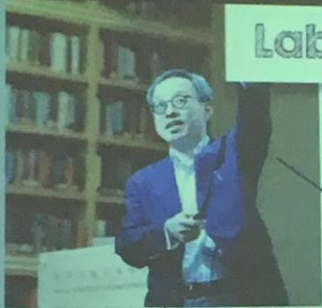


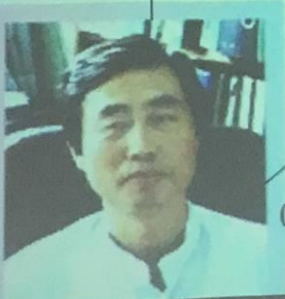
# **Yonsei university Cosmology and High Energy physics workshop (YuCHE 2019)**

## **Pictures Day2**

## Lab. for Dark Universe



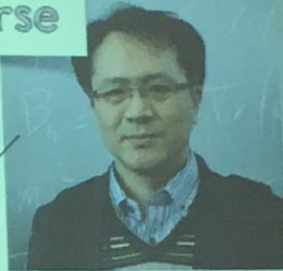
5 papers  
(평균인용수 10.6)



2 papers  
(평균인용수 228)



3 papers  
(평균인용수 37)

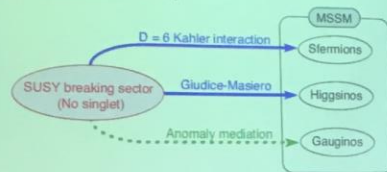


9 papers  
(평균인용수 40.6)



## Anomaly mediation + pure gravity mediation

[Giudice, Luty, Murayama & Rattazzi; Randall & Sundrum; Ibe, TM & Yanagida; Ibe & Yanagida; Arkani-Hamed et al.]

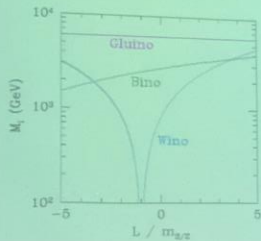


- The most general supergravity Lagrangian with Planck-suppressed interactions
- No singlet in SUSY breaking sector
- Scalar masses are of the order of the gravitino mass  $m_{3/2}$





