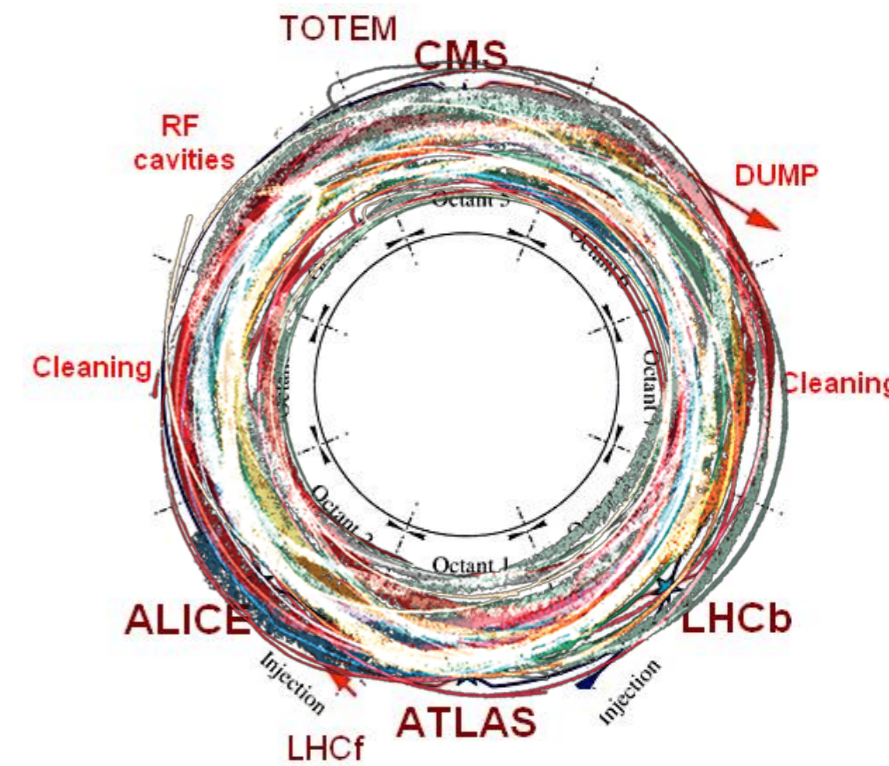


Precision Higgs at



Myeonghun Park



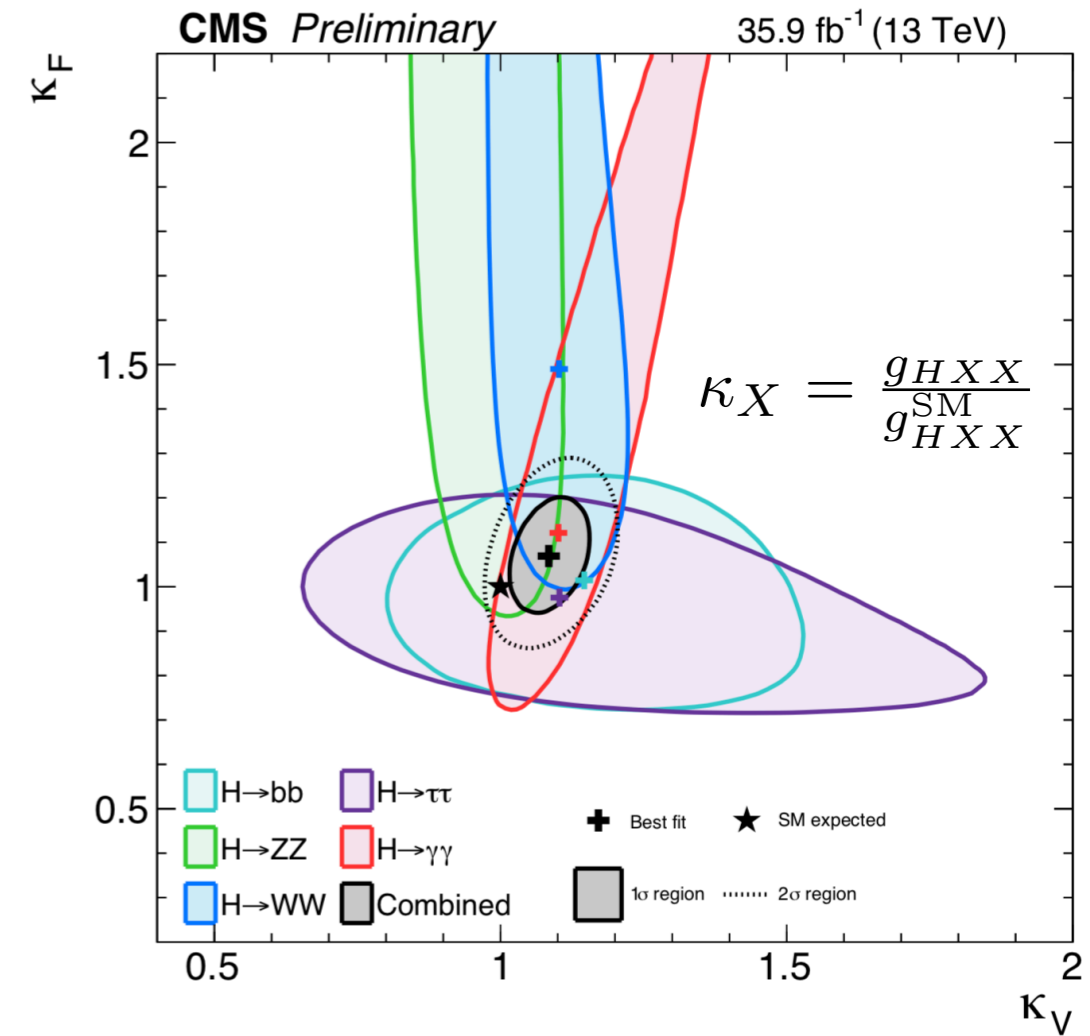
based on arXiv:1812.02679
with **Zhuoni Qian** and SJ Lee

