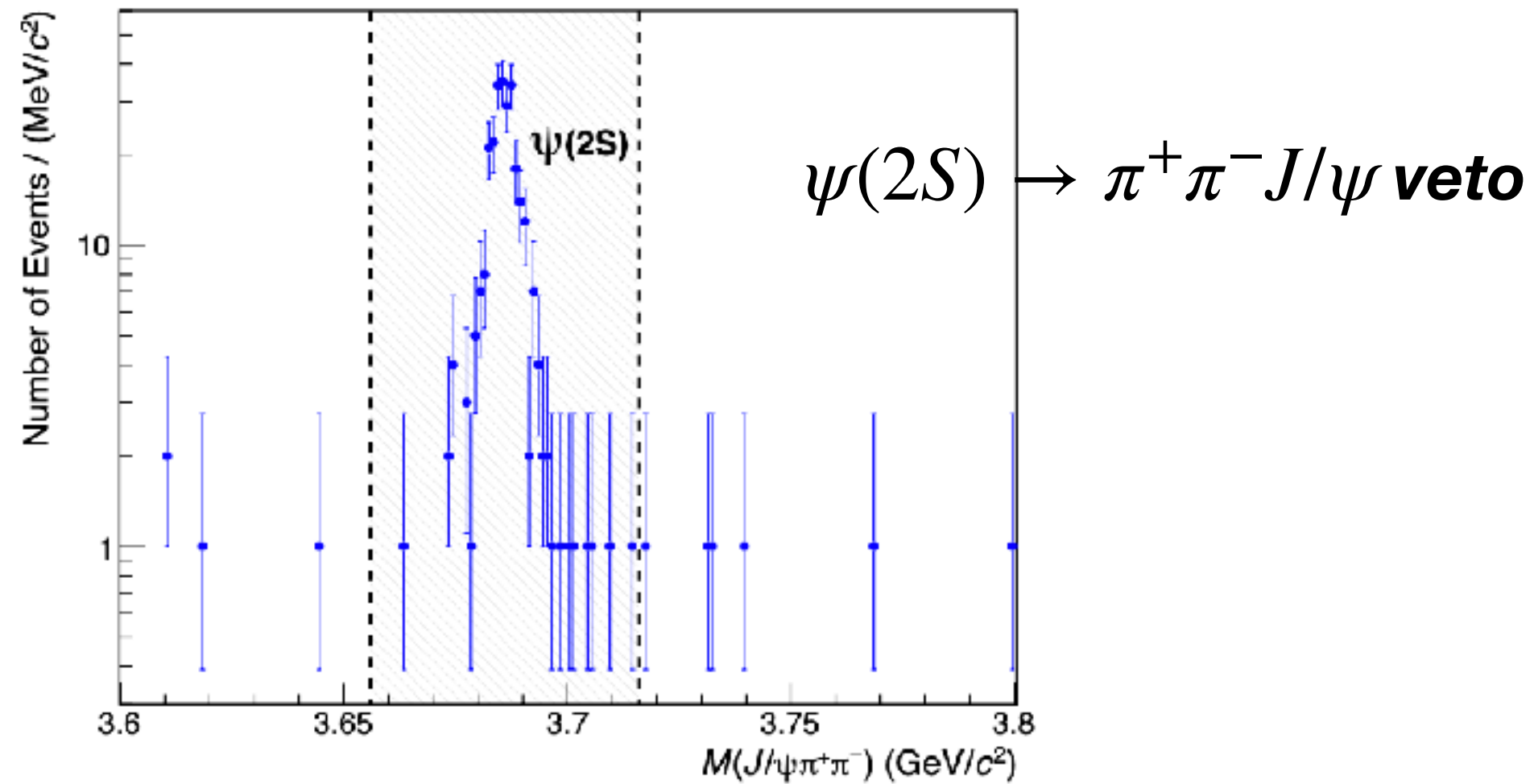
- ✧ $N(track) = 5$ with $\sum Q = \pm 1$
 - ✦ two pions, two leptons(e/μ)
 - ✦ one extra electron/positron (from beam)
 - ✦ No photon with $E > 0.4$ GeV or π^0
- ✧ $X(3872)$ & tagging electron: back to back
 - ✦ azimuthal angle difference within $(\pi \pm 0.1)$
- ✧ Visible transverse momentum < 0.2 GeV/c; measured $\pi^+\pi^-J/\psi$ energy E_{obs}^* consistent with the expectation E_{exp}^*
- ✧ Missing momentum of event projection:

$$p_{z,mis} < -0.4 \text{ GeV}/c^2 \text{ for } e^- \text{ tag}$$

$$p_{z,mis} > +0.4 \text{ GeV}/c^2 \text{ for } e^+ \text{ tag.}$$



Dominant background



$$Q^2 = 2(p_{in} \cdot p_{out} - m_e^2 c^2),$$

where $p_{in/out}$ is the momentum of the incoming (beam) electron and outgoing (tagging) electron

expected number of background:

- $(3 - 5) \times 10^{-2} / (10 \text{ MeV}/c^2)$ from internal bremsstrahlung.
- 0.11 ± 0.10 (0.3 for X(3915)) extrapolated from fit to the background events

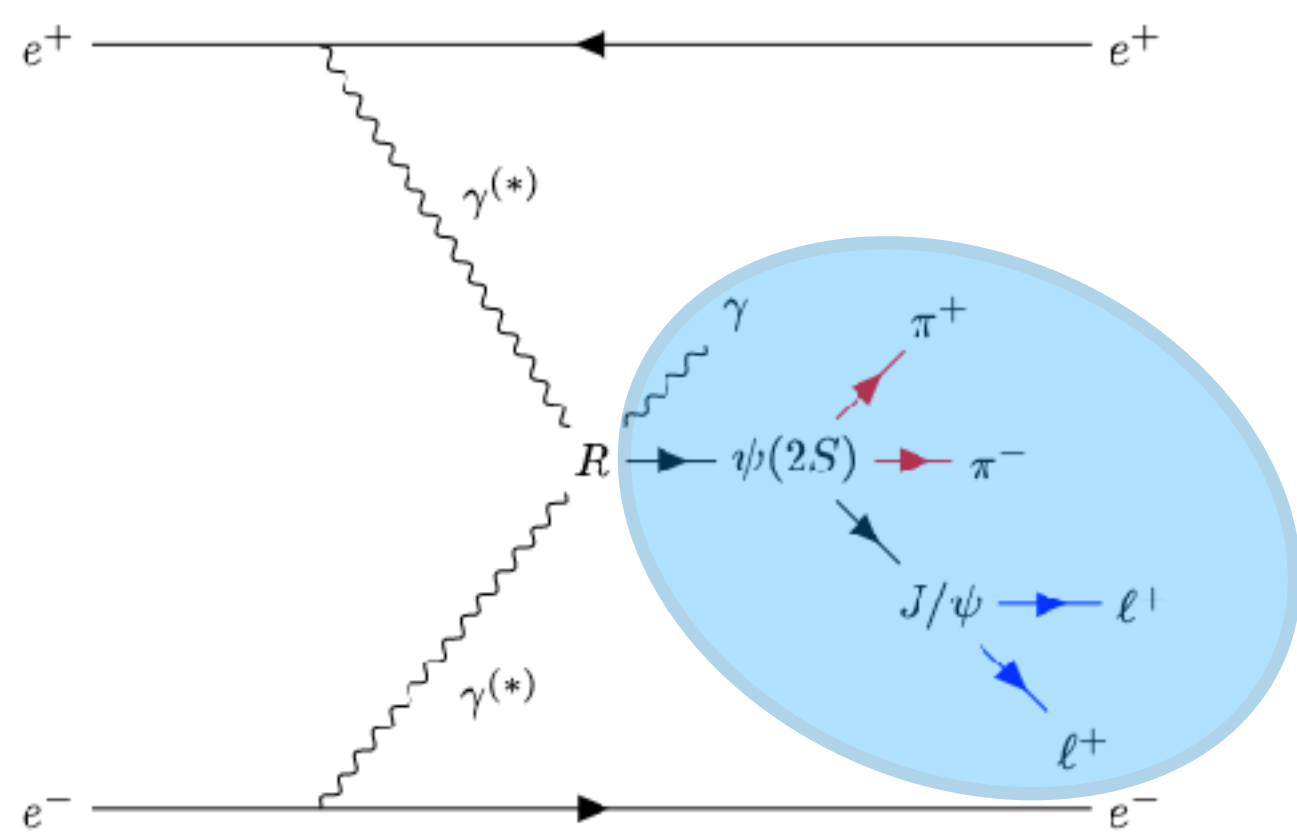
Three events are found in the signal region, with a significance of 3.2σ considering the background.

$$N_{\text{sig}} = 2.9_{-2.0}^{+2.2}(\text{stat.}) \pm 0.1(\text{syst.}) \text{ for } X(3872), N_{\text{sig}} < 2.14 \text{ for } X(3915)$$

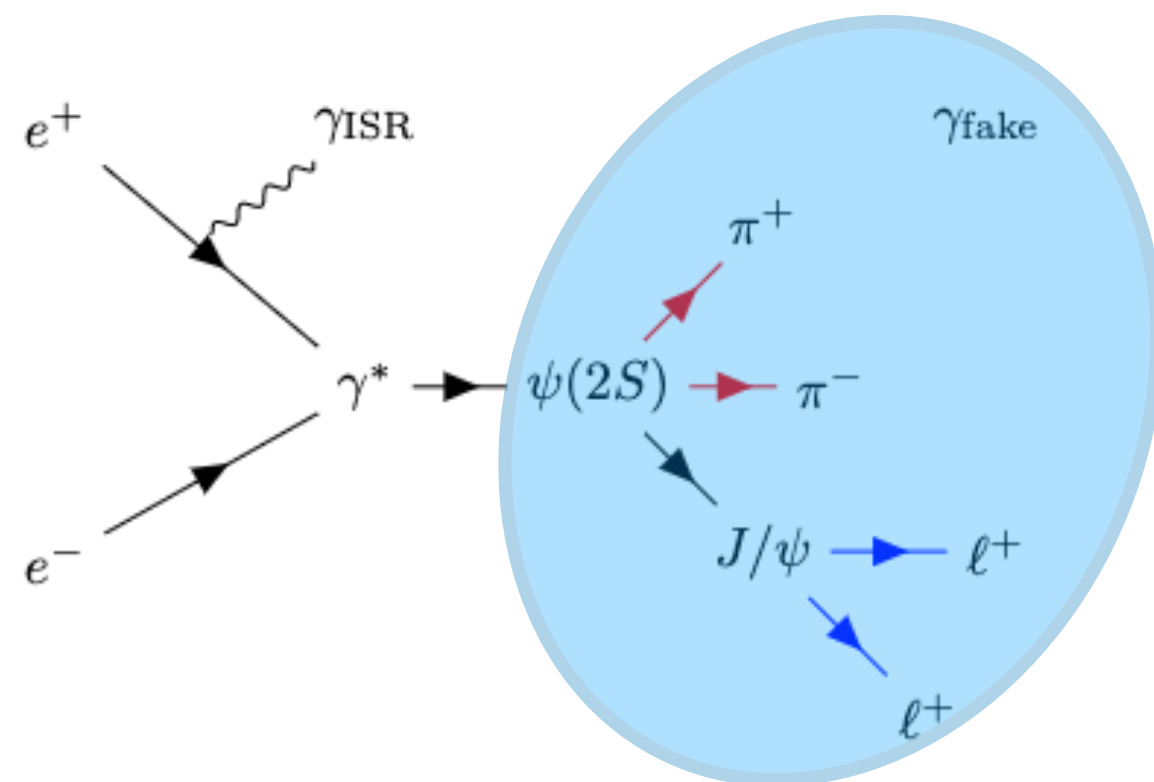
Exotic candidates in $\gamma\gamma \rightarrow \gamma\psi(2S)$

Both 0^{++} and 2^{++} could be produced in the two photon collisions, and can radiatively decay to $\psi(2S)$

Signal process



Dominate background: $e^+e^- \rightarrow \gamma_{ISR}\psi(2S)$

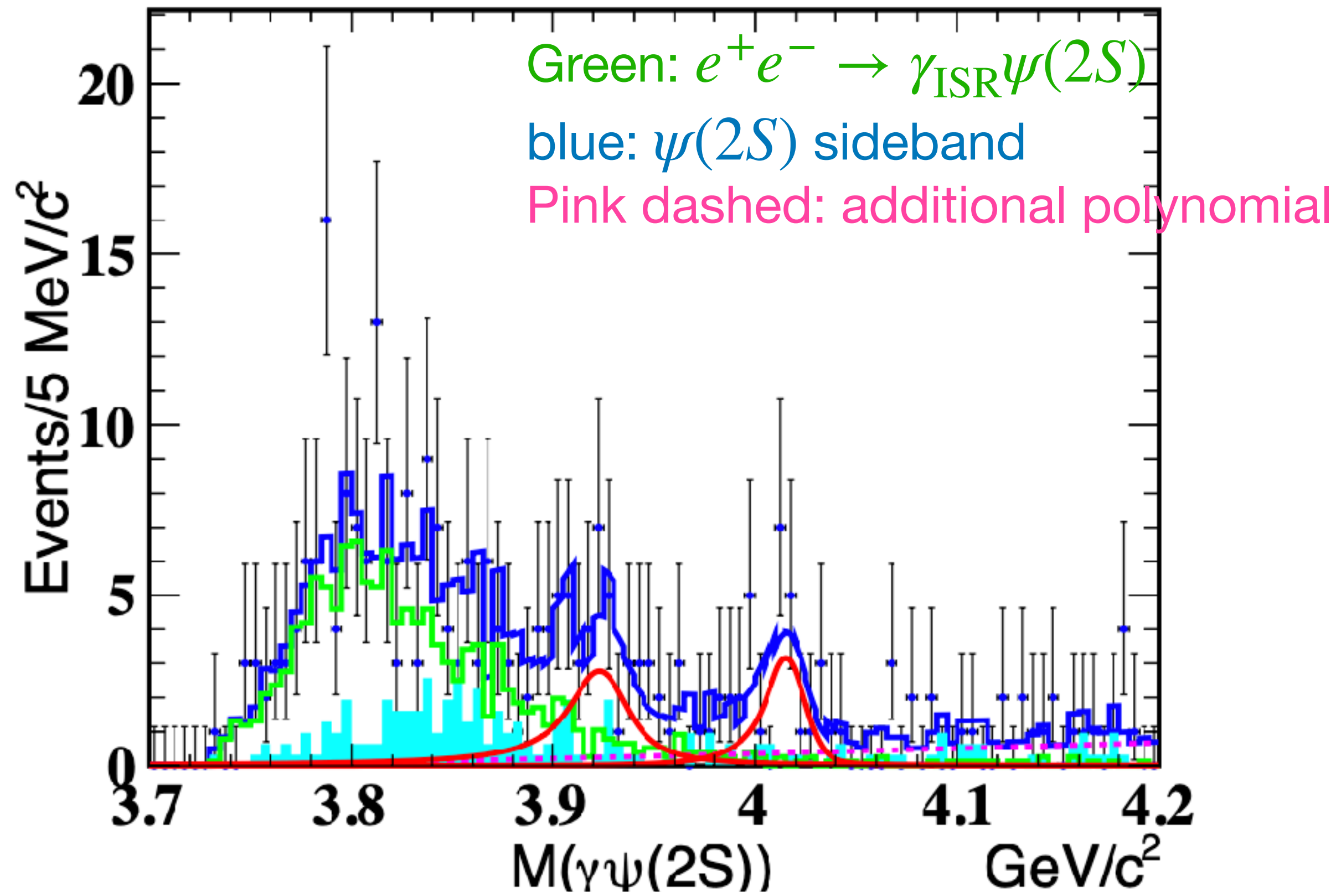


reconstruction:

- * Reconstructing $\pi^+\pi^-\ell^+\ell^-$
 - ◆ J/ψ reconstructed with two leptons (e/μ)
 - ◆ $\psi(2S)$ reconstructed with $\pi^+\pi^-J/\psi$

background suppression:

- * Recoiling mass of $\gamma\psi(2S)$
 - ◆ $M_{\text{rec}}^2(\gamma\psi(2S)) > 10 (\text{GeV}/c^2)^2$
- * Transverse momentum balances
 - ◆ $P_t^*(\psi(2S)) > 0.1 \text{ GeV}/c$
 - ◆ $P_t^*(\gamma\psi(2S)) < 0.2 \text{ GeV}/c$



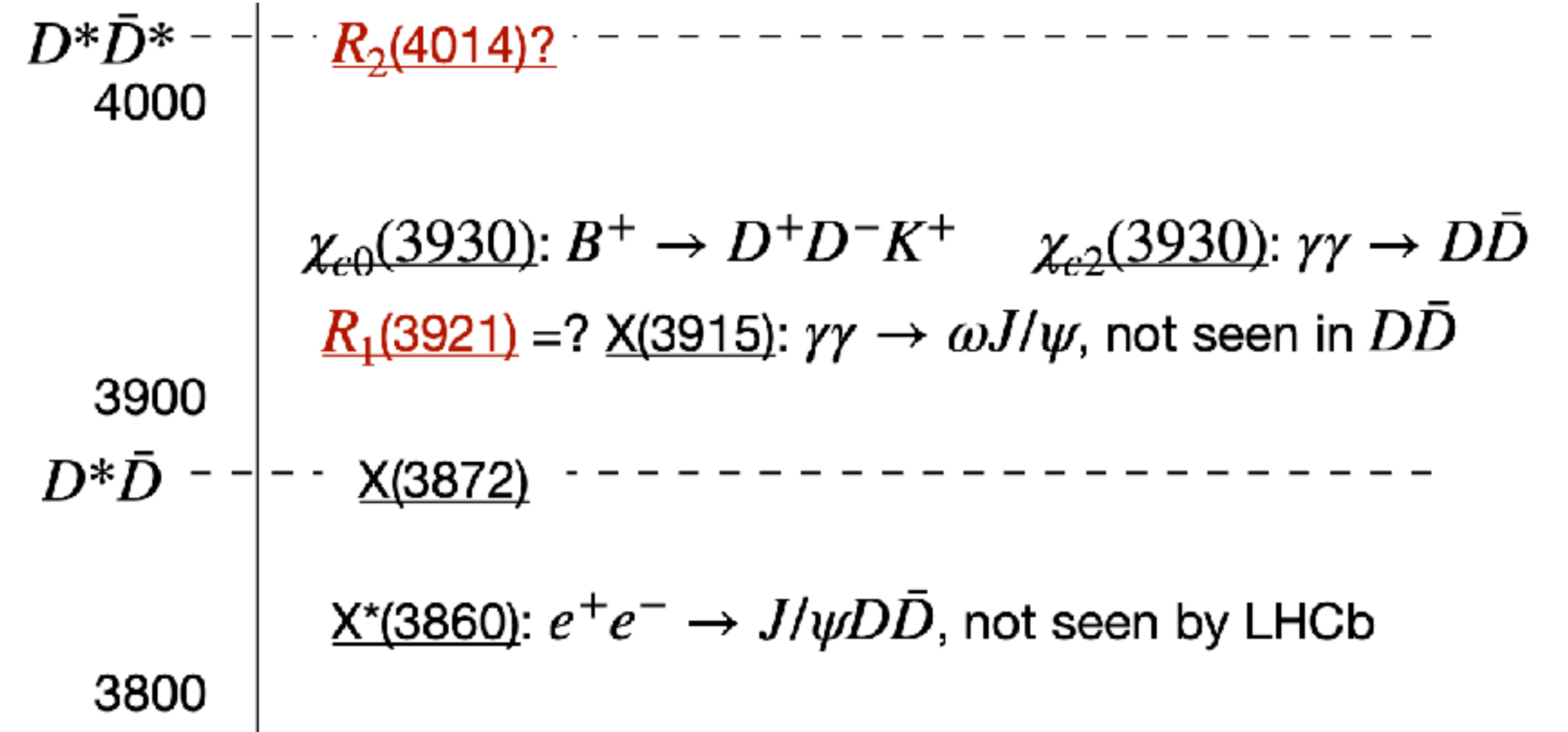
Fitting PDF:

$$f_{\text{sum}} = f_{R_1} + f_{R_2} + f_{\text{ISR}} + f_{\text{bkg}} + f_{\text{SB}}$$

$$f_R \propto \varepsilon \cdot (\text{BW} \otimes \text{CB}).$$

possible interference is ignored

possible nature?



Resonant parameters	$J = 0$	$J = 2$
M_{R_1}	$3922.4 \pm 6.5 \pm 2.0$	
Γ_{R_1}	$22 \pm 17 \pm 4$	
$\Gamma_{\gamma\gamma} \mathcal{B}(R_1 \rightarrow \gamma\psi(2S))$	$9.8 \pm 3.6 \pm 1.2$	$2.0 \pm 0.7 \pm 0.2$
M_{R_2}	$4014.3 \pm 4.0 \pm 1.5$	
Γ_{R_2}	$4 \pm 11 \pm 6$	
$\Gamma_{\gamma\gamma} \mathcal{B}(R_2 \rightarrow \gamma\psi(2S))$	$6.2 \pm 2.2 \pm 0.8$	$1.2 \pm 0.4 \pm 0.2$

Significance of R_1 is 3.1σ considering systematic uncertainty.

Significance of R_2 is 2.8σ after considering LEE.

Excess states = exotics?