### Gaugino masses in the PGM model



$$M_1 \simeq \frac{g_1^2}{16\pi^2} (11m_{3/2} + L)$$

$$M_2 \simeq \frac{g_2^2}{16\pi^2} (m_{3/2} + L)$$

$$M_3 \simeq \frac{g_3^2}{16\pi^2} (-3m_{3/2})$$

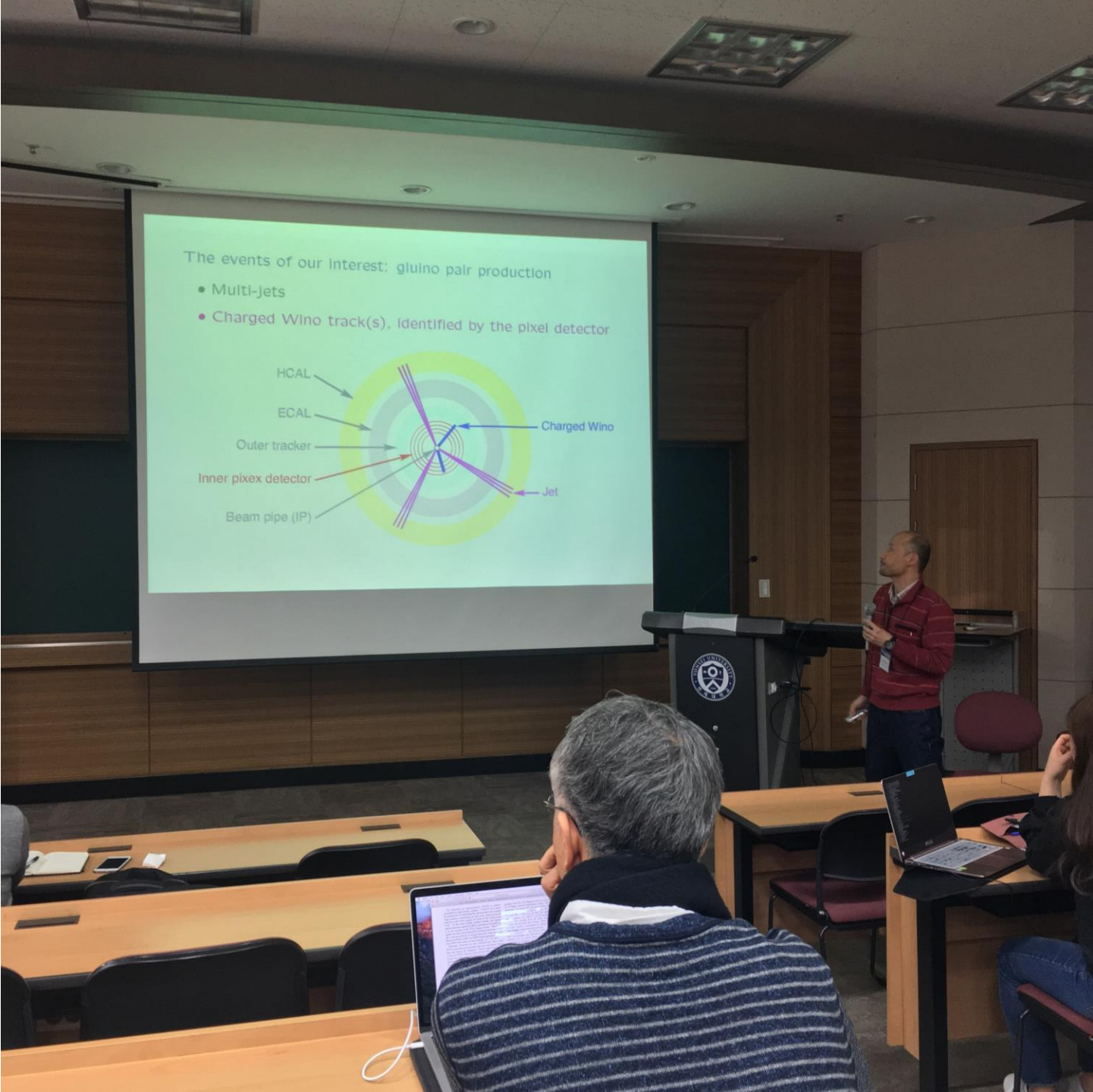
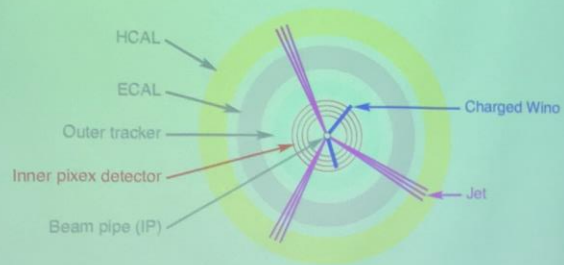
$$L \equiv \mu \sin 2\beta \frac{m_A^2}{\mu^2 - m_A^2} \ln \frac{\mu^2}{m_A^2}$$