based on arXiv:1807.11498 (PRL 2019), arXiv:1903(4?).XXXX
with K. Matchev, KC Kong, **Jeong Han Kim** and **Minho Kim**,

YuCHE 2019

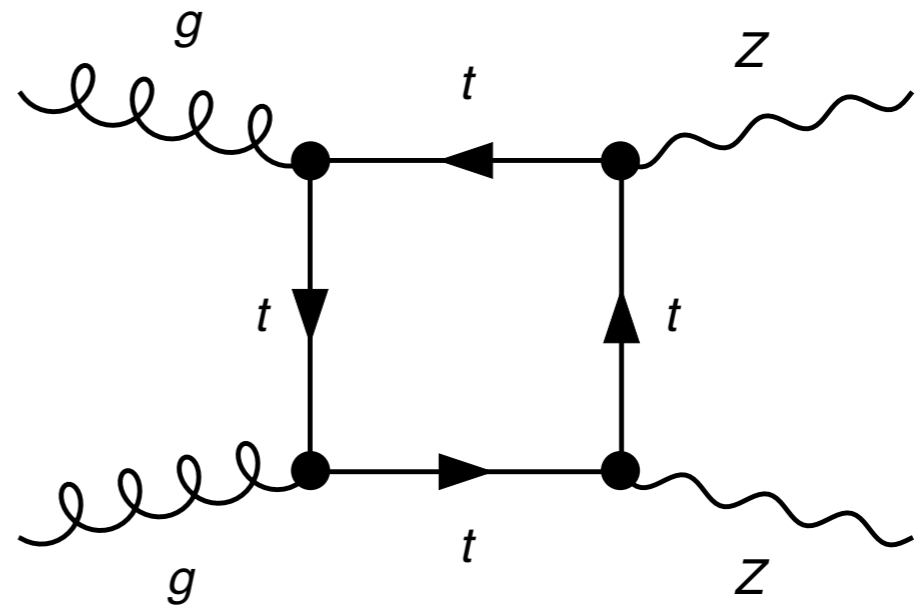
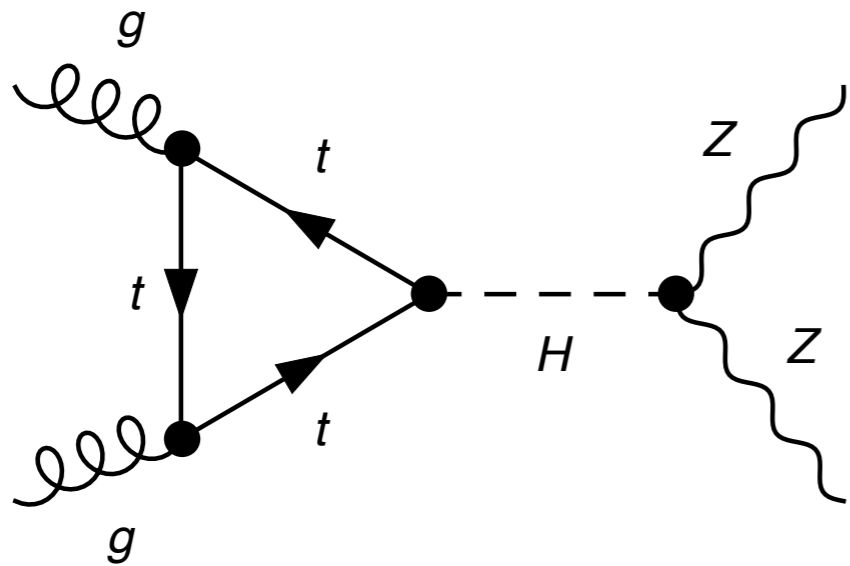
Precision for Higgs

Model	κ_V	κ_b	κ_γ
Singlet Mixing	$\sim 6\%$	$\sim 6\%$	$\sim 6\%$
2HDM	$\sim 1\%$	$\sim 10\%$	$\sim 1\%$
Decoupling MSSM	$\sim -0.0013\%$	$\sim 1.6\%$	$\sim -0.4\%$
Composite	$\sim -3\%$	$\sim -(3 - 9)\%$	$\sim -9\%$
Top Partner	$\sim -2\%$	$\sim -2\%$	$\sim +1\%$



- A singlet S, (for example Higgs portal Dark Matter Model)
- Two Higgs Doublet Model : Enlarged Higgs family
- MSSM, a decoupling limit for SM-like Higgs. Suppressed stop effect for a loop.
- Composite Higgs Model: Higgs as a pseudo-NGB, just like a pion.
- Top Partner, (for Naturalness...)