Observation of $e^+e^- \rightarrow \Upsilon(1,2S)\eta$ at 10.866 GeV

For bottomonia below $B\bar{B}$ threshold, predictions of hadronic transition rates are consistent with measurements.

Measured hadronic transition rates between bottomonia above open bottom threshold are higher than predictions.

e.g. $\frac{\Gamma(\Upsilon(5S) \rightarrow h_b(1P)\pi^+\pi^-)}{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(1S)\pi^+\pi^-)} = 0.46 \pm 0.08^{+0.07}_{-0.12}$ $\frac{\Gamma(\Upsilon(5S) \rightarrow h_b(2P)\pi^+\pi^-)}{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(2S)\pi^+\pi^-)} = 0.77 \pm 0.08^{+0.22}_{-0.17}$ Prediction: $\mathcal{O}(10^{-2})$

Analysis of similar processes is crucial for better understanding of the quark structure of bottomonium states above $B\bar{B}$ threshold.

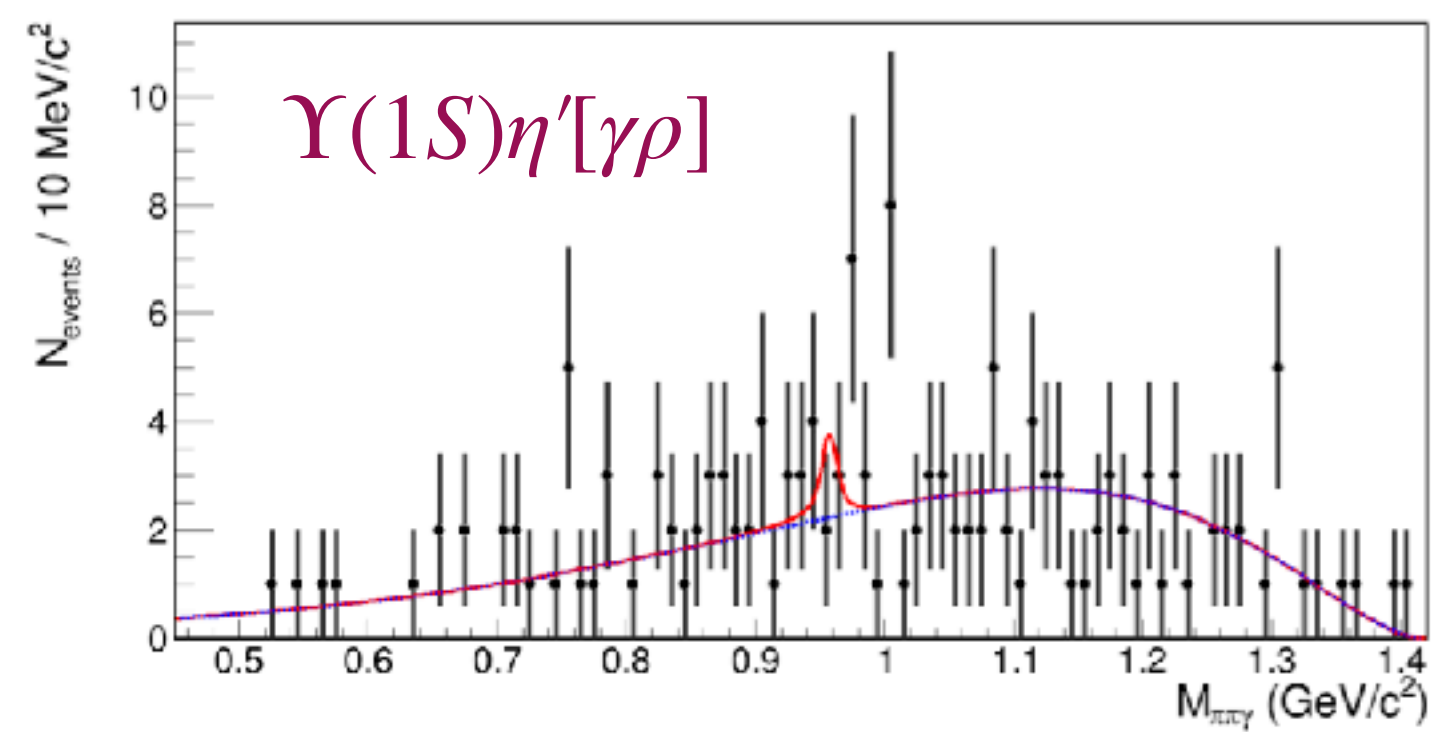
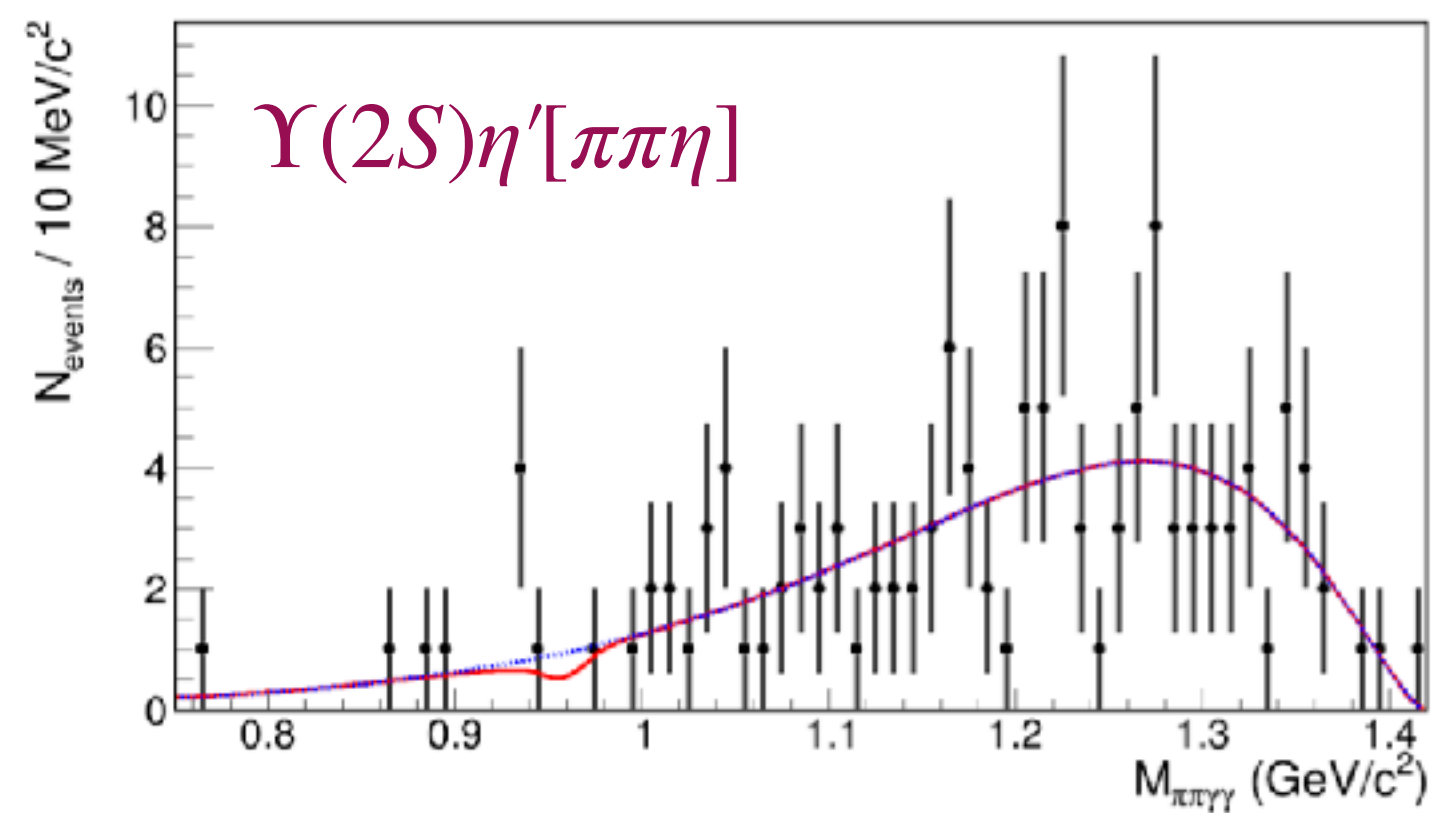
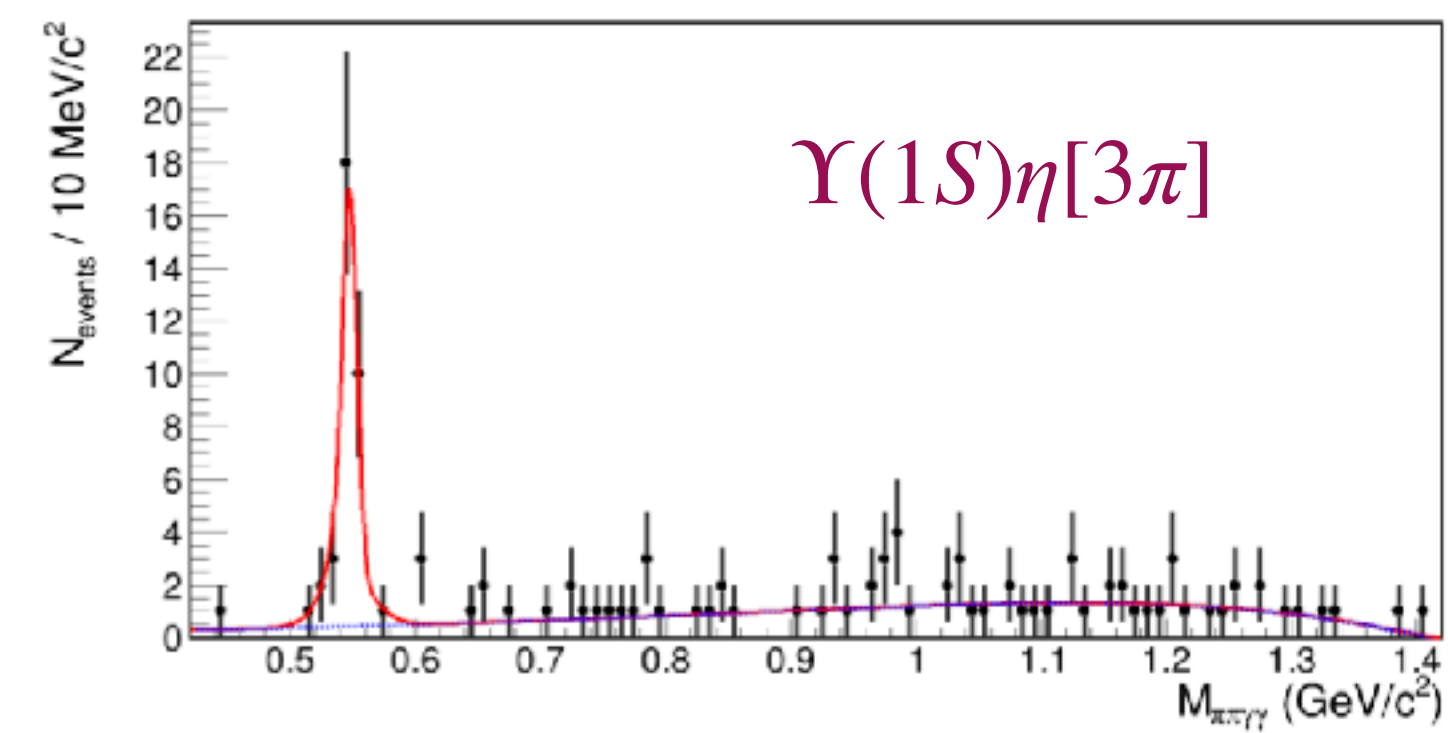
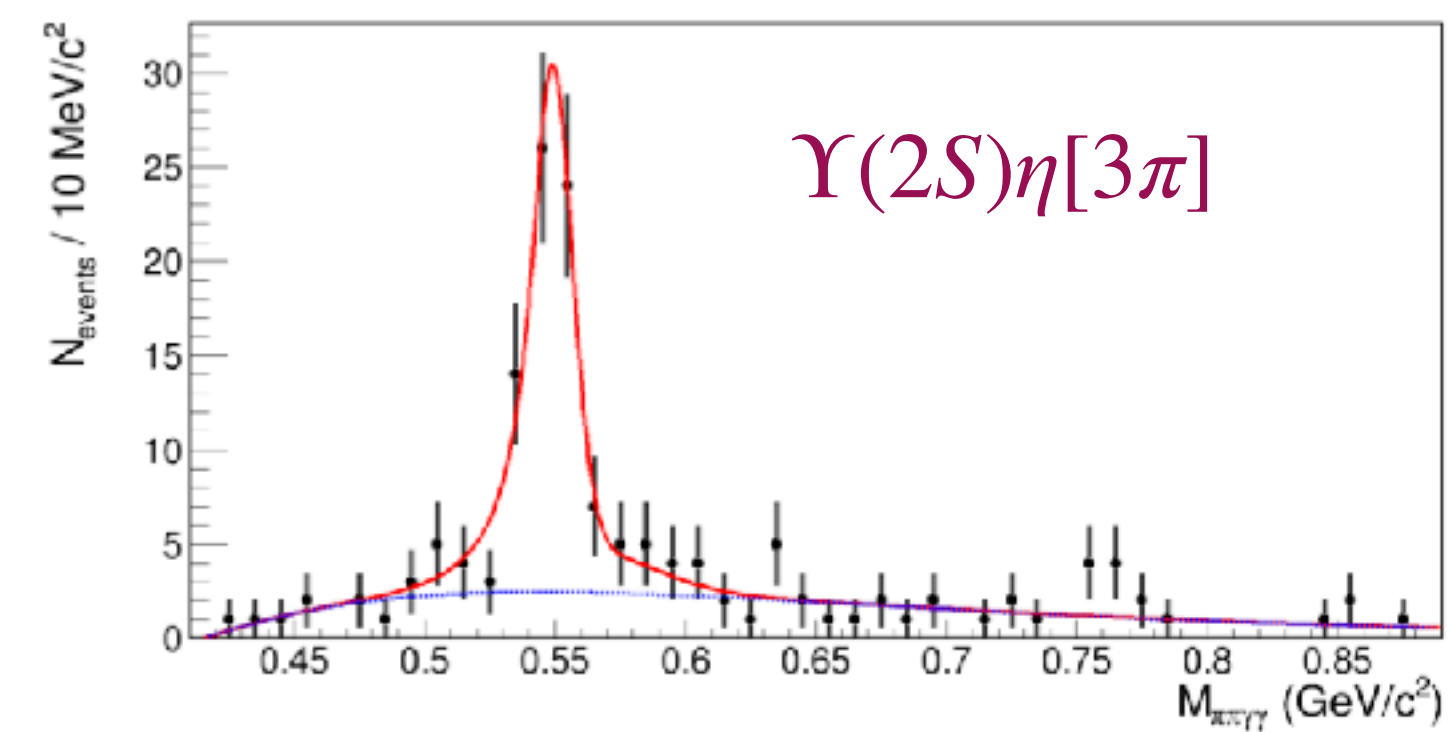
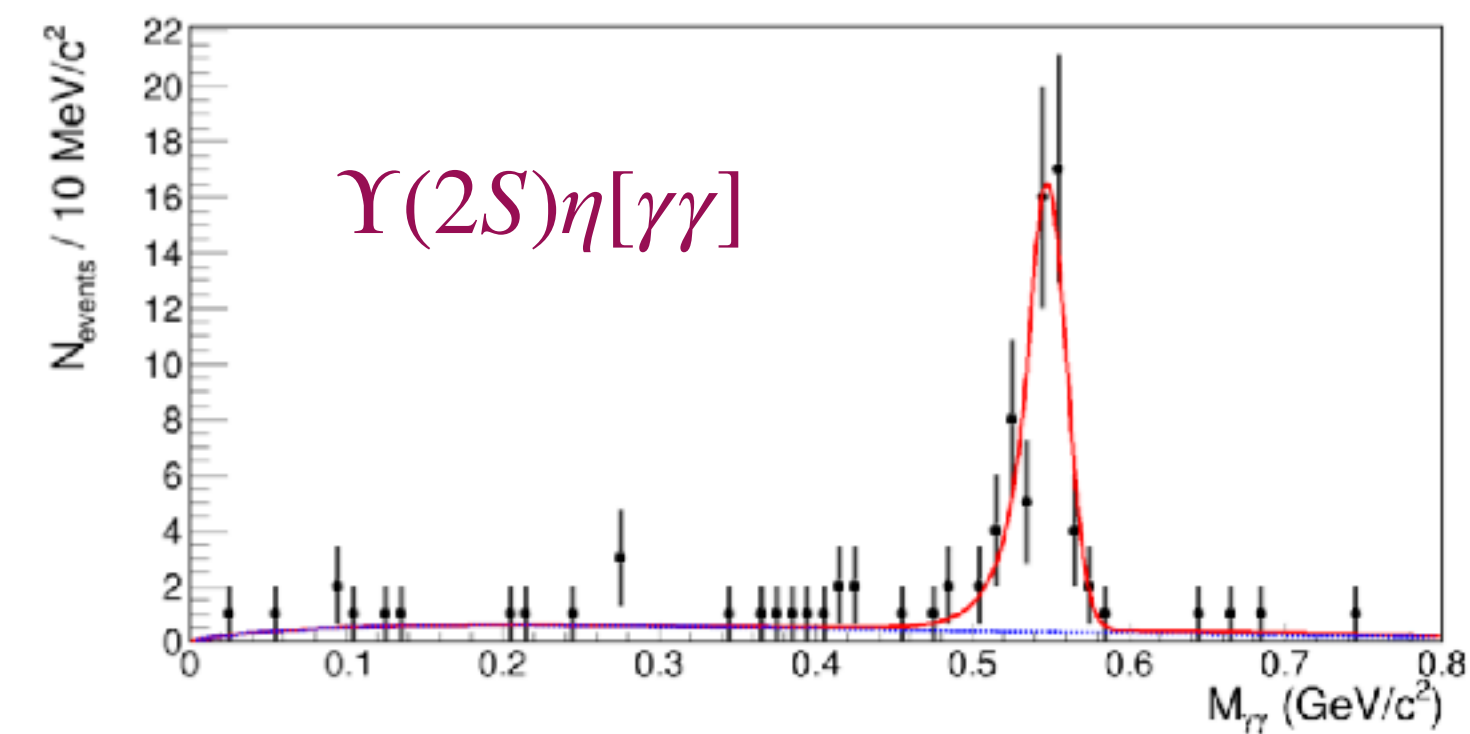
reconstruction:

- * Final states $\pi^+\pi^-\mu^+\mu^-\gamma(\gamma)$: $\Upsilon(2S)\eta[3\pi]$, $\Upsilon(2S)\eta[\gamma\gamma]$, $\Upsilon(1S)\eta[3\pi]$, $\Upsilon(1S)\eta'[\pi\pi\eta]$, $\Upsilon(1S)\eta'[\gamma\rho]$
- * For $\eta \rightarrow \gamma\gamma$
 - ◆ $\Upsilon(2S)$ reconstructed with $\pi^+\pi^-J/\psi$
- * For $\eta \rightarrow \pi^+\pi^-\pi^0$
 - ◆ $\Upsilon(1,2S)$ reconstructed with two leptons

Data sample:

118 fb⁻¹ at $\Upsilon(5S)$

21 fb⁻¹ energy scan in 10.63 ~ 11.02 GeV



Significant $\Upsilon(1S)\eta$ and $\Upsilon(2S)\eta$ signals

- 10.2σ for $e^+e^- \rightarrow \Upsilon(1S)\eta$
- 16.5σ for $e^+e^- \rightarrow \Upsilon(2S)\eta$

Assuming process only from $\Upsilon(5S)$:

$$\mathcal{B}(\Upsilon(5S) \rightarrow \Upsilon(1S)\eta) = (0.85 \pm 0.15 \pm 0.08) \times 10^{-3},$$

$$\mathcal{B}(\Upsilon(5S) \rightarrow \Upsilon(2S)\eta) = (4.13 \pm 0.41 \pm 0.37) \times 10^{-3},$$

$$\mathcal{B}(\Upsilon(5S) \rightarrow \Upsilon(1S)\eta') < 6.9 \times 10^{-5}, \text{ } CL = 90\%.$$

Corresponding fractions are:

- $\frac{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(2S)\eta)}{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(2S)\pi^+\pi^-)} = 0.52 \pm 0.06$ (stat.)
- $\frac{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(1S)\eta)}{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(1S)\pi^+\pi^-)} = 0.18 \pm 0.03$ (stat.)
- $\frac{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(1S)\eta')}{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(1S)\eta)} < 0.14$