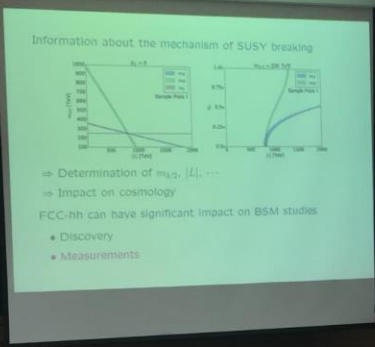
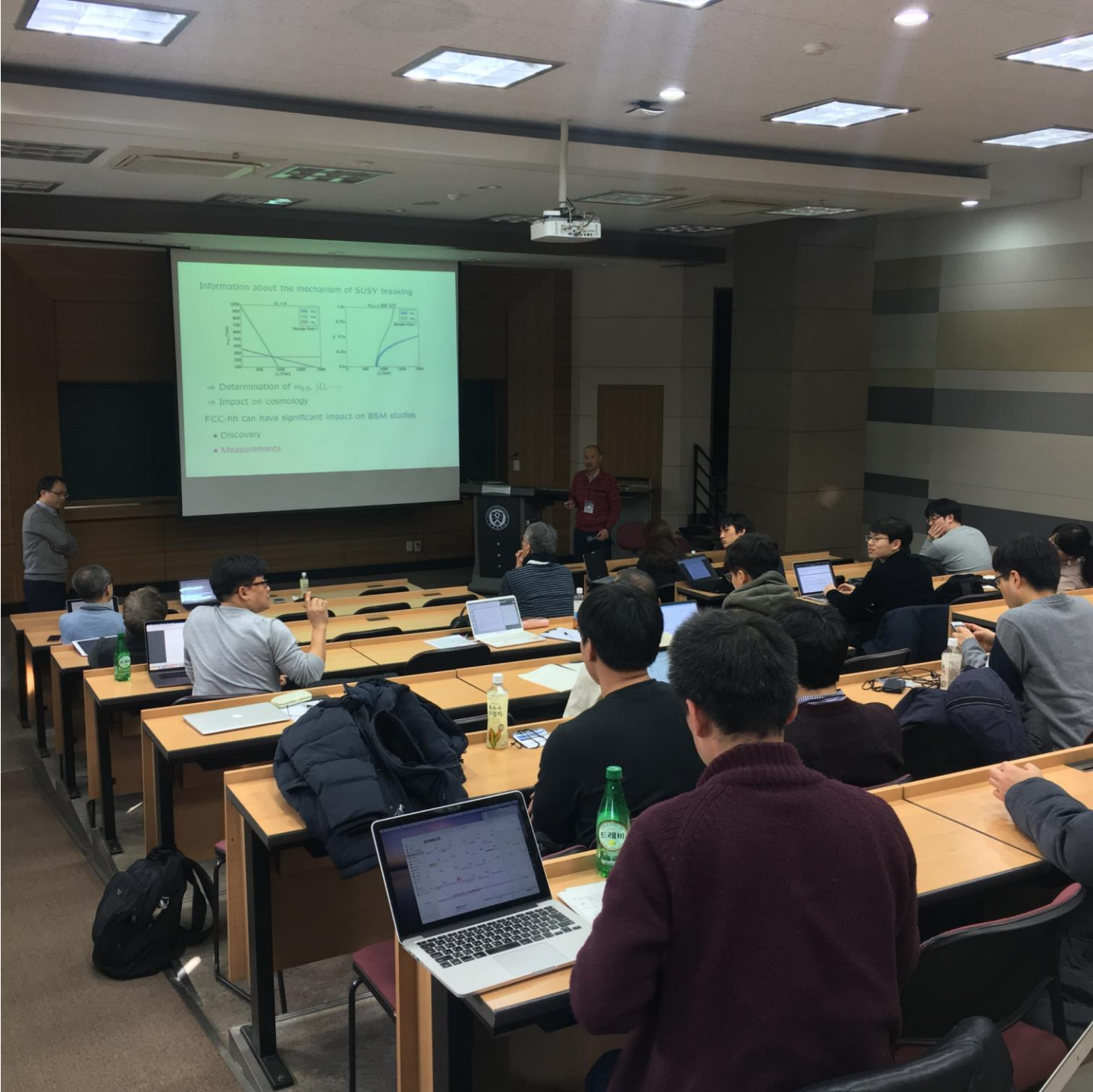
- Wino (gaugino for  $SU(2)_L$ ) is likely to be the LSP

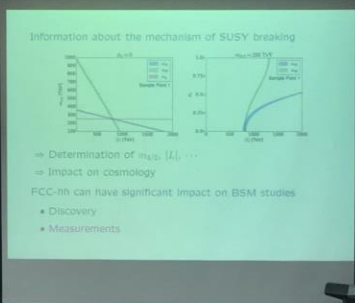
$$\left| \frac{10g_1^2}{3g_3^2} m_{\tilde{g}} - \frac{g_1^2}{g_2^2} m_{\tilde{W}} \right| \lesssim m_{\tilde{B}} \lesssim \frac{10g_1^2}{3g_3^2} m_{\tilde{g}} + \frac{g_1^2}{g_2^2} m_{\tilde{W}}$$

The events of our interest: gluino pair production

- Multi-Jets
- Charged Wino track(s), Identified by the pixel detector

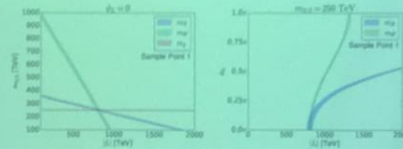








Information about the mechanism of SUSY breaking



- ⇒ Determination of  $m_{3/2}$ ,  $|L|$ , ...
- ⇒ Impact on cosmology

FCC-hh can have significant impact on BSM studies

- Discovery
- Measurements

## PQ mechanism

- Introduce a QCD anomalous global U(1) symmetry which is spontaneously broken at a high scale  $F_a \rightarrow$  axion.
- QCD anomaly induces an axion-gluon-gluon coupling:

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4}G_{\mu\nu}^a G^{a\mu\nu} + (\theta + c_G \frac{a}{F_a}) \frac{g_s^2}{32\pi^2} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} \Rightarrow \bar{\theta} \equiv \theta + c_G \frac{a}{F_a}$$

- QCD condensation generates axion potential driving the theta value to vanish:  $V \sim \Lambda_{\text{QCD}}^4 (1 - \cos \bar{\theta}) \Rightarrow \langle \bar{\theta} \rangle = 0$

- Axion mass is predicted to be  $m_a \sim \frac{\Lambda_{\text{QCD}}^2}{F_a} \sim 10^{-3} \left( \frac{10^{10} \text{GeV}}{F_a} \right) \text{eV}$

Version: 1/2022/23

© Chan "Elegant" Anson



## Axion-pion mixing

- For QCD axion, the mixing is negligible ( $f_p \ll f_a, m_a \ll m_\pi$ ).
- For ALP, it can lead to sizable contribution to  $K^+ \rightarrow \pi^+ a$  &  $a \rightarrow \gamma\gamma$  induced from  $K^+ \rightarrow \pi^+ \pi^0$  &  $\pi^0 \rightarrow \gamma\gamma$ :

$$\Gamma(K^+ \rightarrow \pi^+ a) \approx \left( c_\pi \frac{f_\pi}{f_a} \frac{m_a^2}{m_\pi^2 - m_a^2} \right)^2 \Gamma(K^+ \rightarrow \pi^+ \pi^0)$$

$$\Gamma(a \rightarrow \gamma\gamma) \approx \left( c_\pi \frac{f_\pi}{f_a} \frac{m_a^2}{m_\pi^2 - m_a^2} \right)^2 \left( \frac{m_a}{m_\pi} \right)^3 \Gamma(\pi^0 \rightarrow \gamma\gamma) \Rightarrow (\theta_{a\gamma})_{\text{mix}} = \frac{c_\pi m_a^2}{\pi m_\pi^2 - m_a^2} \frac{1}{f_a}$$

- $B(K^+ \rightarrow \pi^+ a) < 10^{-10}$  puts a limit:

$$f_a > 4 \left( \frac{c_\pi m_a^2}{m_\pi^2 - m_a^2} \right) \text{TeV} \Rightarrow (\theta_{a\gamma})_{\text{mix}} < 5.8 \times 10^{-7} \text{GeV}^{-1} \text{ for } m_a < 110 \text{MeV}$$

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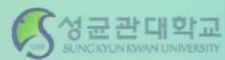




# Constraining neutrino-DM interaction with IceCube I80922A

[work on progress with Carsten Rott and Jongkuk Kim]

Ki -Young Choi



26 February, 2019 at Yonsei University

Ki-Young Choi, Sungkyunkwan University, Korea





## Summary of part 2

A rapid cooling of Cas A Neutron Star (NS) has been observed.

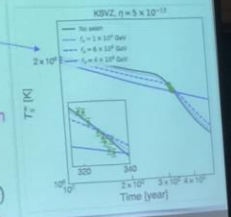
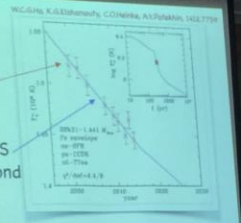
It can be explained within the standard NS cooling scenario. (i.e., without physics beyond the Standard Model)



If there is an extra cooling, such as an axion emission, the cooling is modified.

We studied the Cas A NS cooling with an axion emission, and obtained a new bound on the axion decay constant,  $f_a > O(10^8)$  GeV. (comparable to the existing SNI987A bound)

K.Hamaguchi, N.Nagata, K.Yanagi, J.Zheng, 1806.07151



## Cas A NS Cooling with axion

$$C \frac{dT}{dt} = -L_\nu - L_a$$

axion emission

### What we did:

- followed NS cooling **with axion emission** (Brems. and PBF), by modifying a public code NSCool.
- APR EoS.
- NS mass  $M = 1.4 M_{\text{sun}}$ .
- gap models:
  - ▶  $n-1S_0$  gap: SFB (doesn't matter)
  - ▶  $p-1S_0$  gap: CCDK (doesn't matter as far as large enough)
  - ▶  $n-3P_2$  gap: **gap height  $\Delta \propto T_c$  and width: free parameter.**



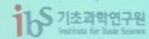




# ENERGETIC ALPS FROM DECAYING DARK MATTER

Kyu Jung Bae

Center for Theoretical Physics of the Universe,



based on

arXiv: 1806.08569

with A. Kamada (IBS-CTPU), H.-J. Kim (KAIST)

YuCHE 2019 @Yonsei U.