Expected

~ 0.03

~ 0.005

~ 12

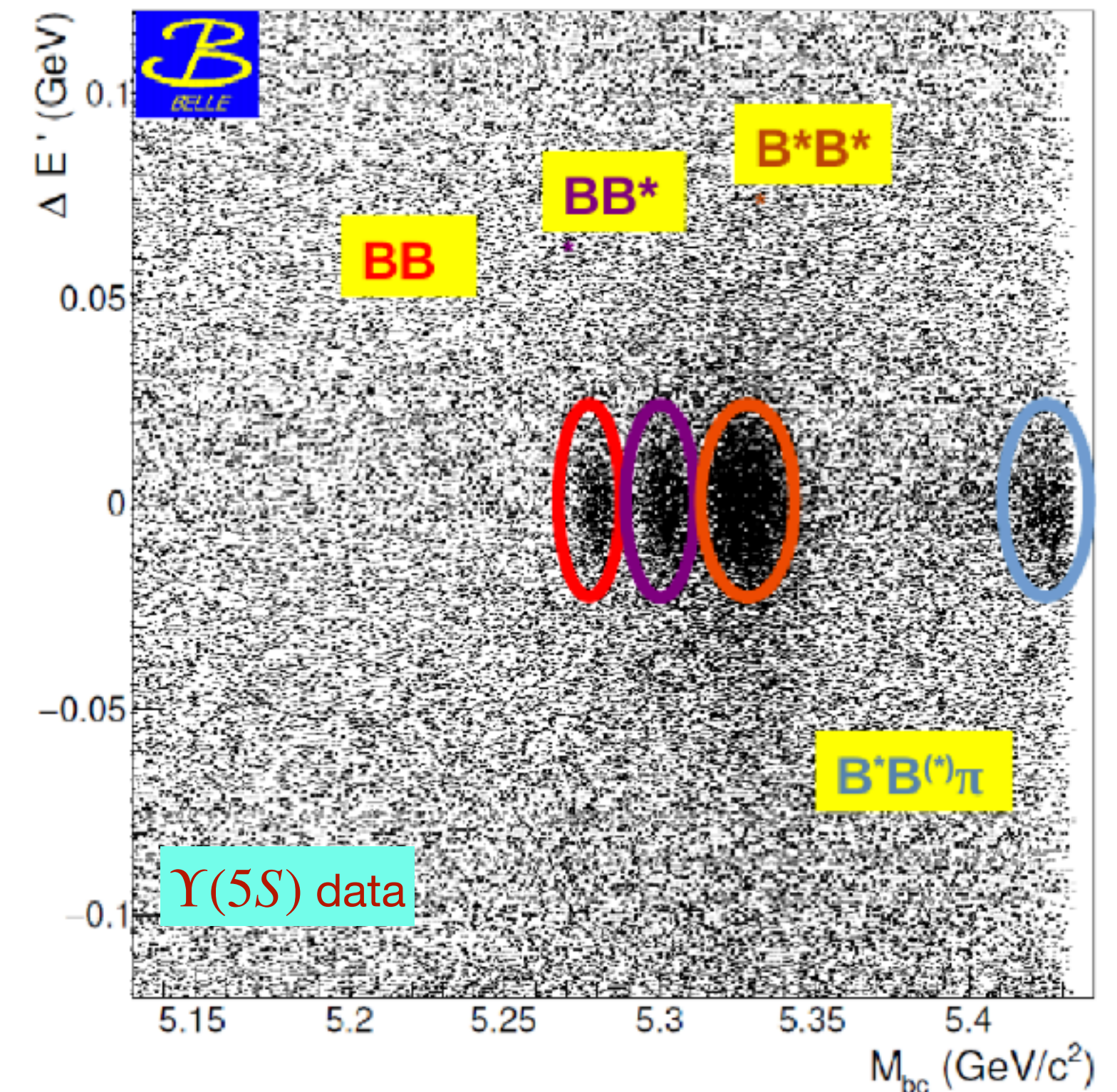
Cross section of $e^+e^- \rightarrow B^{(*)}\bar{B}^{(*)}$

Data samples:

- 571 fb⁻¹ on the $\Upsilon(4S)$ resonance
- 121 fb⁻¹ on the $\Upsilon(5S)$ resonance
- 16 fb⁻¹ distributed evenly in 16 points within 10.63 ~ 11.02 GeV

Full reconstruction of one B meson using FEI, a tool developed for tagging B meson in the $\Upsilon(4S) \rightarrow B\bar{B}$ decays using multivariate analysis for event selection.

[Comput. Softw. Big Sci. 3, 6 (2019)]

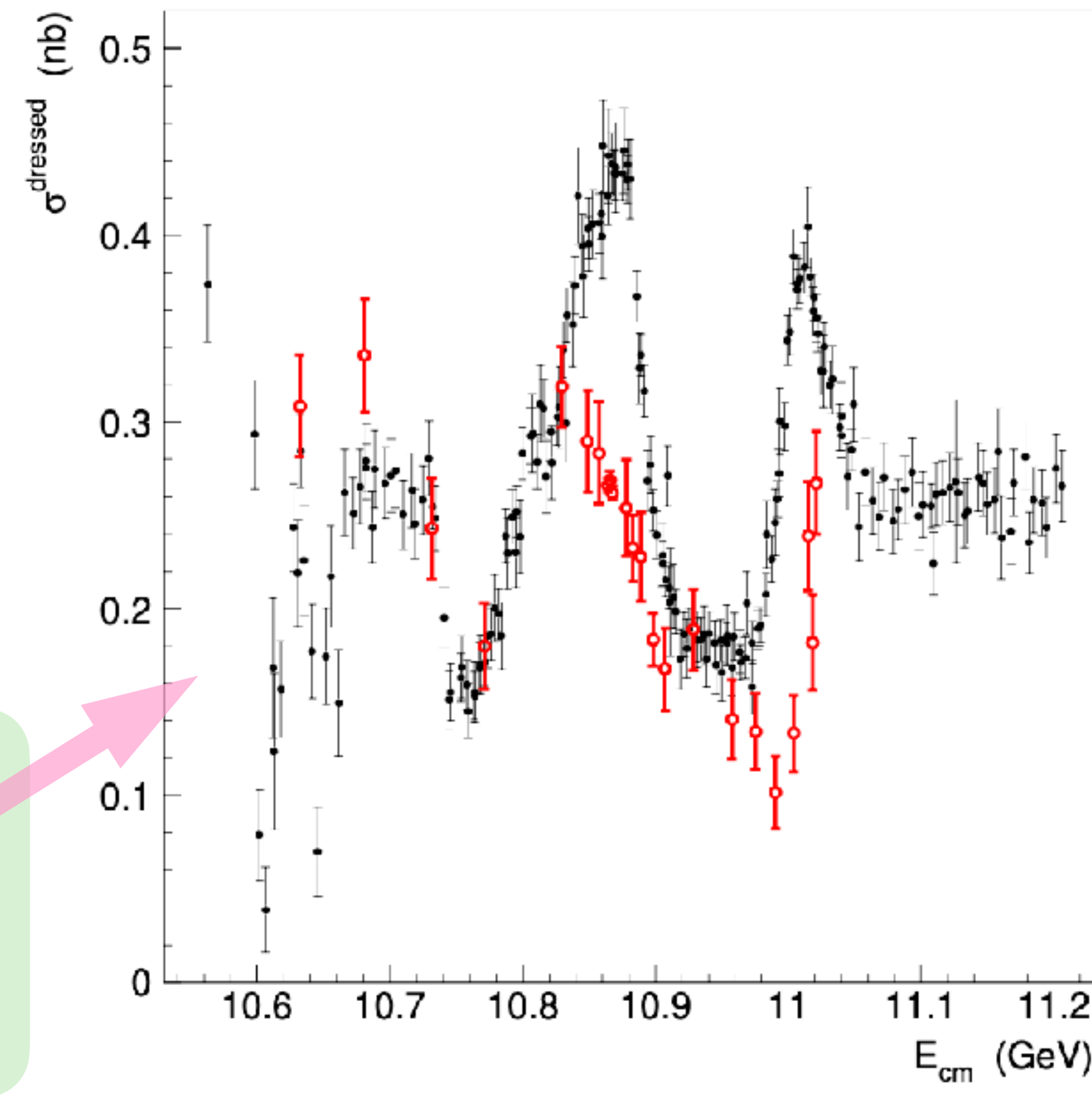
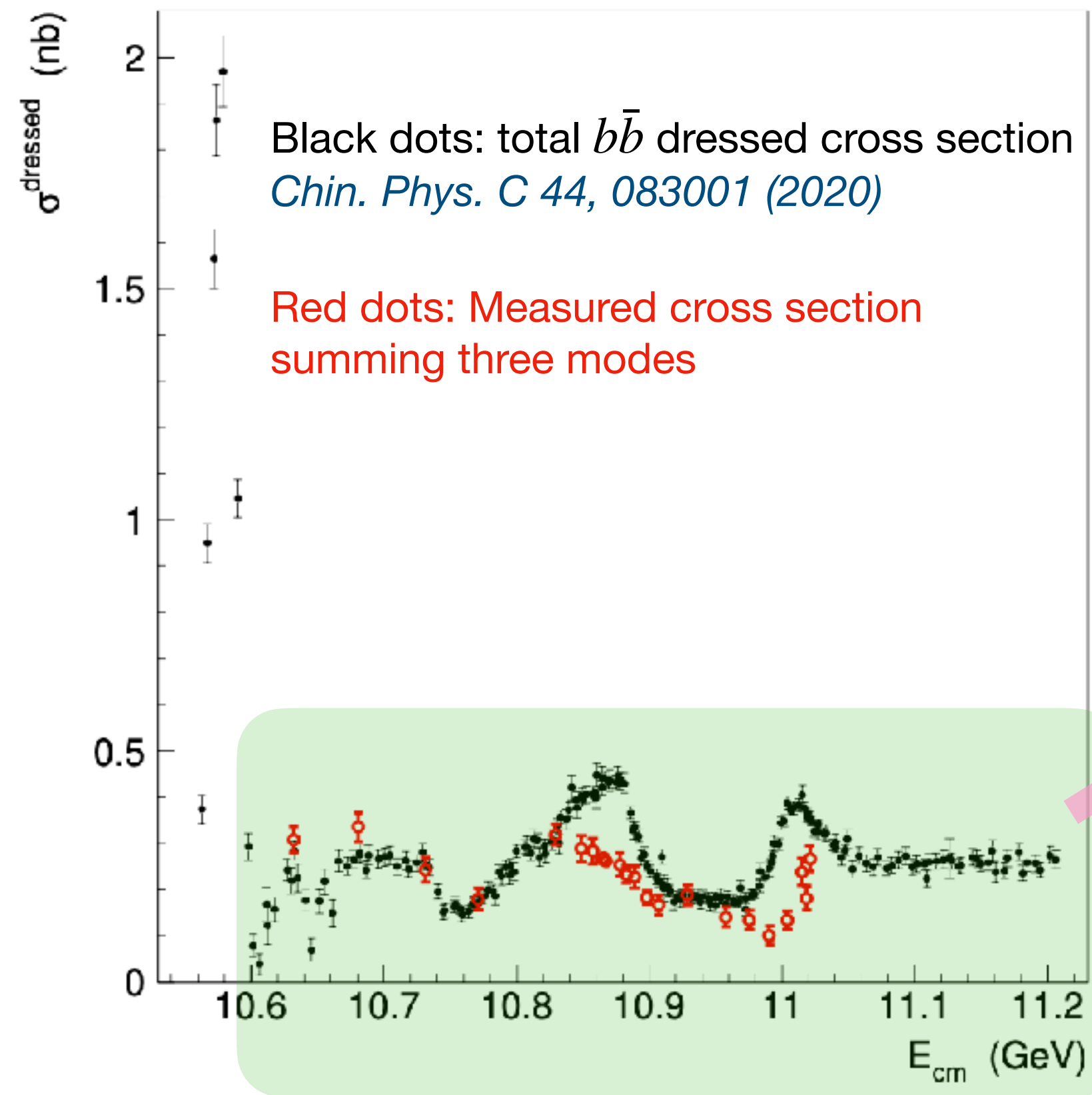
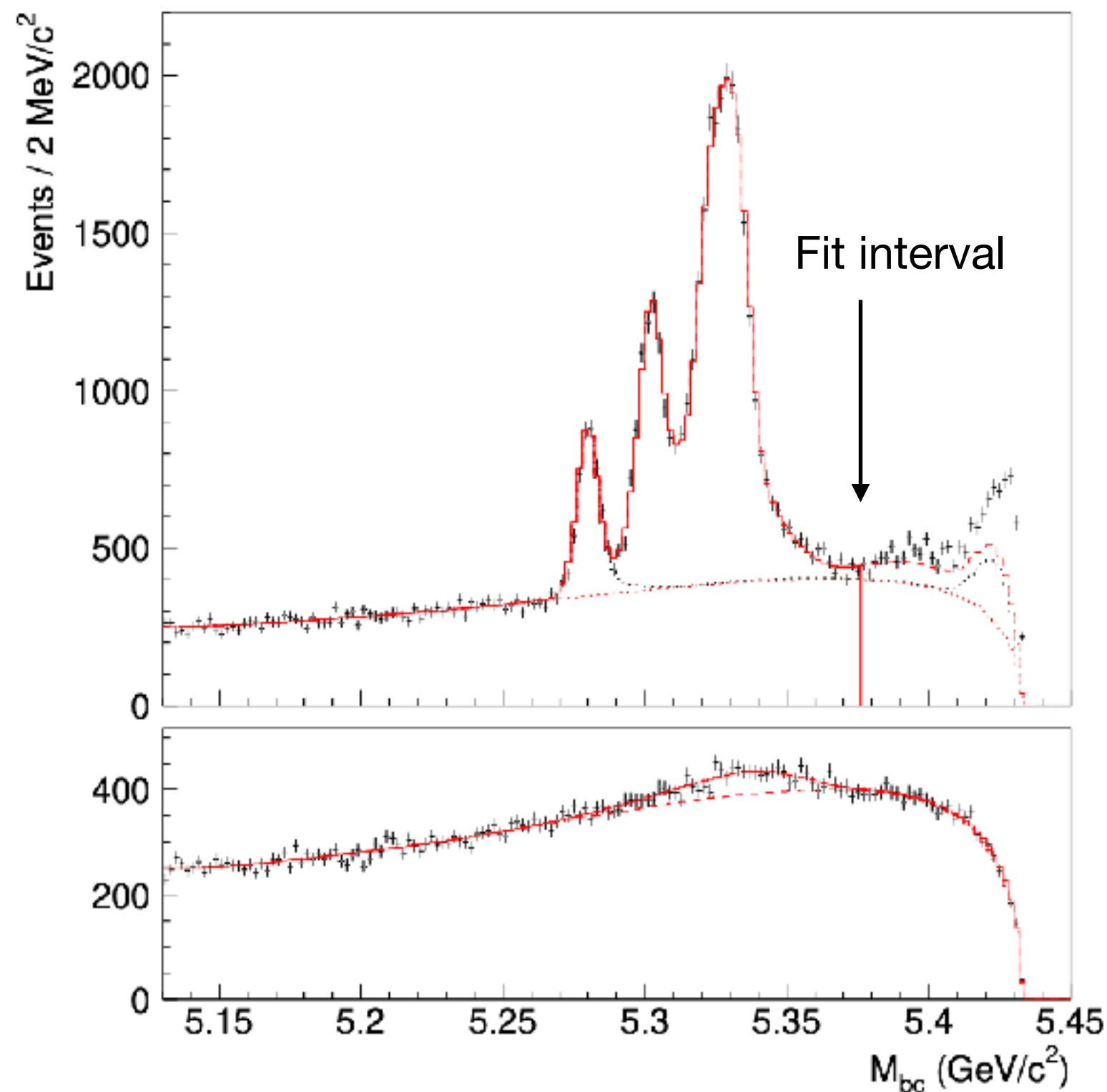


$$M_{bc} \equiv \sqrt{(E_{beam,CM})^2 - (p_{B,CM})^2}$$

$$\Delta E' \equiv \Delta E - M_{bc} + M_B$$

where

$$\Delta E \equiv E_{B,CM} - E_{beam,CM}$$



$$\sigma^{\text{dressed}} = \frac{N}{(1 + \delta_{\text{ISR}}) L \epsilon'}$$

- Fit M_{bc} distributions, selected in $\Delta E'$ signal region and $\Delta E'$ side-bands, for data taken at $\Upsilon(5S)$
- A similar fit is done for all 16 data points in the scan, to study the energy dependence of the cross sections

- ◆ No clear $\Upsilon(5S)$ signals from cross section measurement of $B^{(*)}\bar{B}^{(*)}$
- ◆ Excess is $B_s^{(*)}\bar{B}_s^{(*)}$, $B^{(*)}\bar{B}^{(*)}n\pi$, bottomonia + light hadrons, which contradicts expectation of $\Upsilon(5S) \rightarrow B^{(*)}\bar{B}^{(*)}$ dominantly.
- ◆ Cross sections do not peak at $\Upsilon(5S)$ mass.

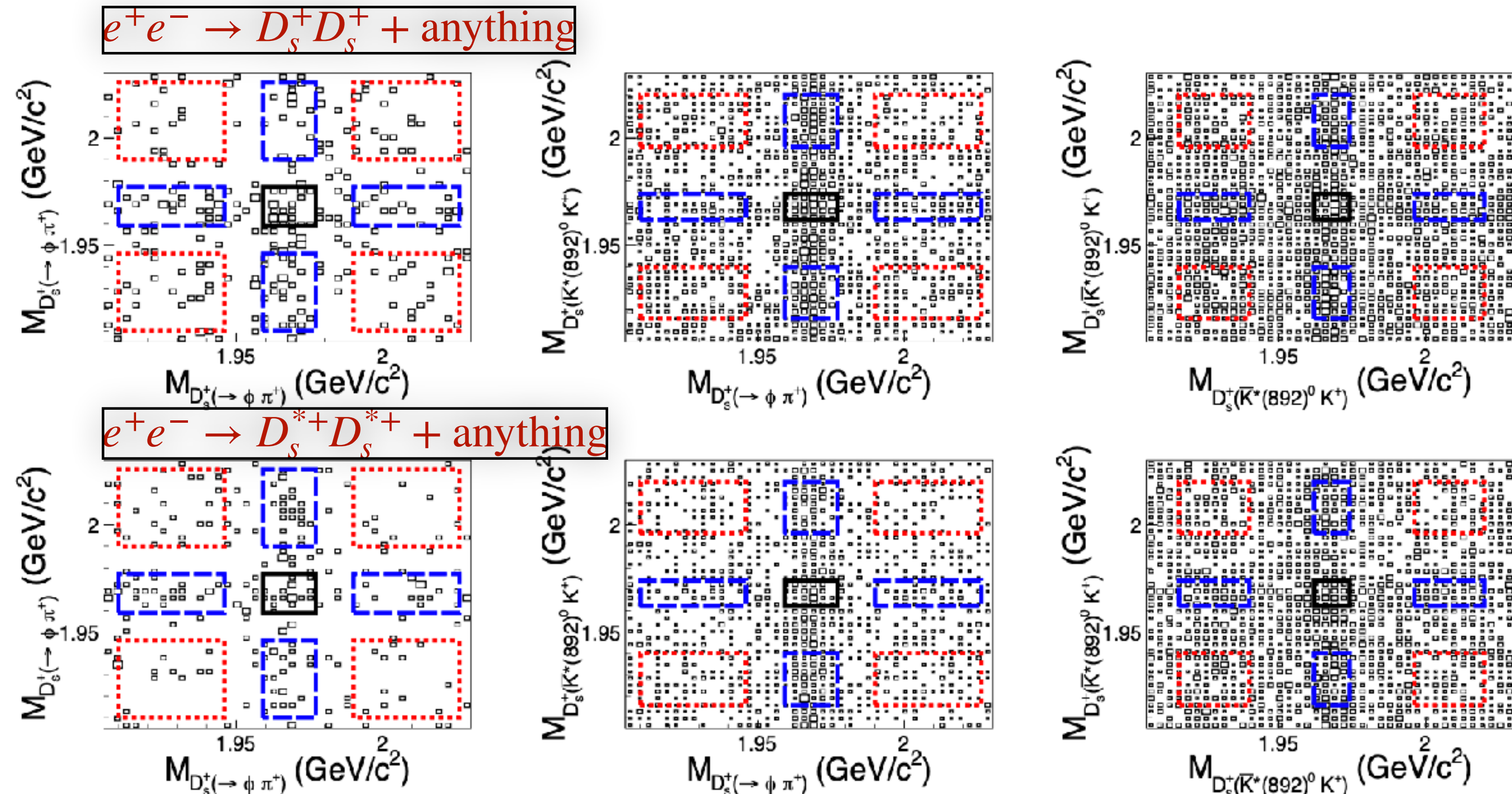
Search for tetraquark states $X_{cc\bar{s}\bar{s}}$ in $D_s^+ D_s^+$, $D_s^{*+} D_s^{*+}$

predictions

Mode	IJ^P	Mass (MeV/ c^2)	Width (MeV)
$X_{cc\bar{s}\bar{s}} \rightarrow D_s^+ D_s^+$	00^+	4902	3.54
$X_{cc\bar{s}\bar{s}} \rightarrow D_s^{*+} D_s^{*+}$	02^+	4821	5.58
	02^+	4846	10.68
	02^+	4775	23.26

Data samples:

- 5.74 fb⁻¹ on the $\Upsilon(1S)$ resonance
- 24.7 fb⁻¹ on the $\Upsilon(2S)$ resonance
- 89.5 fb⁻¹ at $\sqrt{s} = 10.52$ GeV
- 711 fb⁻¹ at $\sqrt{s} = 10.58$ GeV
- 121.4 fb⁻¹ at $\sqrt{s} = 10.867$ GeV