Feb. 26, 2019





# Unitary inflaton as dark matter and radiation

Hyun Min Lee  
Chung-Ang University, Korea



Based on: HML, Phys. Rev. D78, 015020 (2018);  
S. Choi, Y. Kang, HML, K. Yamashita, arXiv: 1902.03781 [hep-ph].

Cosmology and High Energy Physics Workshop  
Yonsei University, Feb 26, 2019.



# Reheating and dark matter

Reheating temperature:

$$BR = \frac{\Gamma_{\chi \rightarrow \gamma\gamma}}{\Gamma_{\chi \rightarrow \gamma\gamma} + \Gamma_{\chi \rightarrow bb}} = \frac{11.5\lambda_\chi^2}{11.5\lambda_\chi^2 + \lambda_{H\chi}^2}$$

$$\Gamma_{\chi\gamma} = \Gamma_{\chi \rightarrow b\bar{b}} \cdot \left(\frac{1}{1-BR}\right) \approx H_{dec} = \sqrt{\frac{\lambda_\chi \chi_{eff}(t_{dec})}{12 M_P}}; \quad \frac{\pi^2 g_*(T_{RH})}{30} T_{RH}^4 = (1-BR) \cdot \rho_\chi(t_{dec})$$

$$\begin{aligned} \rightarrow T_{RH} &= 0.002 \left(\frac{100}{g_*(T_{RH})}\right)^{1/4} \lambda_{H\chi}^2 \lambda_\chi^{-3/4} (1-BR)^{-3/4} M_P \\ &= (4.4 \times 10^6 \text{ GeV}) \left(\frac{100}{g_*(T_{RH})}\right)^{1/4} \left(\frac{\lambda_{H\chi}}{10^{-4}}\right)^2 R^2 (1-BR)^{-3/4} \left(\frac{r}{0.01}\right)^{-3/4} \end{aligned}$$

Dark matter abundance:  $\lambda_{\chi H} \sim \lambda_{\gamma H} / \xi_1^2 \lesssim 10^{-7}$  out of equilibrium

$$\Omega_\chi h^2 = (\Omega_\chi h^2)_{\text{FIMP}} + (\Omega_\chi h^2)_{\text{RH}} \quad \text{: non-thermal production}$$

$$\text{Higgs decay: } (\Omega_\chi h^2)_{\text{FIMP}} = 0.12 \left(\frac{100}{g_*(m_h)}\right)^{3/2} \left(\frac{\lambda_{H\chi}}{4.4 \times 10^{-7}}\right)^2 \left(\frac{m_\chi}{1 \text{ eV}}\right)$$

$$\text{Inflaton decay: } (\Omega_\chi h^2)_{\text{RH}} = 7.3 R \left(\frac{r}{0.01}\right)^{-1/4} \cdot BR \cdot \left(\frac{m_\chi}{1 \text{ eV}}\right)$$

# Dark Matter and WIMP Baryogenesis in Scotogenic Model

Yonsei university Cosmology and High Energy physics  
workshop (YuCHE 2019)  
2019.02.26-27

Sin Kyu Kang

Seoul-Tech



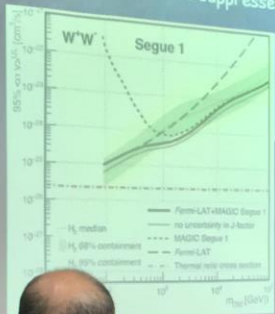
Based on arXiv:1806.04689 (to appear in JHEP)  
with D Borah, A Dasgupta



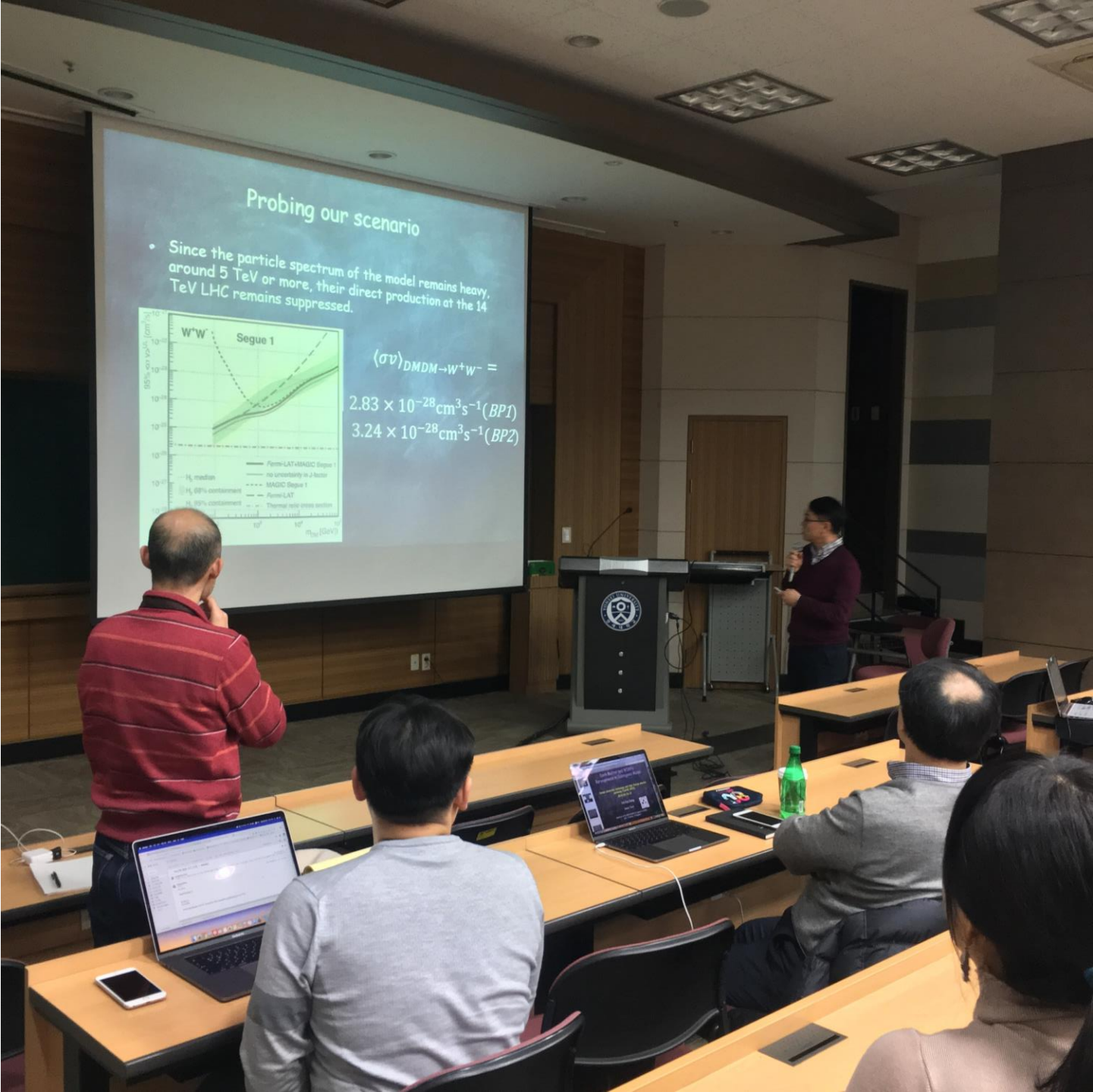


## Probing our scenario

- Since the particle spectrum of the model remains heavy, around 5 TeV or more, their direct production at the 14 TeV LHC remains suppressed.



$$\langle\sigma v\rangle_{DMDM\rightarrow W^+W^-} =$$
$$2.83 \times 10^{-28} \text{ cm}^3 \text{ s}^{-1} (BP1)$$
$$3.24 \times 10^{-28} \text{ cm}^3 \text{ s}^{-1} (BP2)$$



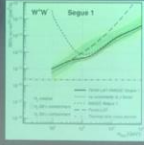






### Probing our scenario

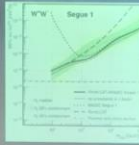
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$$2.83 \times 10^{-28} \text{cm}^2 \text{s}^{-1} (BP1)$$

$$3.24 \times 10^{-29} \text{cm}^2 \text{s}^{-1} (BP2)$$

Yonsei university Cosmology and High Energy physics  
workshop (YuCHE 2019)

# Murdering the 4th generation 2HDM

Jeonghyeon Song  
(Konkuk University, Korea)

w/ S.K.Kang, Z. Qian, Y.W. Yoon  
Yonsei Univ, 2019.2.26






Introduction  
A model for light DM  
Direct detection of light dark baryons

## WIMP miracle for correct thermal relics

- ▶ Weakly interacting massive particles (Lee+Weinberg, 1977)



DM SM  
DM SM  
annihilation

$$\langle \sigma v \rangle = \frac{\alpha^2}{m_{DM}^2}$$
$$m_{DM} \sim 100 \text{ GeV}$$

5/25