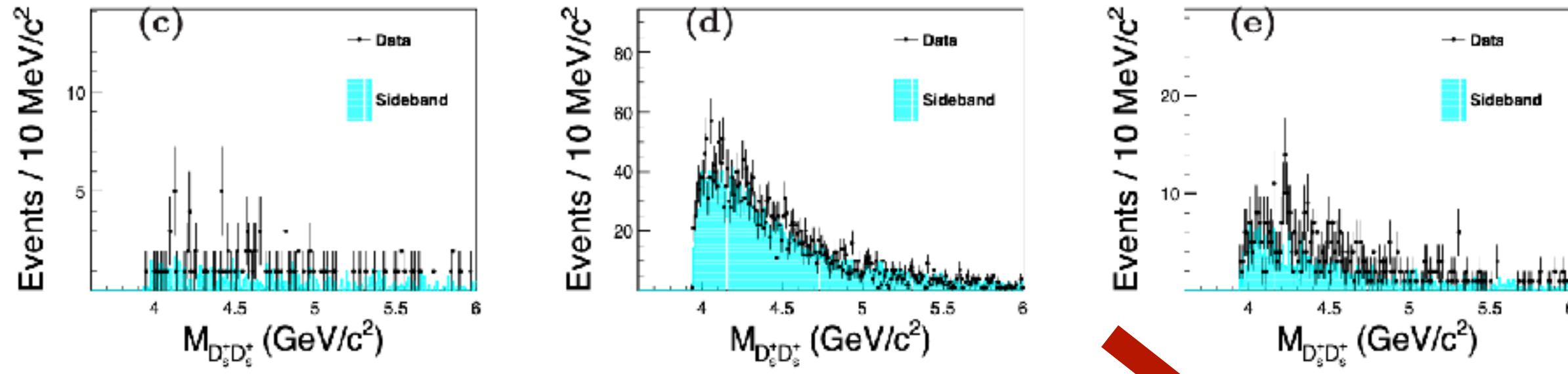
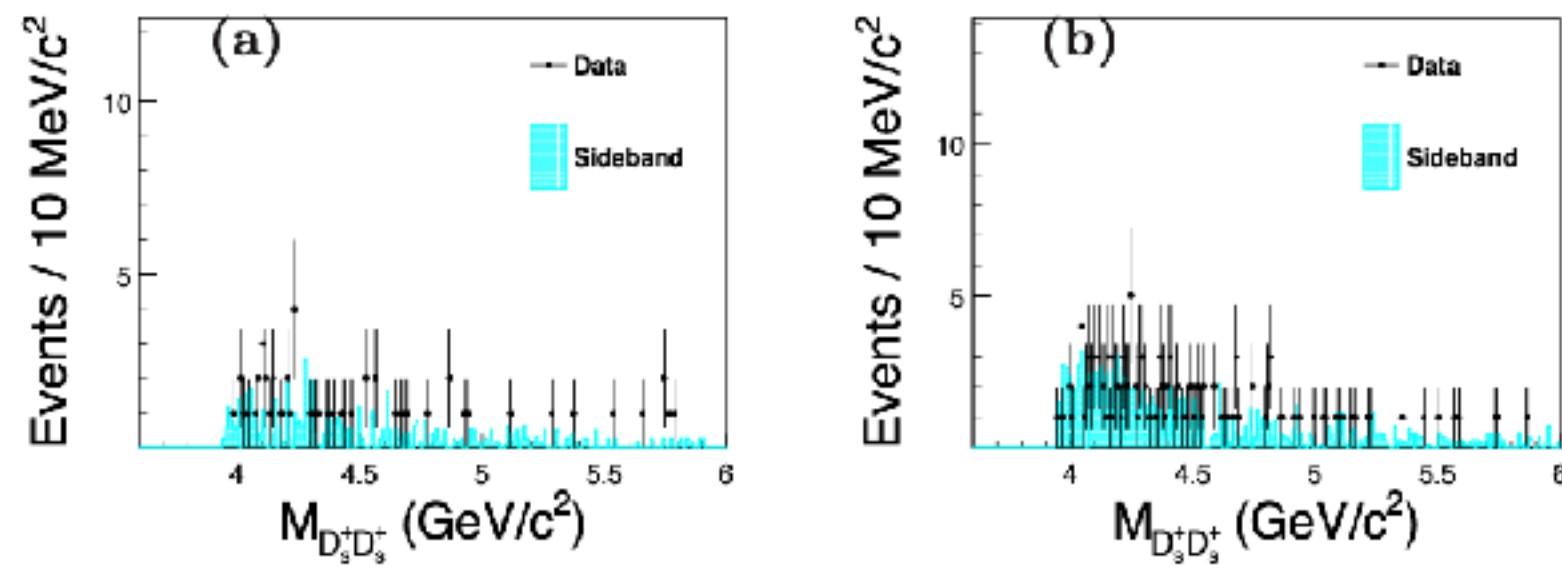


Use two modes to reconstruct D_s^+

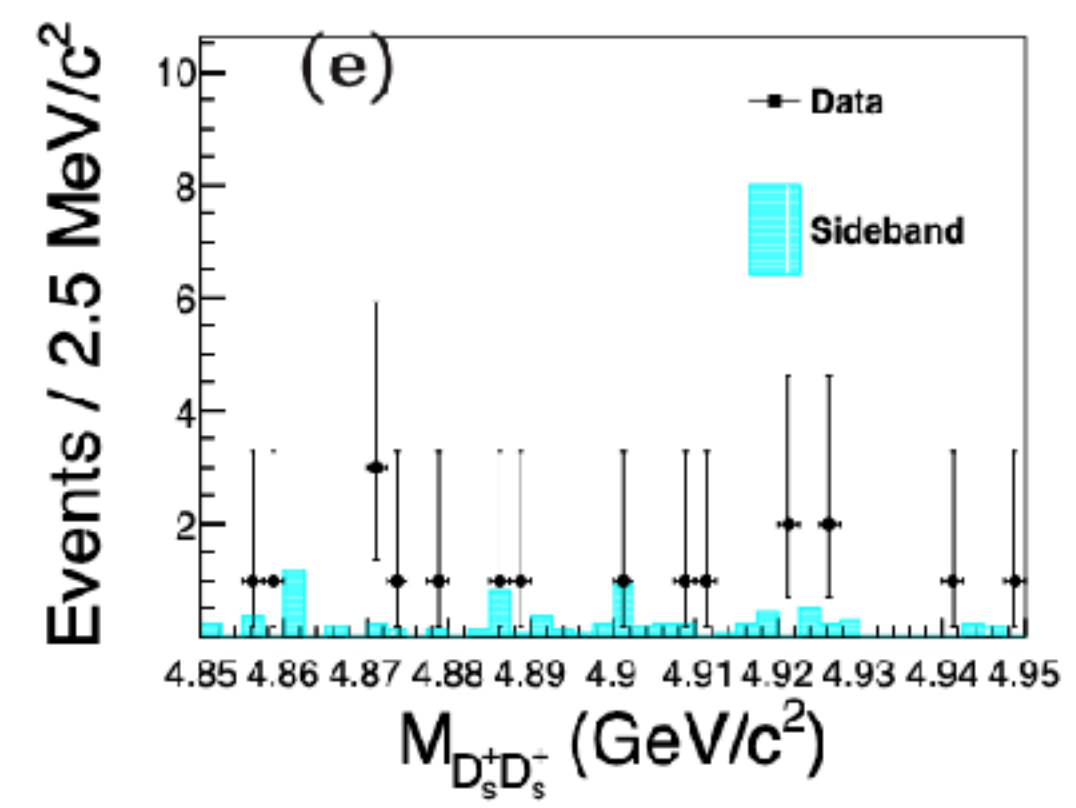
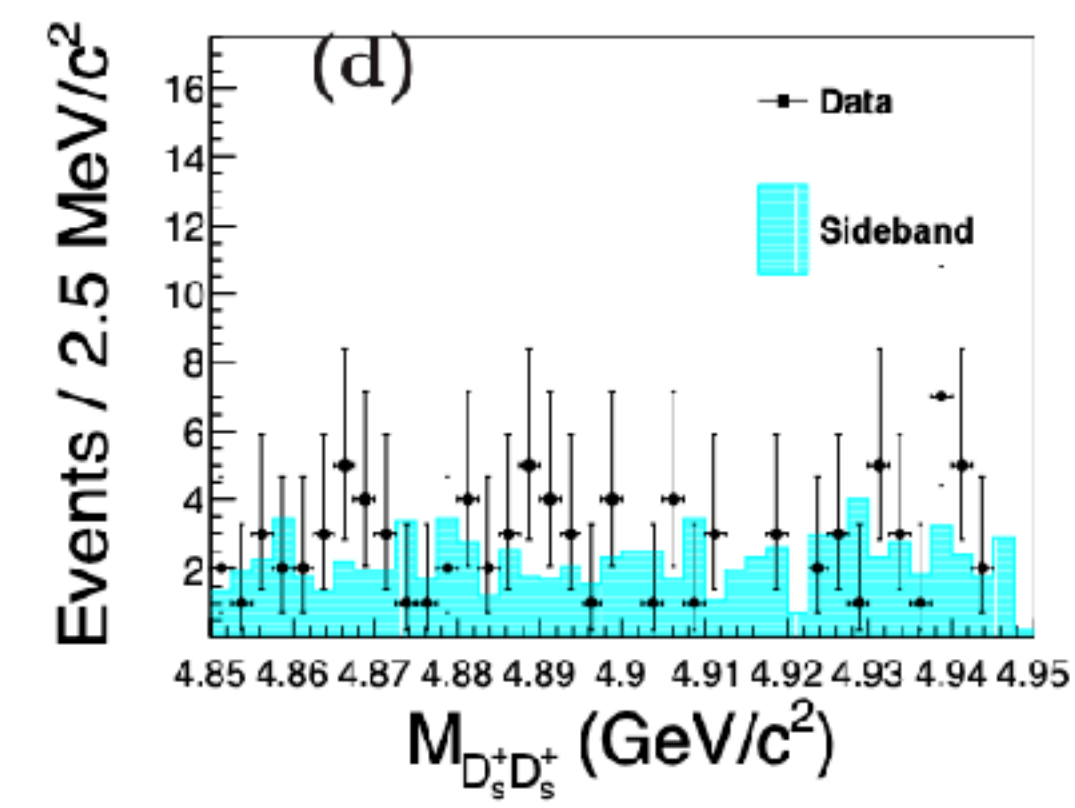
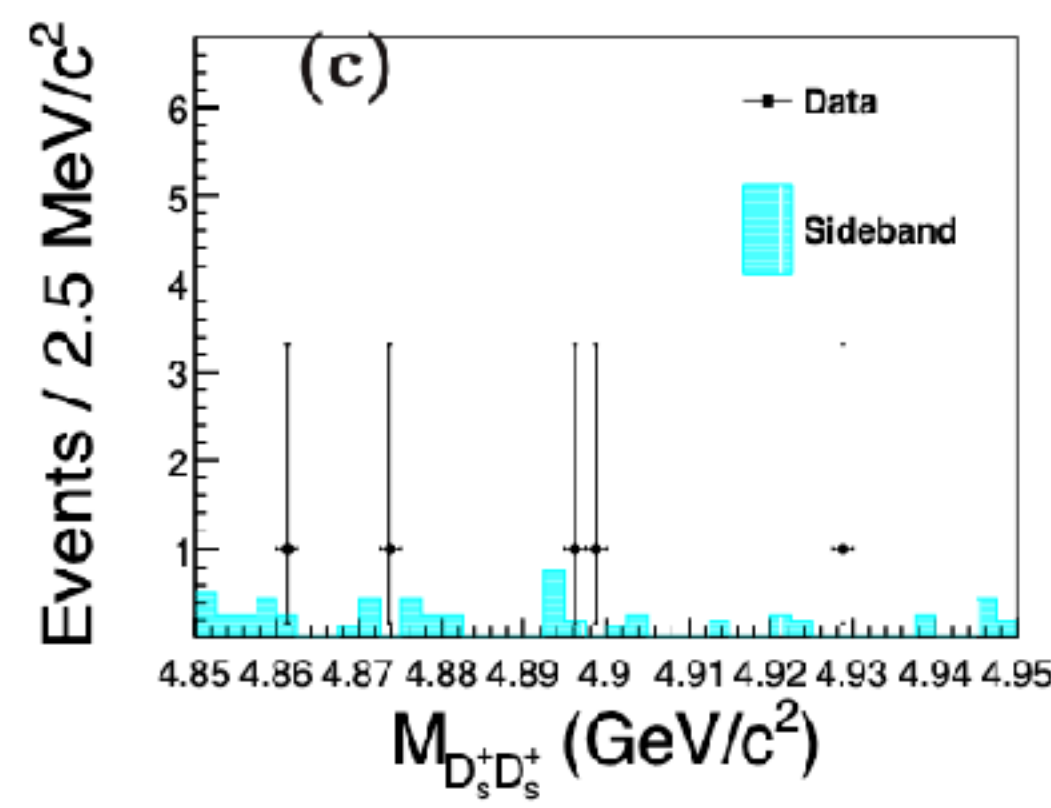
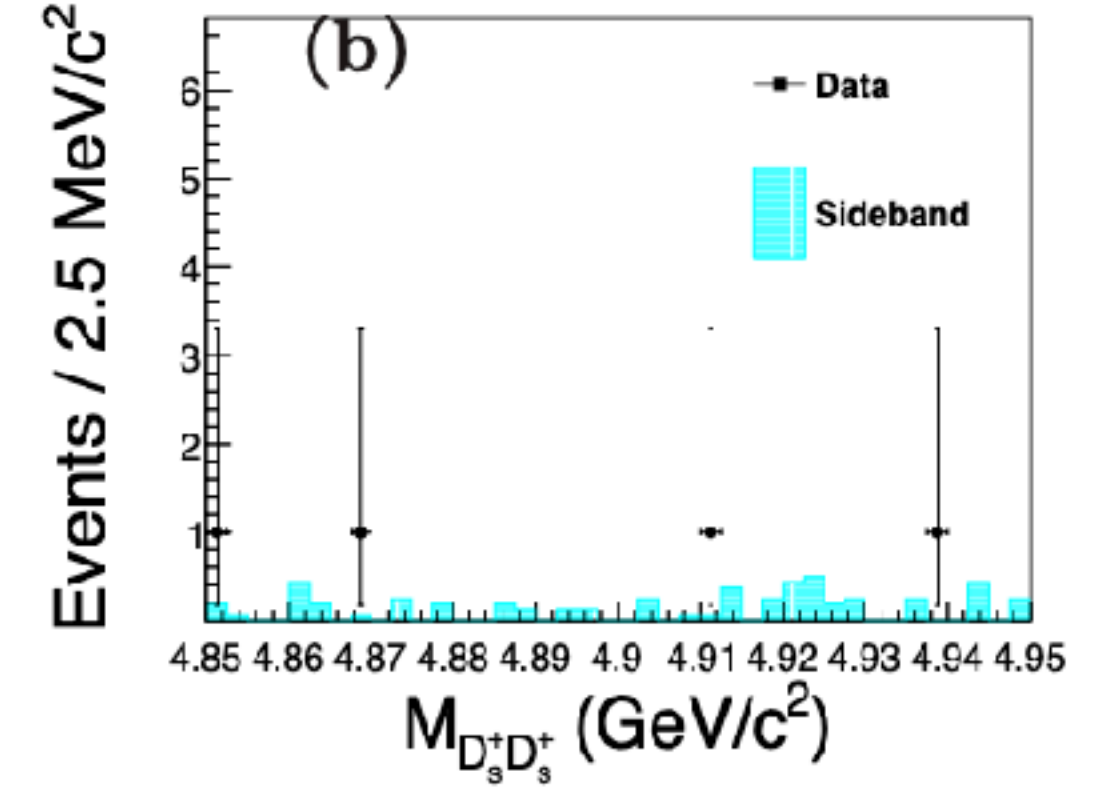
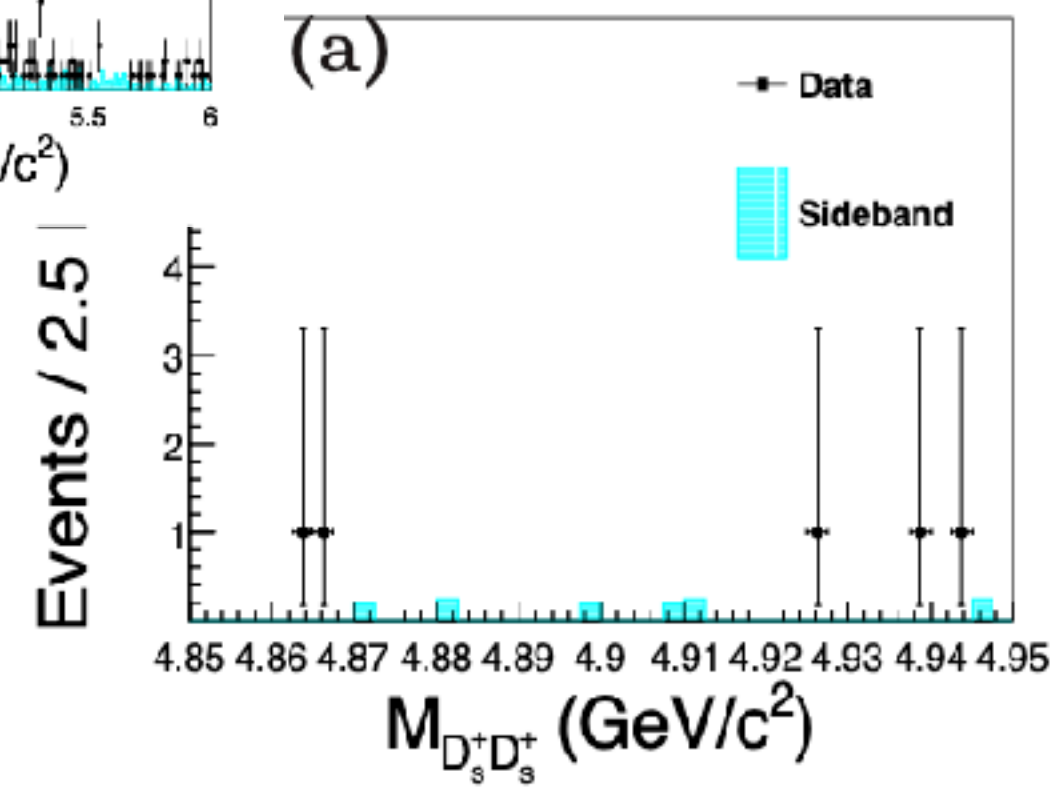
- $D_s^+ \rightarrow K^*(892)K^+$
- $D_s^+ \rightarrow \phi\pi^+$

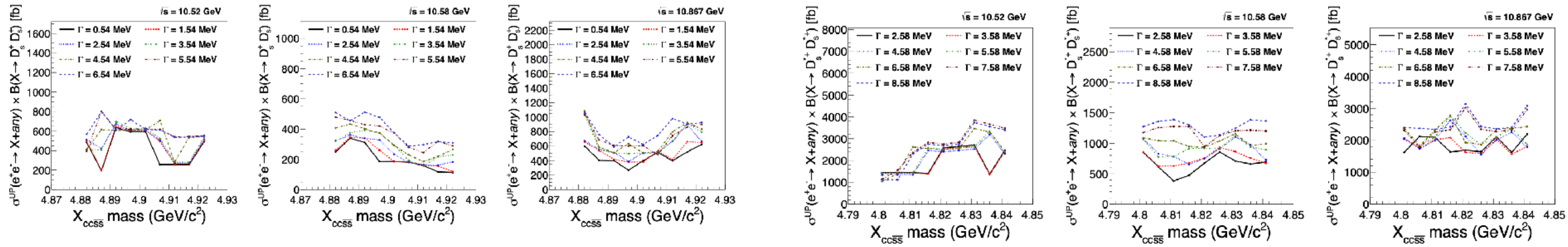
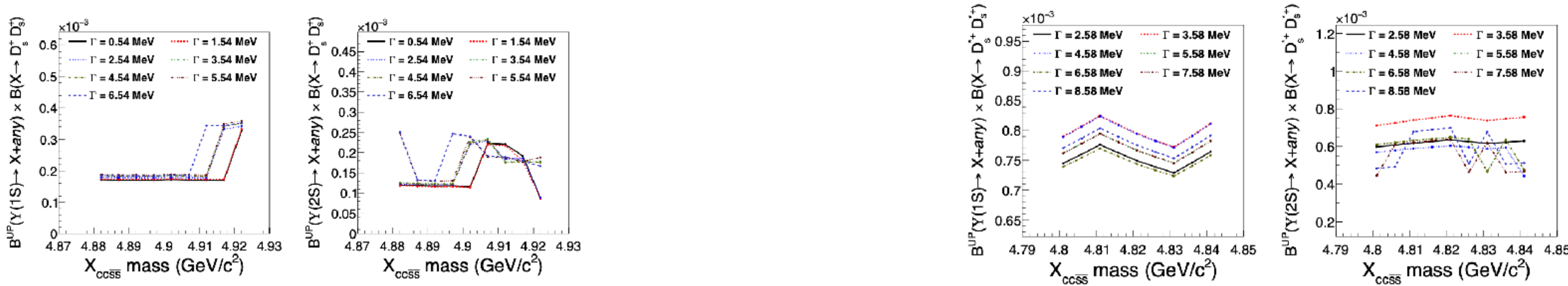
$$e^+e^- \rightarrow D_s^+ D_s^+ + \text{anything}$$

No peaking background in *sideband*



Focus on prediction regions for $X_{cc\bar{s}\bar{s}}$





No evident signal from $D_s^+ D_s^+$ or $D_s^{*+} D_s^{*+}$.
 Upper limits in 90% C.L. are estimated in different mass and width assumptions.

Others

Search for $\eta_{c2}(1D)$ in the $e^+e^- \rightarrow \gamma\eta_{c2}(1D)$

Phys.Rev.D 104 (2021) 012012

Study of $\chi_{bJ}(nP) \rightarrow \omega\Upsilon(1S)$ at Belle

arXiv: 2108.03497

Search for a doubly-charged DDK bound state

Phys.Rev.D 102 (2020) 11, 112001

Evidence of $\Omega_c \rightarrow \pi^+ \Omega(2012)^- \rightarrow \pi^+ (\bar{K} \Xi)^-$

- (Belle 2018) observation of $\Omega(2012)^- \rightarrow K^- \Xi^0, K_S^0 \Xi^-$
- Interpretations?
 - Is it a $(\bar{K} \Xi(1530))^-$ molecule, or not?
 - If molecule \rightarrow large decay width of $\Omega(2012)^- \rightarrow (\bar{K} \pi \Xi)^-$
- Prediction
 - $\Omega(2012)^-$ would be much more visible in $\Omega_c^0 \rightarrow \pi^+ (\bar{K} \Xi)^-$

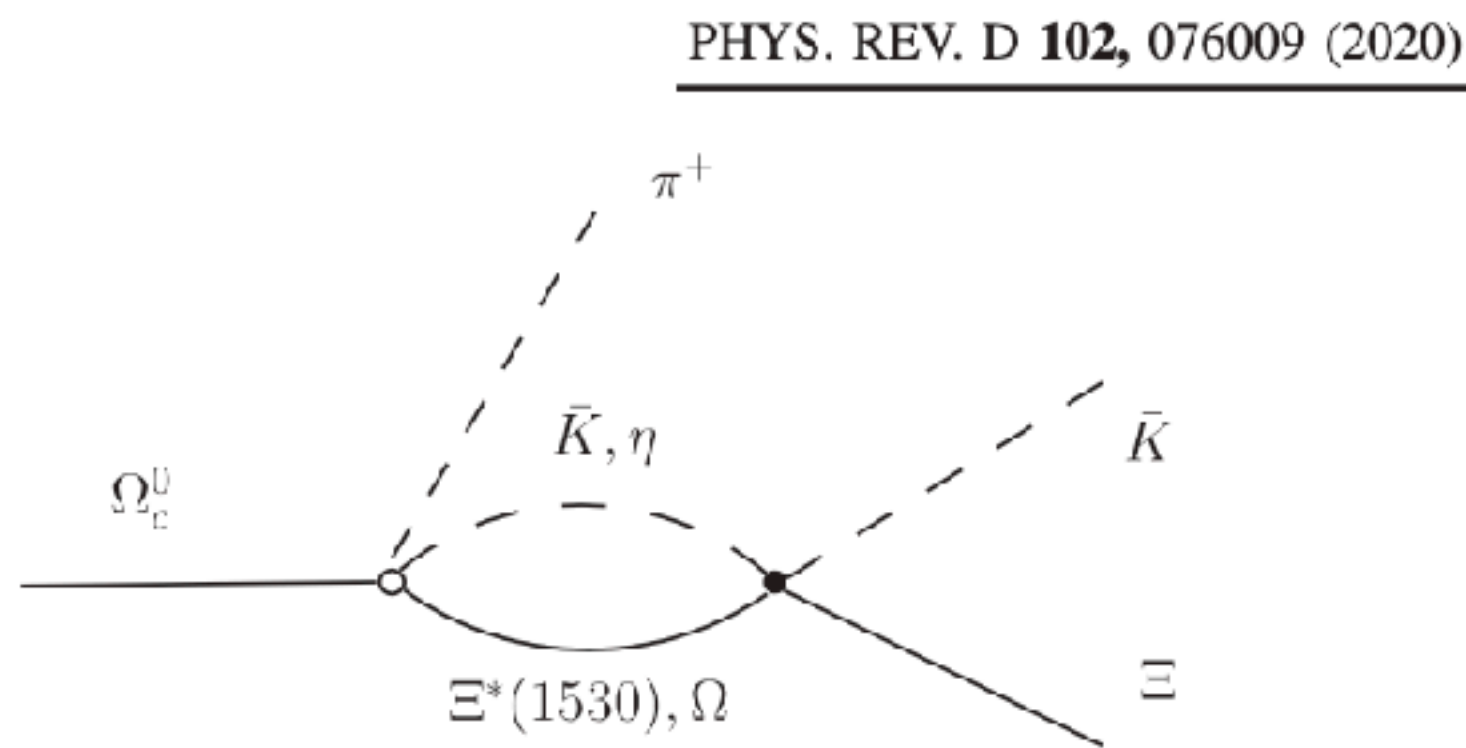


FIG. 2. Diagram for the meson-baryon final-state interaction for the $\Omega_c^0 \rightarrow \pi^+ \Omega(2012)^- \rightarrow \pi^+ (\bar{K} \Xi)^-$ decay.

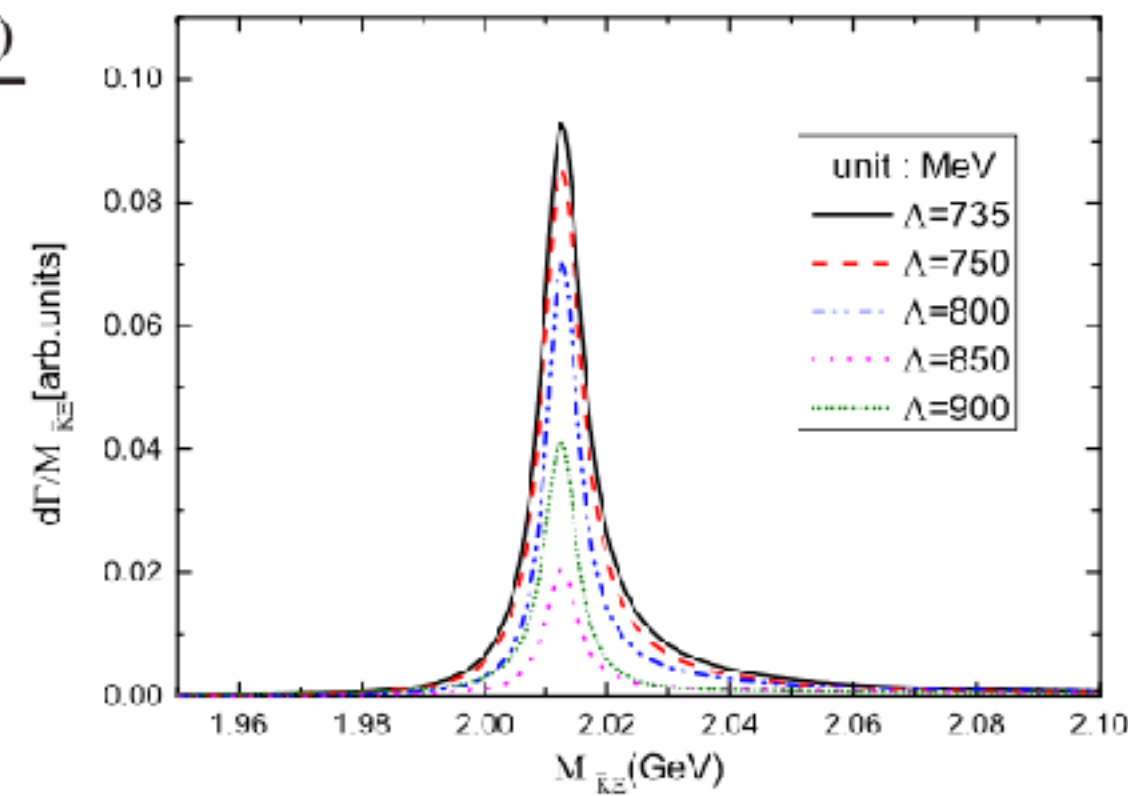


FIG. 3. $\bar{K} \Xi$ invariant mass distributions of the $\Omega_c^0 \rightarrow \pi^+ (\bar{K} \Xi)^-$ decay.

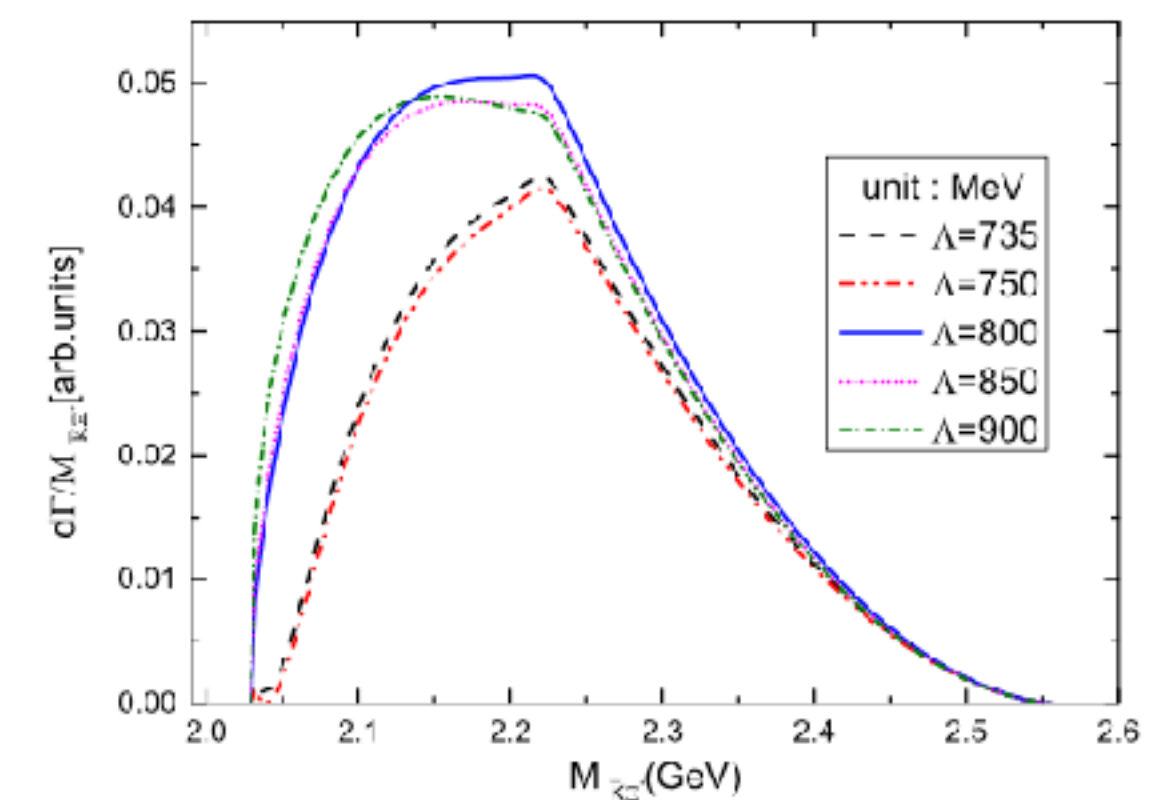
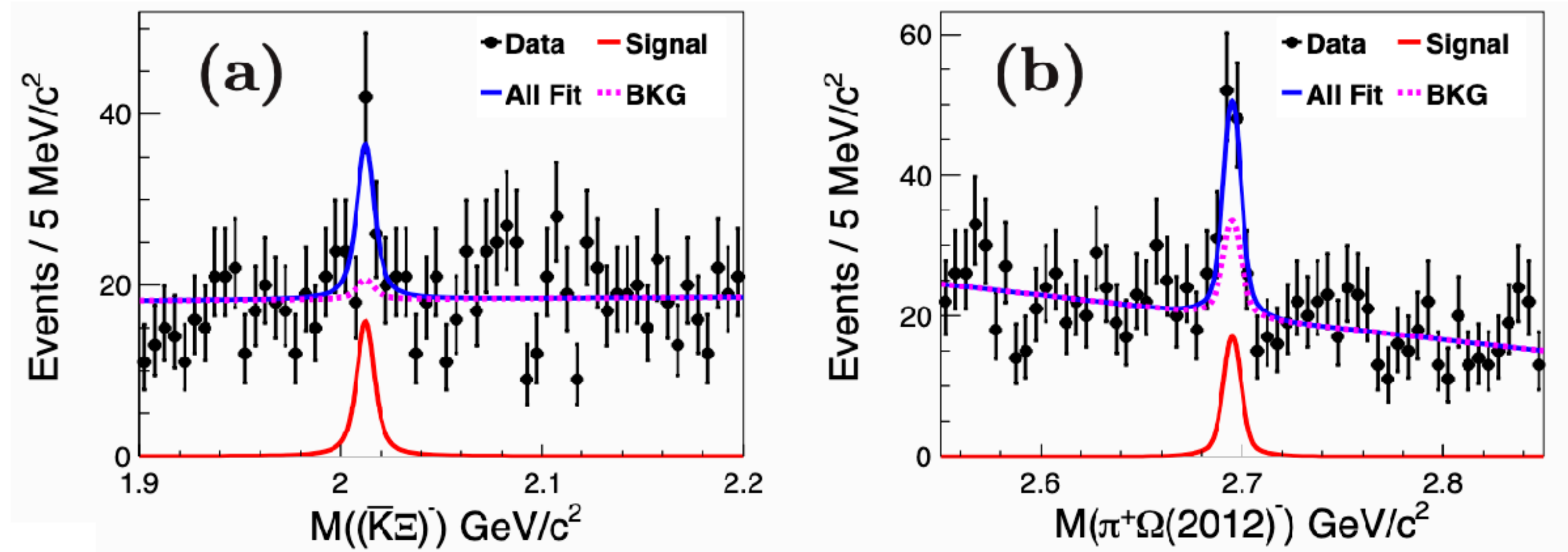


FIG. 4. $\bar{K} \Xi^*$ invariant mass distribution of the $\Omega_c^0 \rightarrow \pi^+ (\bar{K} \Xi^*)^-$ decay.

2D fit to $M(\bar{K}\Xi)$ vs $M(\pi^+\Omega(2012)^-)$ simultaneously for $M(K^-\Xi^0)$ and $M(K_S^0\Xi^-)$.

Entire Belle data – 980 fb⁻¹



$$N_{\text{fit}} = 46.6 \pm 12.3$$

Signal significance: 4.2 σ
(including systematic uncertainties)

$$\frac{\mathcal{B}(\Omega_c^0 \rightarrow \pi^+\Omega(2012)^- \rightarrow \pi^+(\bar{K}\Xi)^-)}{\mathcal{B}(\Omega_c^0 \rightarrow \pi^+\Omega^-)} = 0.220 \pm 0.059 \pm 0.035$$

$$\frac{\mathcal{B}(\Omega_c^0 \rightarrow \pi^+\Omega(2012)^- \rightarrow \pi^+K^-\Xi^0)}{\mathcal{B}(\Omega_c^0 \rightarrow \pi^+K^-\Xi^0)} = (9.6 \pm 3.2 \pm 1.8)\%$$

$$\frac{\mathcal{B}(\Omega_c^0 \rightarrow \pi^+\Omega(2012)^- \rightarrow \pi^+K_S^0\Xi^-)}{\mathcal{B}(\Omega_c^0 \rightarrow \pi^+K_S^0\Xi^-)} = (5.5 \pm 2.8 \pm 0.7)\%$$

Spin parity of $\Xi_c(2970)^+$

- \exists no determination of spin-parity of Ξ_c baryons

Ξ_c^+	$I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$	$\Xi_c(2970)$	$I(J^P) = \frac{1}{2}(??)$
J^P has not been measured; $\frac{1}{2}^+$ is the quark-model prediction.		was $\Xi_c(2980)$	
Mass $m = 2467.71 \pm 0.23$ MeV (S = 1.3)		$\Xi_c(2970)^+$ $m = 2964.3 \pm 1.5$ MeV (S = 3.9)	
Mean life $\tau = (456 \pm 5) \times 10^{-15}$ s		$\Xi_c(2970)^0$ $m = 2967.1 \pm 1.7$ MeV (S = 6.7)	

- We study $\Xi_c(2970)^+$ because

- wide variety (of J^P) and controversial
- many predicted states within ~50 MeV

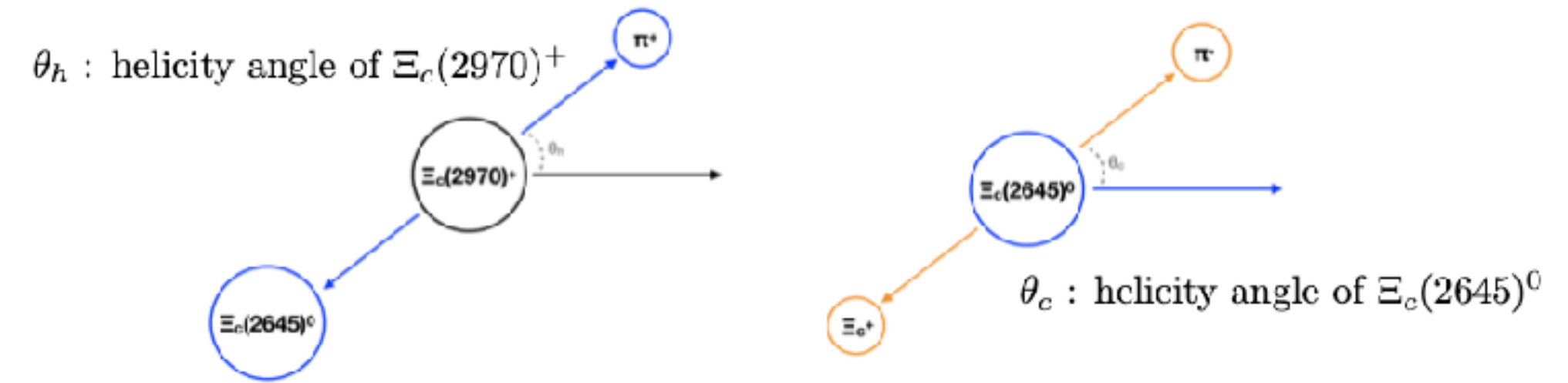
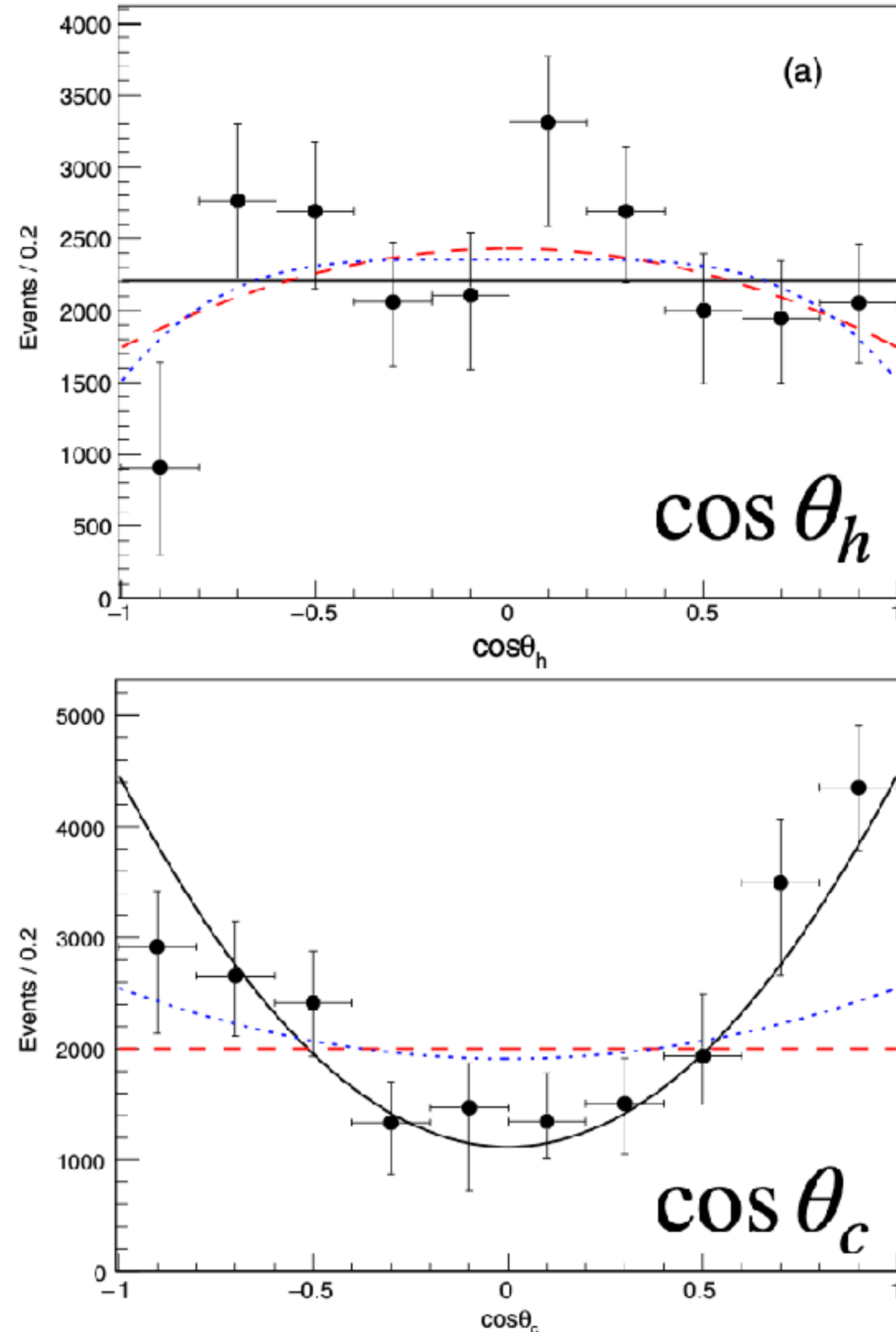
- J^P has not been assigned by PDG
- Belle, with high-stat. $\Xi_c(2970)^+ \rightarrow \Xi_c(2645)^0 \pi^+ \rightarrow \Xi_c^+ \pi^+ \pi^-$ and clean e^+e^- setting, is an ideal place to measure J^P of $\Xi_c(2970)^+$

- \exists many theory predictions

- $J^P = 1/2^+, 3/2^+, 5/2^+, 5/2^-$, and even predictions of (-) parity
- so, we want to measure it and help decipher the nature of the state

Spin & Parity of $\Xi_c(2970)^+$

Phys.Rev.D 103 (2021) 11, L111101



J^P	Partial wave	$W(\theta_c)$
$1/2^+$	P	$1 + 3\cos^2\theta_c$
$1/2^-$	D	$1 + 3\cos^2\theta_c$
$3/2^+$	P	$1 + 6\sin^2\theta_c$
$3/2^-$	S	1
$5/2^+$	P	$1 + (1/3)\cos^2\theta_c$
$5/2^-$	D	$1 + (15/4)\sin^2\theta_c$

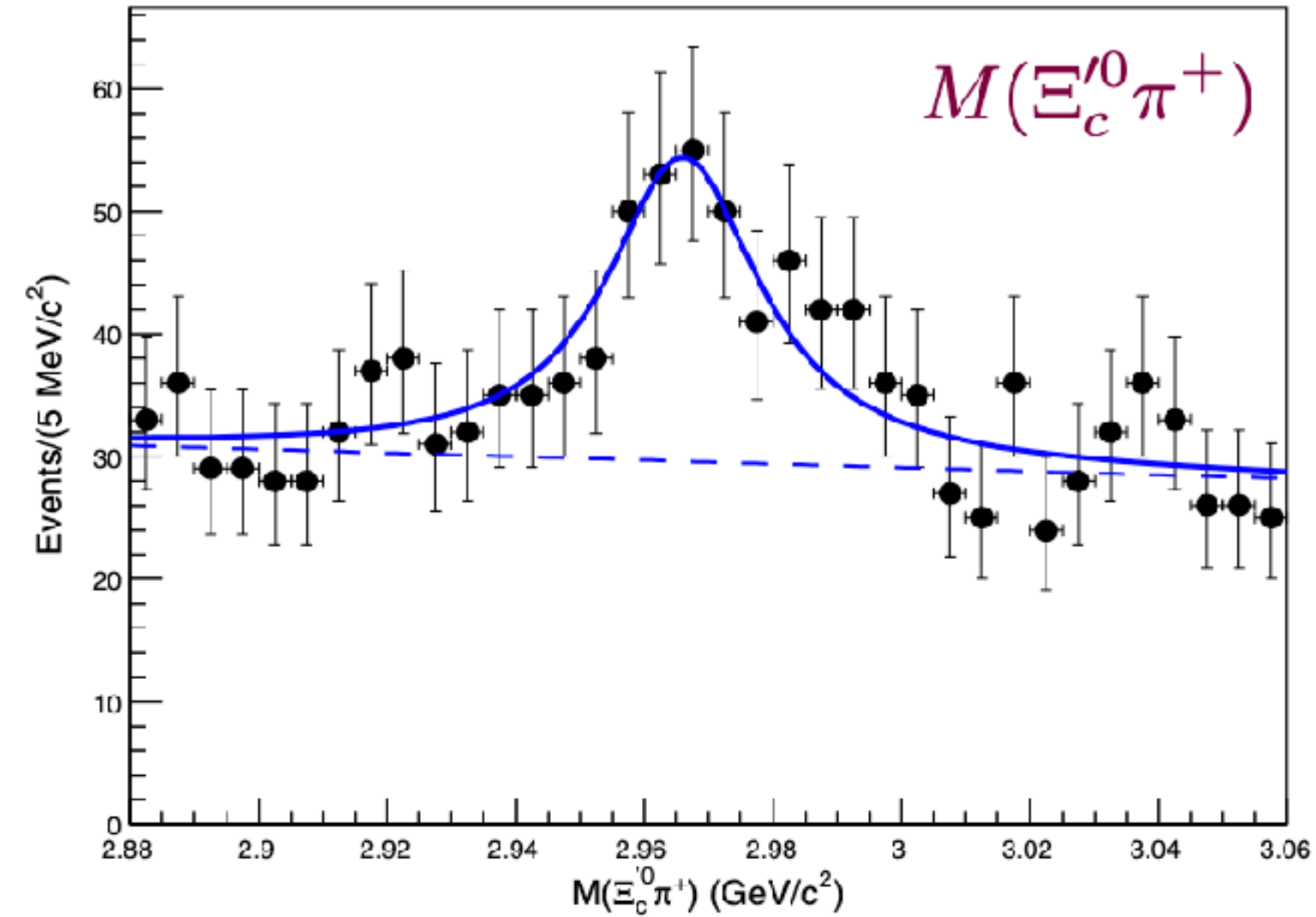
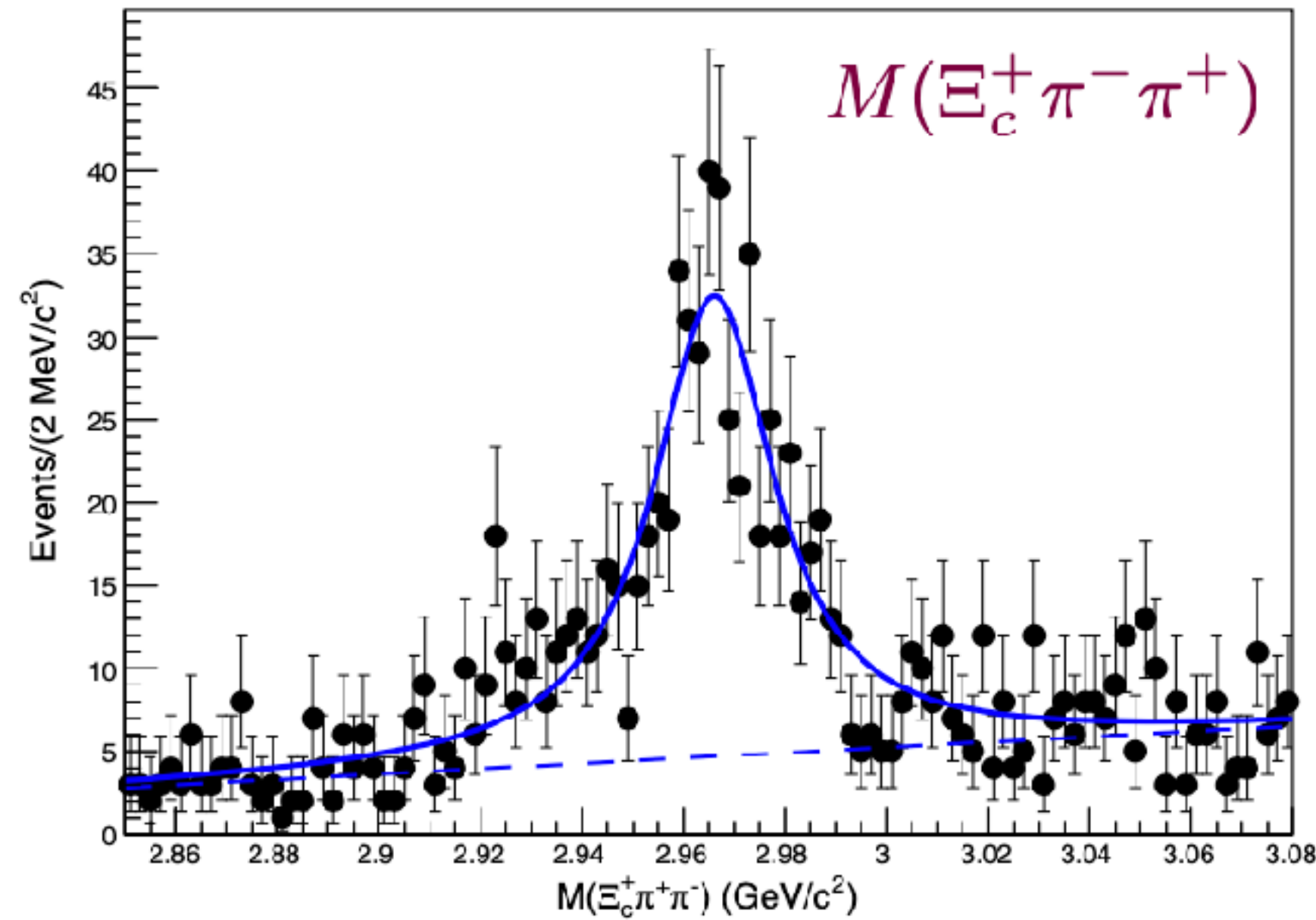
J^P	$1/2^\pm$	$3/2^-$	$5/2^+$
$\chi^2/\text{n.d.f.}$	6.4/9	32.2/9	22.3/9
Exclusion level (s.d.)	...	5.5	4.8

- most consistent with spin=1/2 hypothesis
- also excludes $\Xi_c(2645)$ spin of 1/2 ($\because \cos \theta_c$ not flat)

Spin & Parity of $\Xi_c(2970)^+$

Phys.Rev.D 103 (2021) 11, L111101

$$R = \frac{\mathcal{B}[\Xi_c(2970)^+ \rightarrow \Xi_c(2645)^0 \pi^+]}{\mathcal{B}[\Xi_c(2970)^+ \rightarrow \Xi_c'{}^0 \pi^+]} \quad \text{--- sensitive to parity}$$



$$R = 1.67 \pm 0.29_{-0.09}^{+0.15} \pm 0.25 \text{ (IS)}$$

Heavy-quark spin sym. prediction

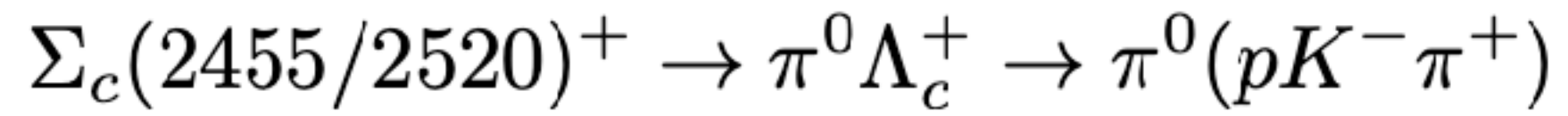
Parity	+	+	-	-
Brown-muck spin s_ℓ	0	1	0	1
R	1.06	0.26	0	$\ll 1$

$\therefore (+)$ parity assignment is favored

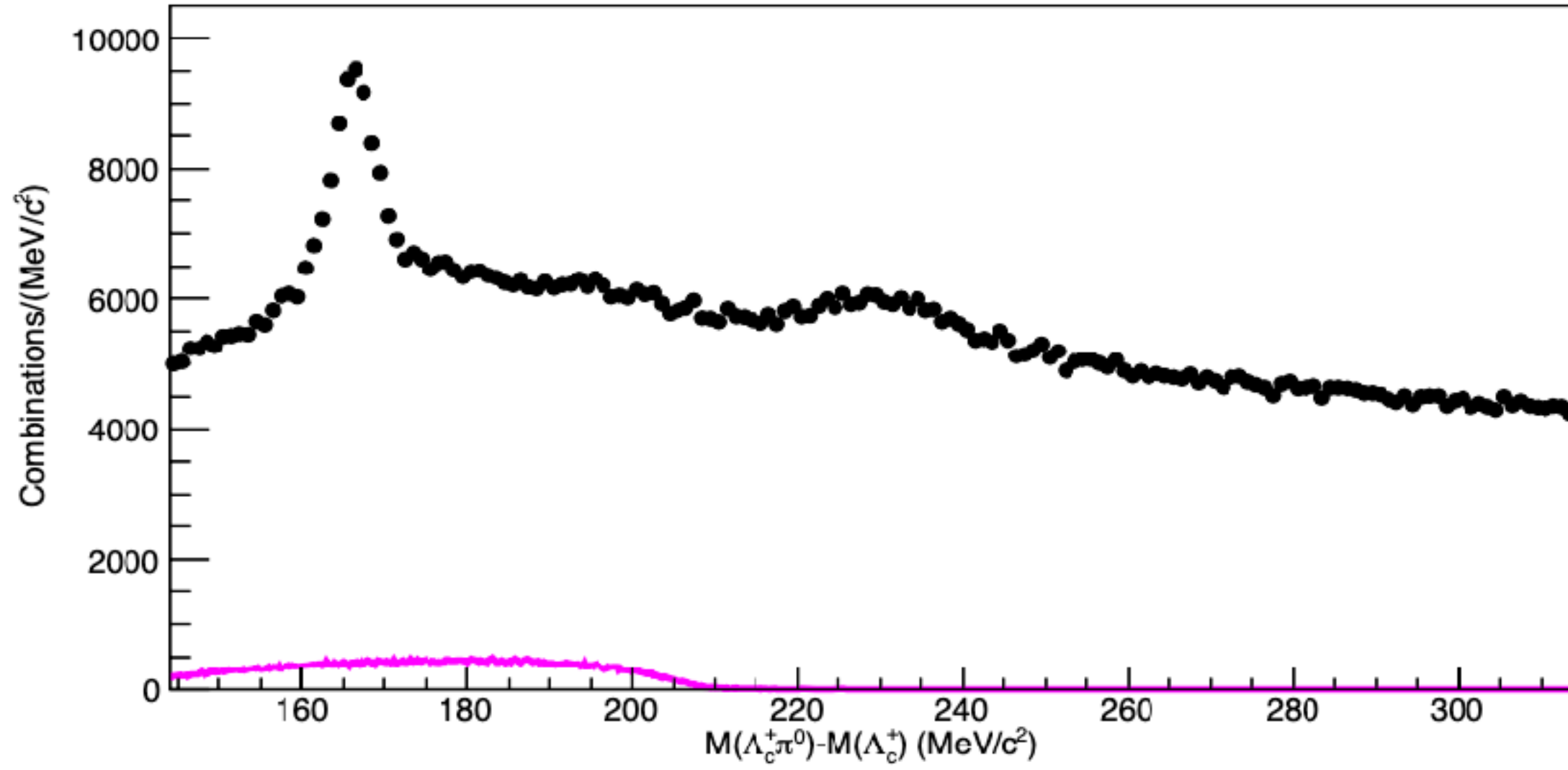
Mass and width of $\Sigma_c^{(*)+}$

- Σ_c baryons = a c quark + spin-1 light diquark (uu , ud or dd)
 - lowest : $\Sigma_c(2455)$ triplet, with $J^P = (1/2)^+$ → decay to $\Lambda_c^+\pi$
 - next up: $\Sigma_c(2520)$ triplet, with $J^P = (3/2)^+$ → decay to $\Lambda_c^+\pi$
 - $\Sigma_c^{++/0}$ mass, width — well measured for both charges, but
 - Σ_c^+ — mass only from CLEO II, and limit only for widths
- Mass measurements of the two isotriplets
 - allow tests of models of isospin mass splittings
- Predictions
 - most mass models: $m(\Sigma_c^+) < m(\Sigma_c^{0/++})$
 - natural width models: $\Gamma(\Sigma_c^+) > \Gamma(\Sigma_c^{0/++})$

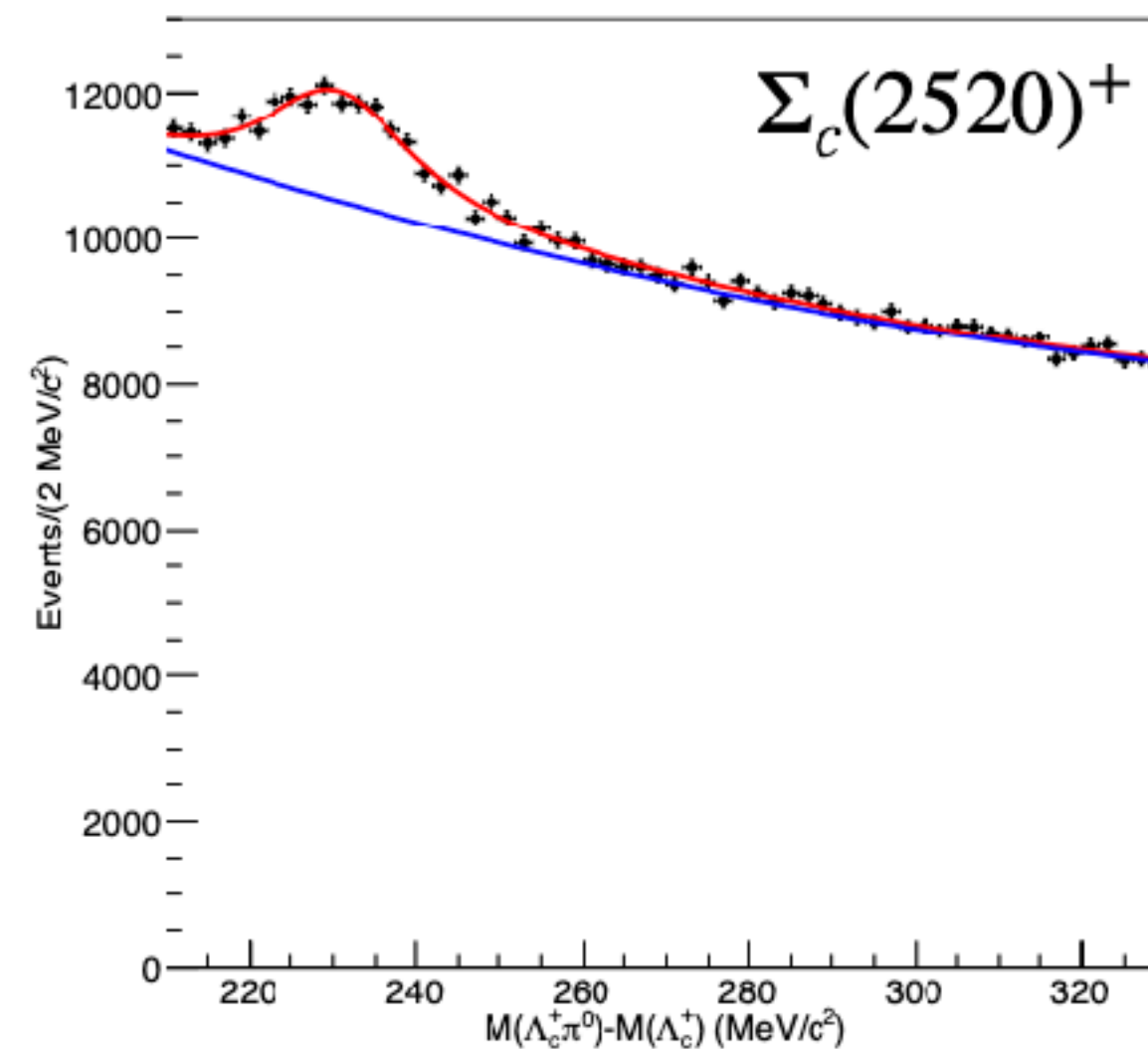
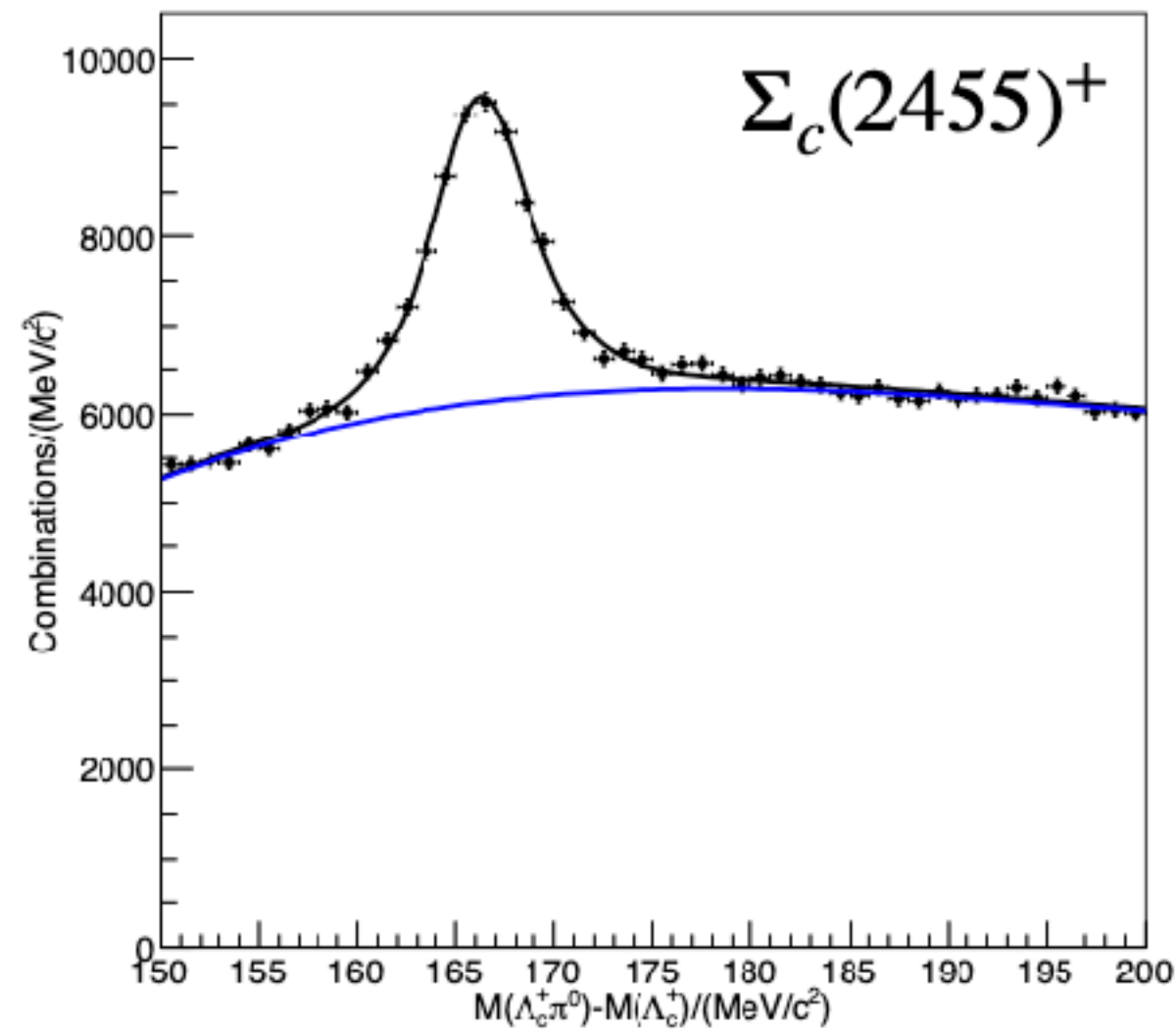
[4] G.-S. Yang and H.-C. Kim, Phys. Lett. B 808, 135619 (2020).



Phys.Rev.D 104 (2021) 5, 052003



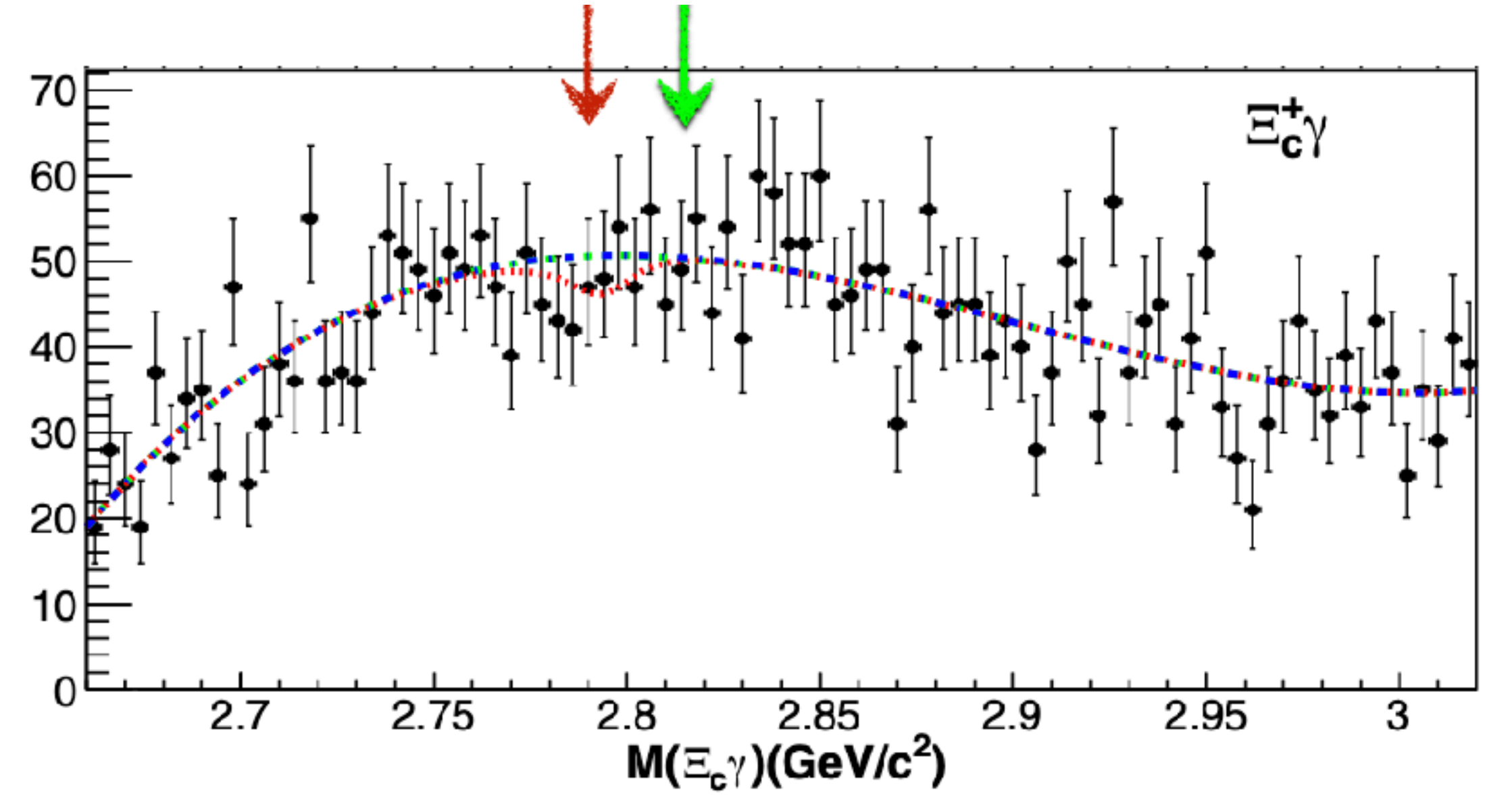
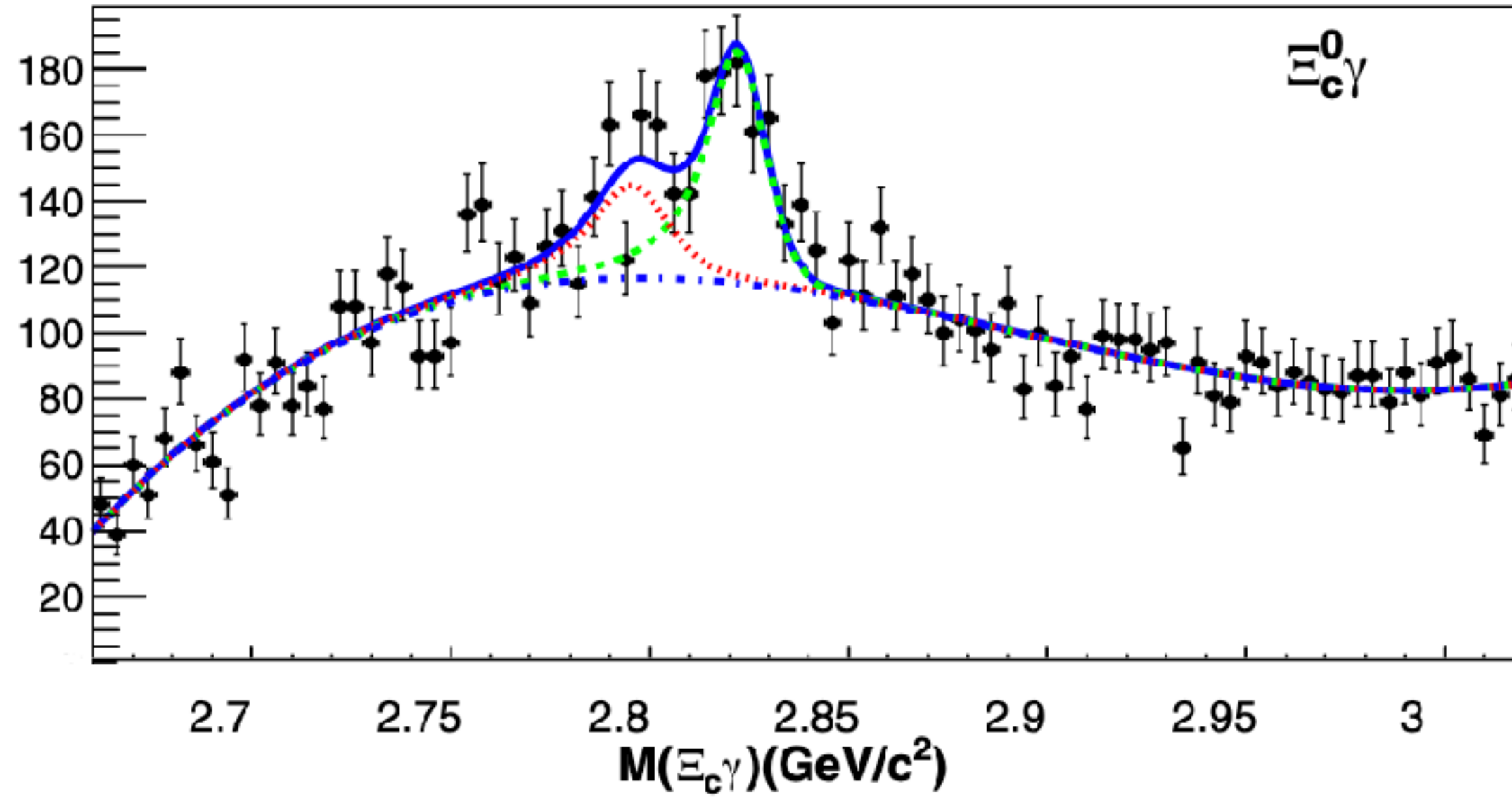
	$\Sigma_c(2455)^+$	$\Sigma_c(2520)^+$
ΔM [MeV/c ²]	$166.17 \pm 0.05^{+0.16}_{-0.07}$	$230.9 \pm 0.5^{+0.5}_{-0.1}$
Γ [MeV/c ²]	$2.3 \pm 0.3 \pm 0.3$	$17.2^{+2.3+3.1}_{-2.1-0.7}$



- First measurement of widths of $\Sigma_c(2455)^+$, $\Sigma_c(2520)^+$
- Much improved precision of $m(\Sigma_c(2455)^+)$, $m(\Sigma_c(2520)^+)$
- Measured masses and widths are consistent with theory predictions

Radiative decay of $\Xi_c(2790/2815)$

- Recently measured $\Xi_c(2790)^{+/\ 0}$ & $\Xi_c(2815)^{+/\ 0}$ masses and widths
 - In the picture of $(c + ud, us)$, these are typically interpreted as $L = 1$ orbital excitations (“ λ ”).
 - The nature of these states are identified by mass spectra and decay modes.
- Excited charmed baryons mostly decay via strong interactions.
 - the only observed EM decays : $\Xi'_c \rightarrow \Xi_c \gamma$, $\Omega_c(2770) \rightarrow \Omega_c \gamma$
- Wang, Yao, Zhong, Zhao (PRD 96, 116016 (2017)) predicts
 - assuming λ excitations, large widths of $\Xi_c(2790)^0 \rightarrow \Xi_c^0 \gamma$, $\Xi_c(2815)^0 \rightarrow \Xi_c^0 \gamma$ ($\Gamma \gtrsim 200$ keV)
 - assuming ρ excitations (between the two light quarks), much smaller widths (< 10 keV) for the Ξ_c^+ baryons



$$R_{2790}^0 = \frac{\mathcal{B}[\Xi_c(2790)^0 \rightarrow \Xi_c^0 \gamma]}{\mathcal{B}[\Xi_c(2790)^0 \rightarrow \Xi_c'^+ \pi^- \rightarrow \Xi_c^+ \gamma \pi^-]} = 0.13 \pm 0.03 \pm 0.02$$

$$R_{2790}^+ = \frac{\mathcal{B}[\Xi_c(2790)^+ \rightarrow \Xi_c^+ \gamma]}{\mathcal{B}[\Xi_c(2790)^+ \rightarrow \Xi_c'^0 \pi^+ \rightarrow \Xi_c^0 \gamma \pi^+]} < 0.06 @ 90\% \text{ C.L.}$$

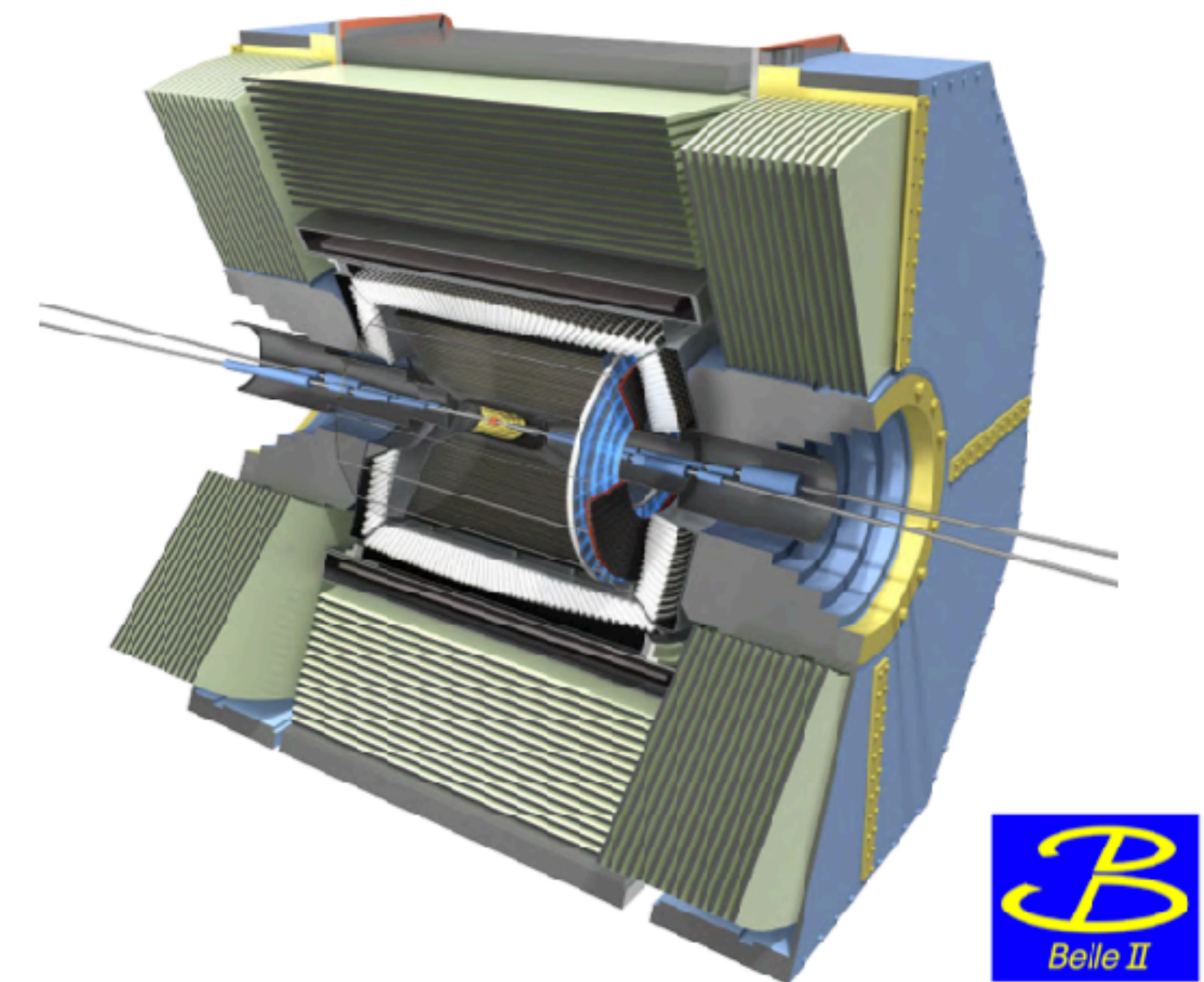
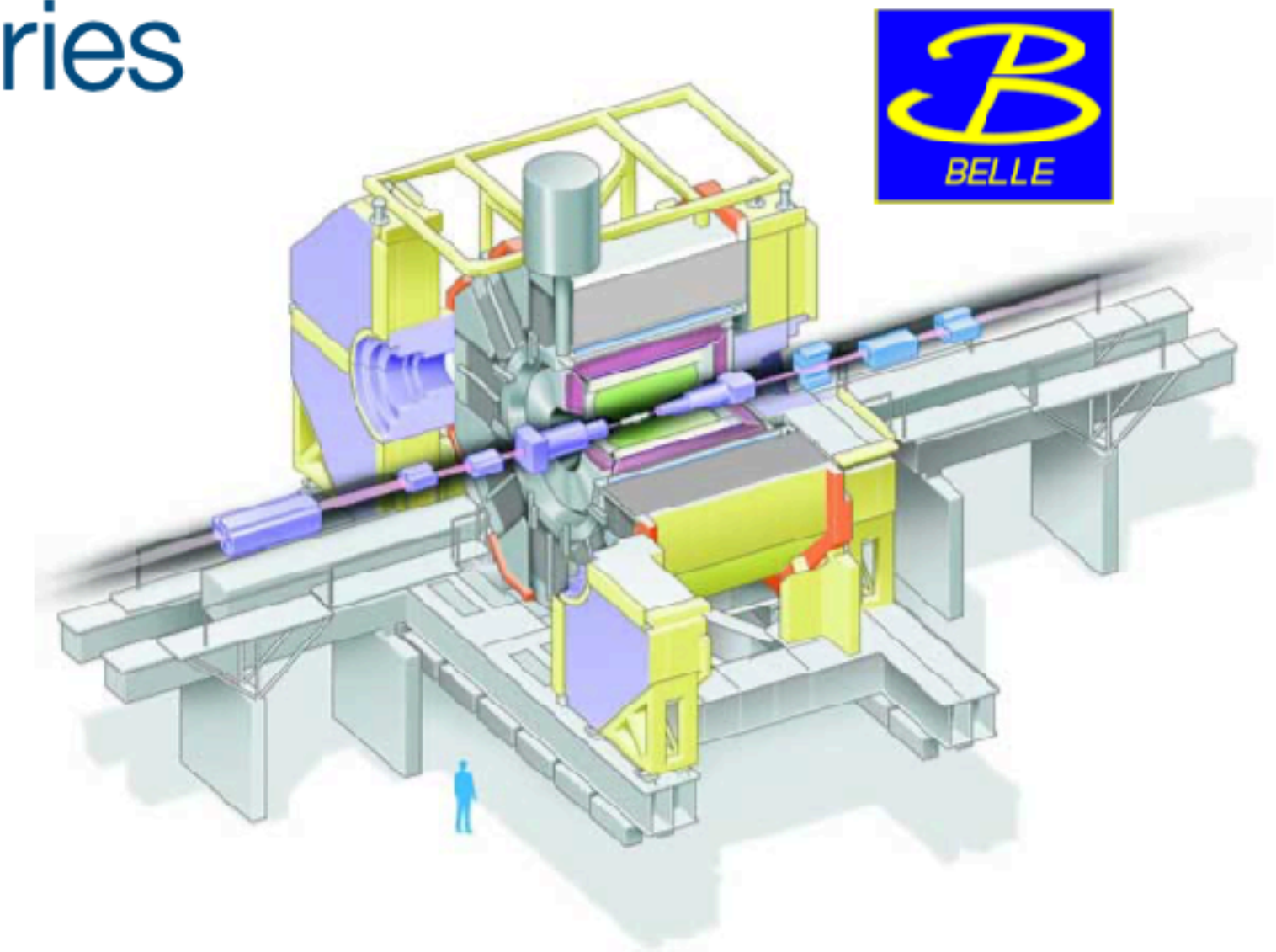
$$R_{2815}^0 = \frac{\mathcal{B}[\Xi_c(2815)^0 \rightarrow \Xi_c^0 \gamma]}{\mathcal{B}[\Xi_c(2815)^0 \rightarrow \Xi_c(2645)^+ \pi^- \rightarrow \Xi_c^0 \pi^+ \pi^-]} = 0.41 \pm 0.05 \pm 0.03$$

$$R_{2815}^+ = \frac{\mathcal{B}[\Xi_c(2815)^+ \rightarrow \Xi_c^+ \gamma]}{\mathcal{B}[\Xi_c(2815)^+ \rightarrow \Xi_c(2645)^0 \pi^+ \rightarrow \Xi_c^+ \pi^+ \pi^-]} < 0.09 @ 90\% \text{ C.L.}$$

- First observation of radiative decays of orbitally excited Ξ_c

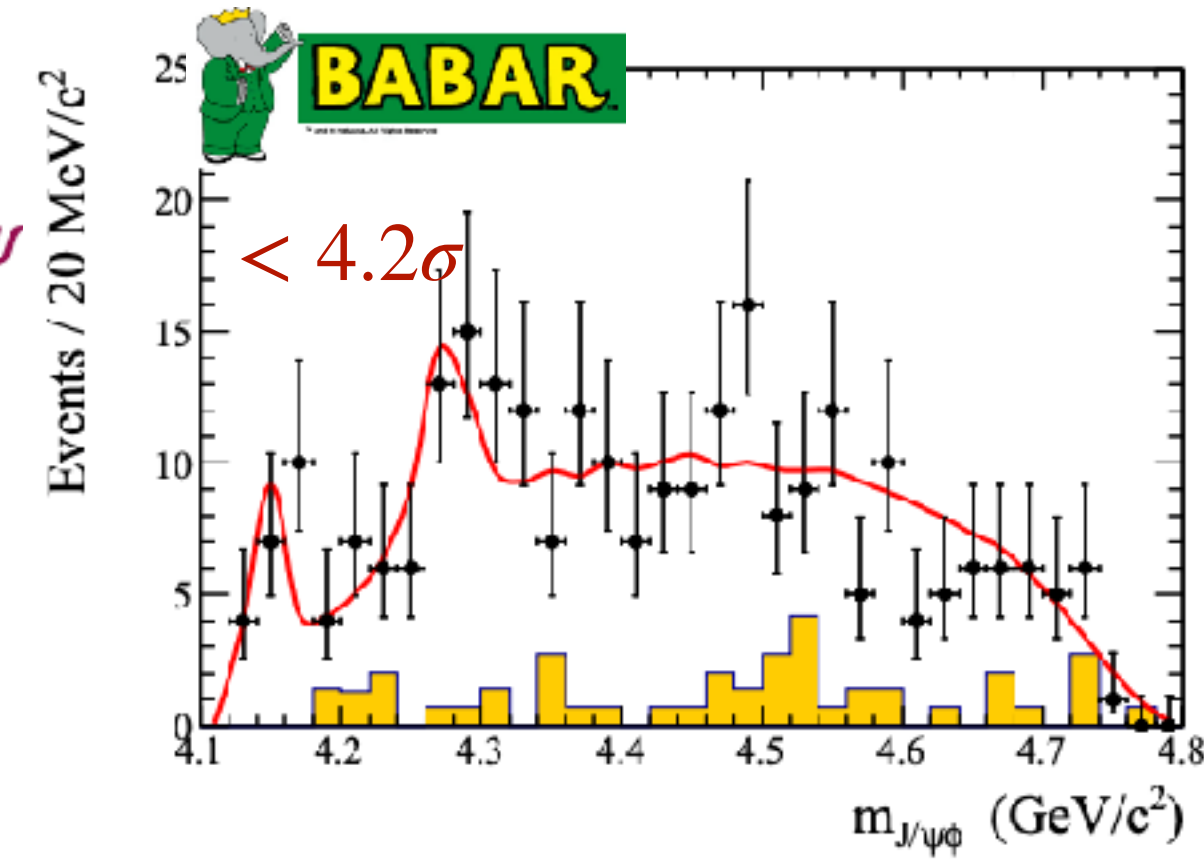
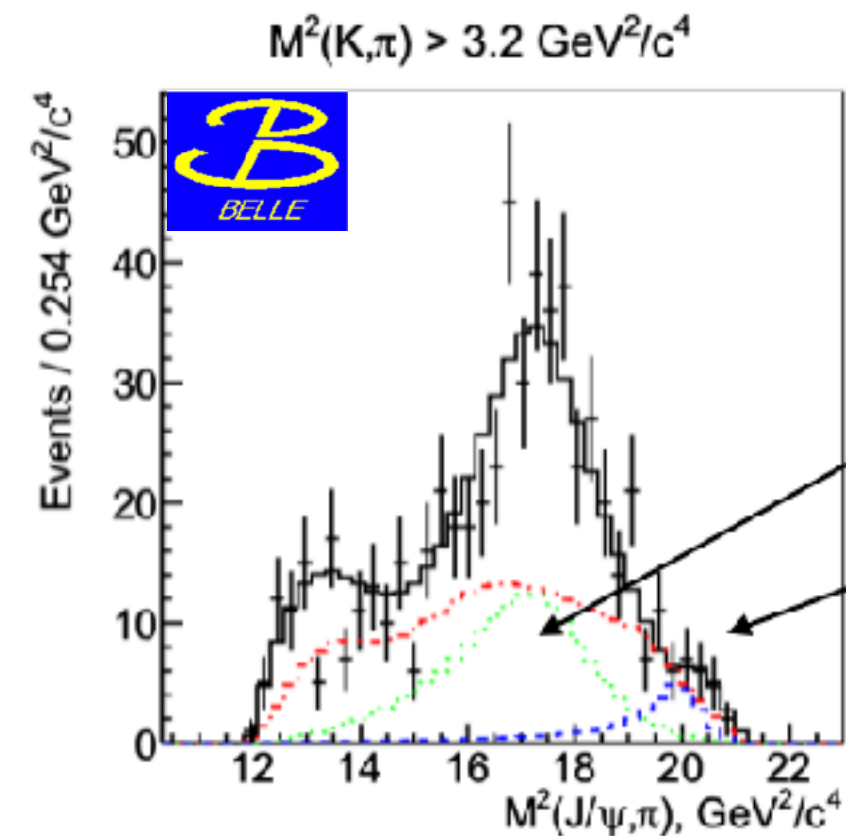
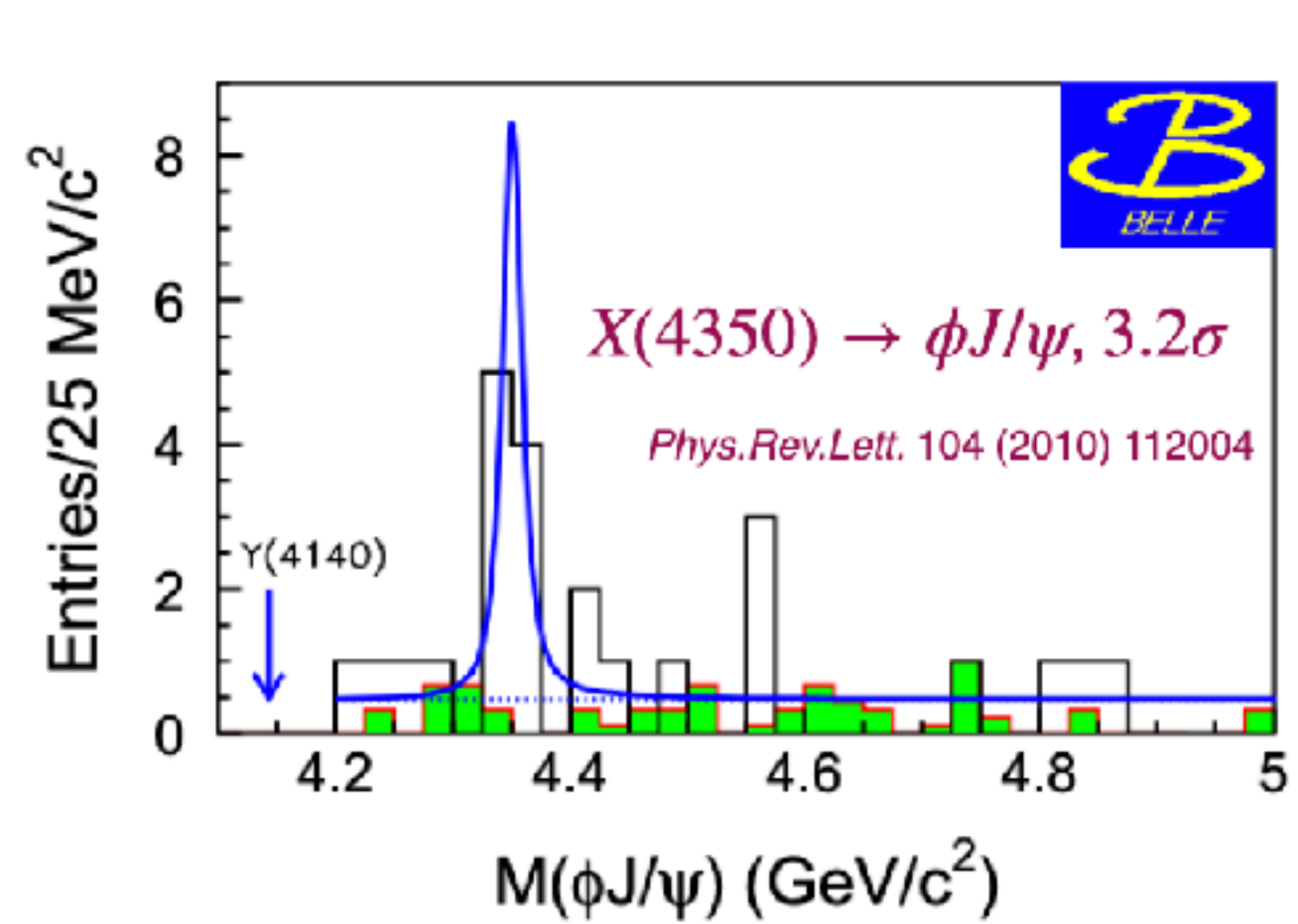
Benefits of hadron spectroscopy at B-factories

- ☑ Efficient reconstruction of neutrals (π^0 , η , ...)
- ☑ Fully reconstruction or recoil system
- ☑ Variety of production mechanisms
- ☑ Large production rate of $b \rightarrow c\bar{c}$
- ☑ Unique dataset

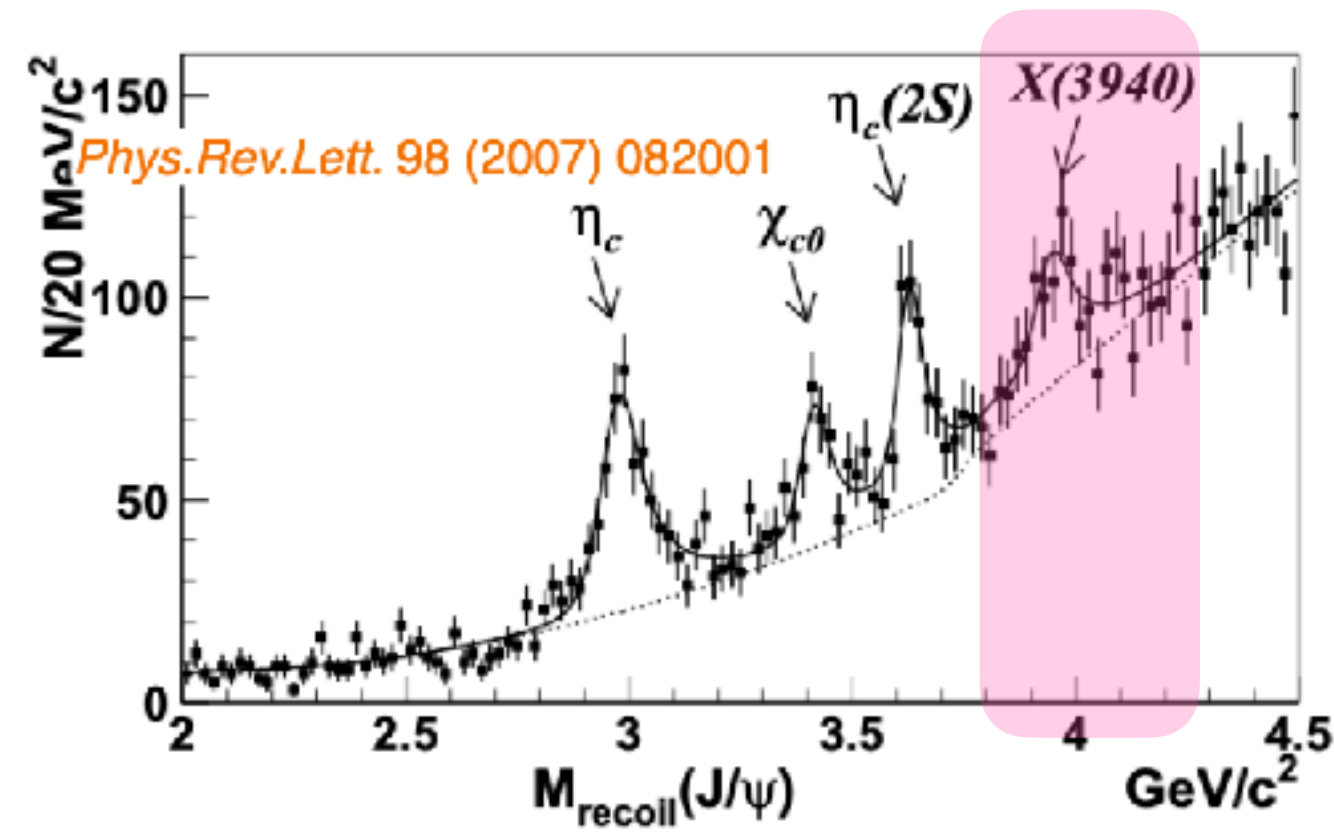


Prospect

Evidence could be clarified, e.g.



Properties measurements with dedicated analysis

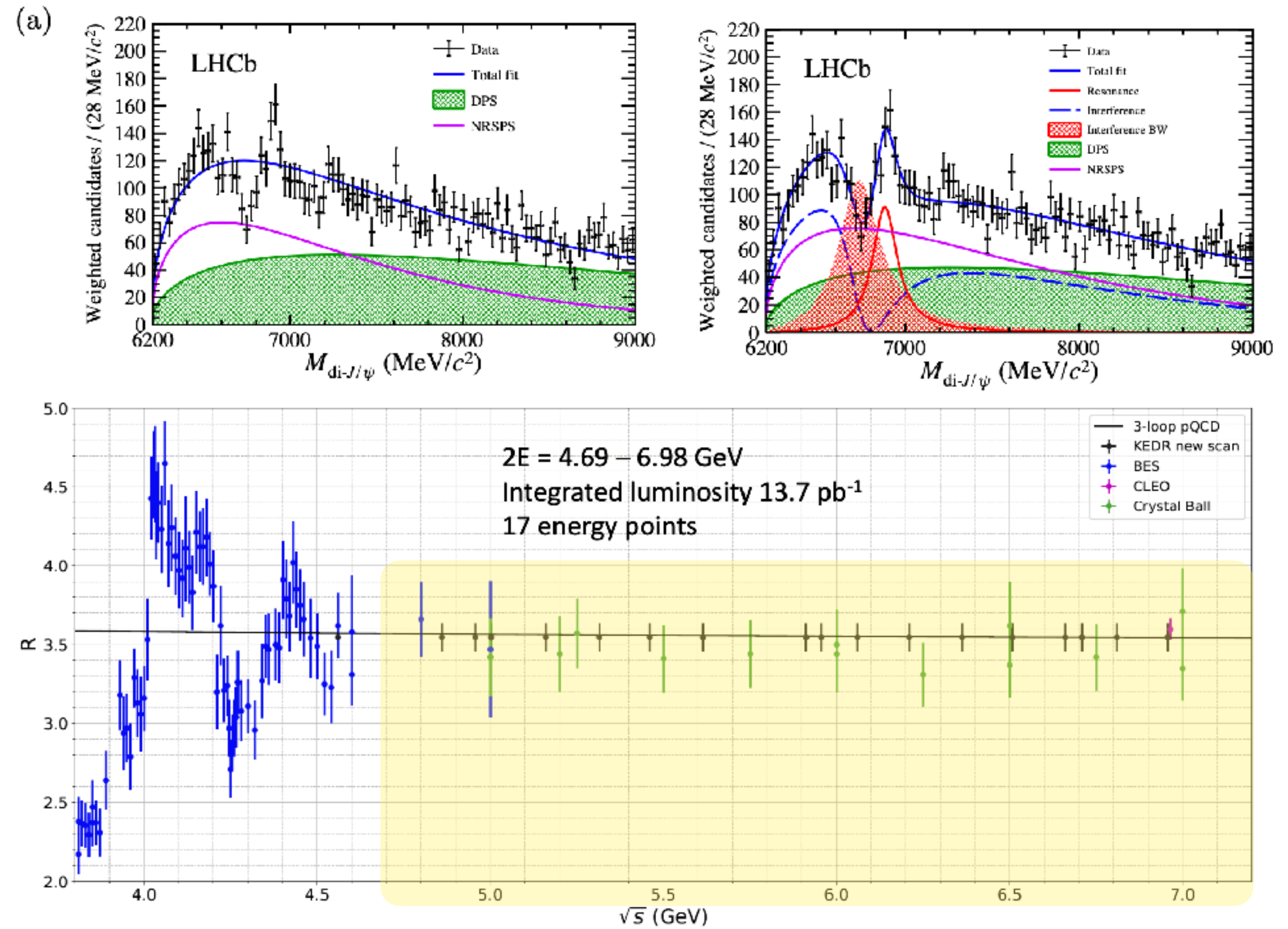
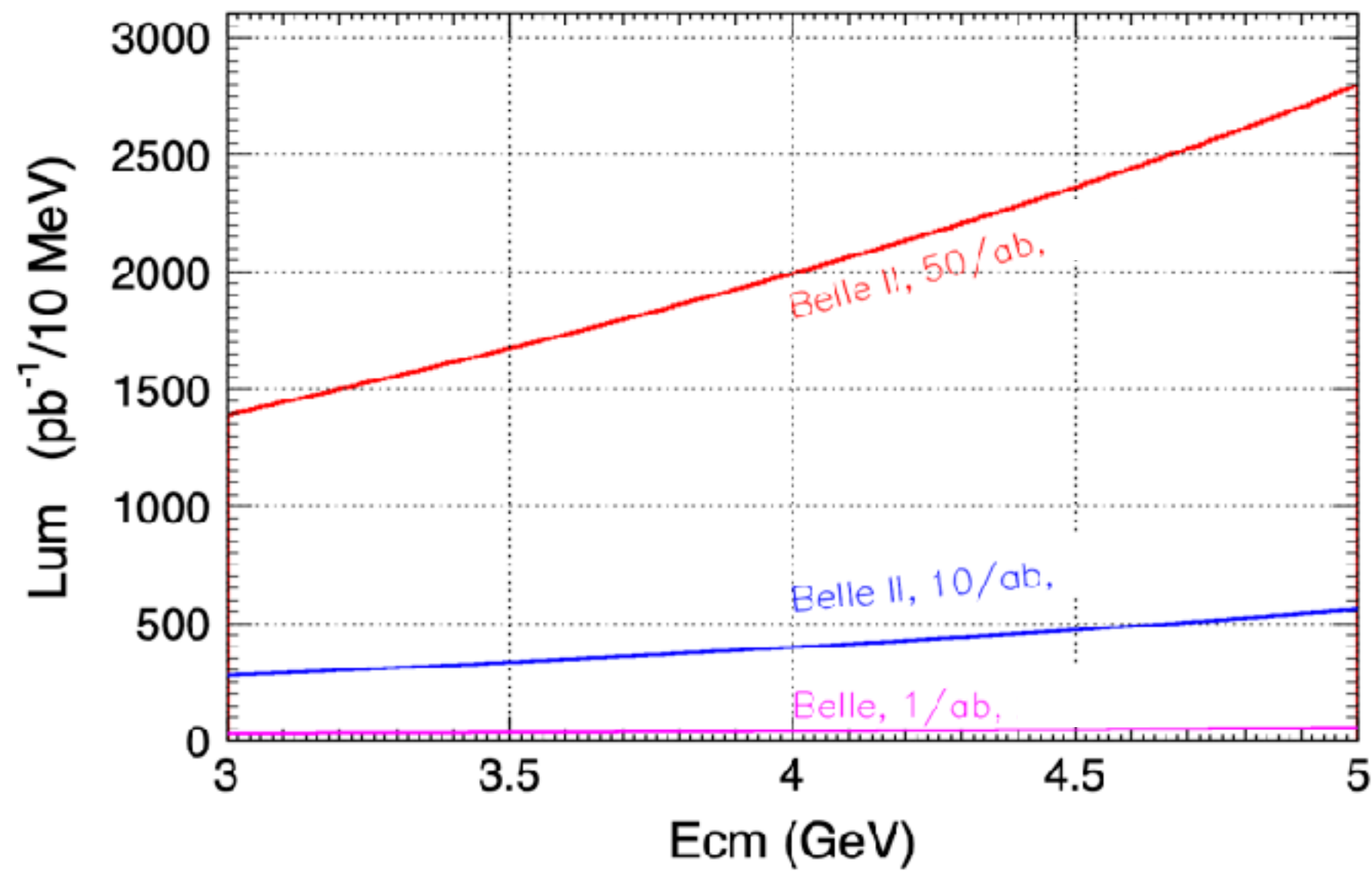


$X(3940) \quad I^G(J^{PC}) = ?^?(?^{??})$

Quantum number of $X(3940)$ remains unknown.

Prospect

Fully cover charmonium region with ISR



- Dedicated study to $e^+e^- \rightarrow \pi^+\pi^-J/\psi, K\bar{K}J/\psi$, etc.
- Z_c production in both e^+e^- annihilation and B decays.
- Doubly charmonium state in, e.g. $e^+e^- \rightarrow \eta_c J/\psi, \chi_c J/\psi$

Other prospects at Belle II

- ★ Very high statistics samples of $\Upsilon(4S)$
 - ★ Dedicate study of $X(3872)$ decays to final states with neutrals, i.e. $D^0\bar{D}^{*0}$.
 - ★ Searching for new charmonium(-like) states in various productions.
- ★ Higher statistics samples of $\Upsilon(5S)$ and $\Upsilon(6S)$
 - ★ Investigate Z_b states: quantum numbers, neutral partners, decay modes...
 - ★ Search for new states
 - ★ Potential laboratory for other bottomonium states like $h_b(3P)$, $\Upsilon(D)$
- ★ Potential to reach higher E_{cms}
- ★ Reach charmonium(-like) states via ISR with huge datasets

*