Precision for Higgs *beyond the Onshell*



$$\mathcal{A}_\Delta \rightarrow \frac{m_t^2}{\hat{s}} \frac{1}{2} \log^2 \left(\frac{m_t^2}{\hat{s}} \right) \left(\frac{\sqrt{\hat{s}}}{m_Z} \right)^2$$

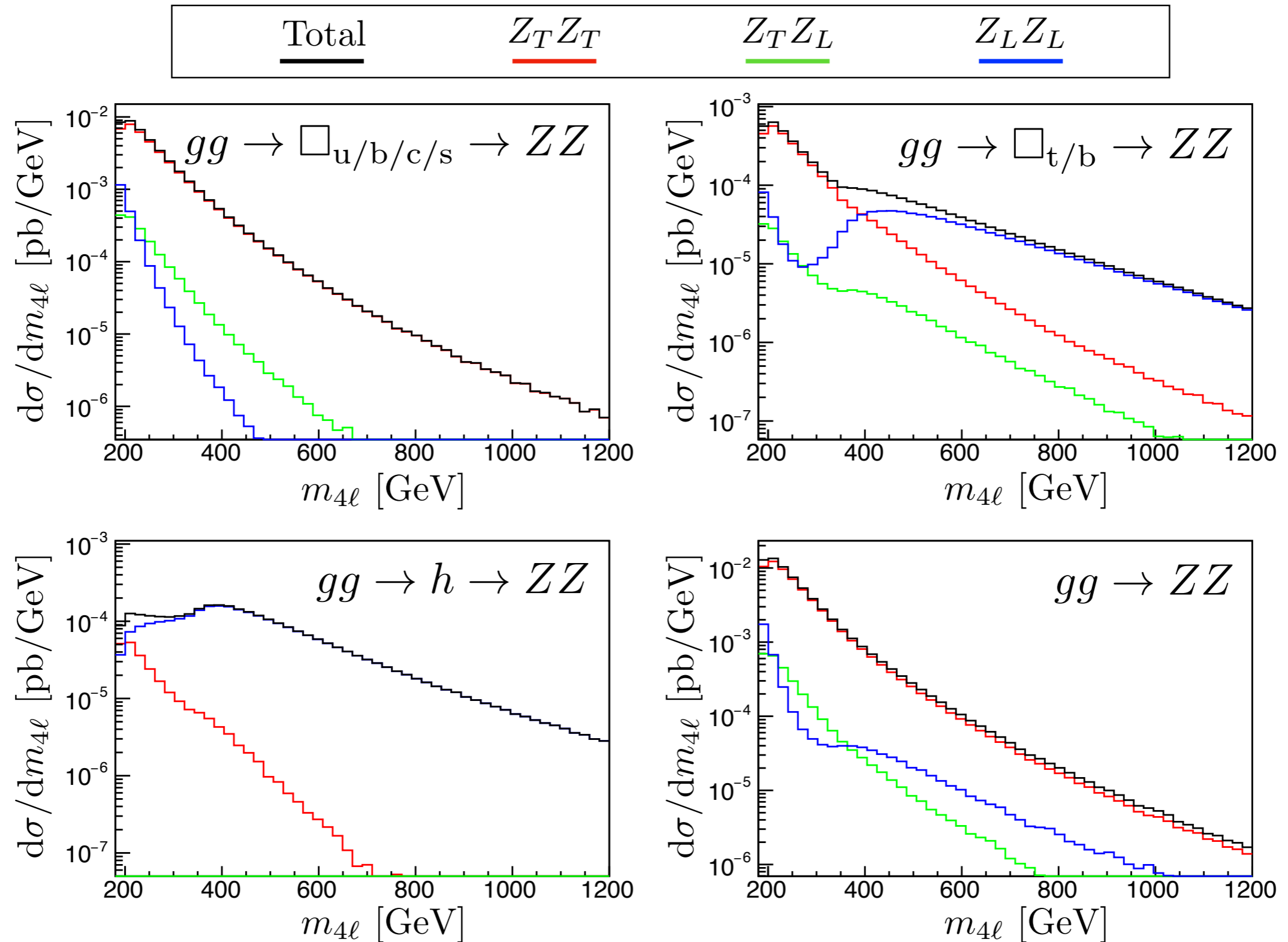
$$= \frac{m_t^2}{2m_Z^2} \log^2 \left(\frac{\hat{s}}{m_t^2} \right)$$

$$\mathcal{A}_\square \rightarrow -8C_A^2 \frac{m_t^2}{\hat{s}} \left(\frac{\hat{s}}{m_Z^2} \right) \log^2 \left(\frac{\hat{s}}{m_t^2} \right)$$

$$= -\frac{m_t^2}{2m_Z^2} \log^2 \left(\frac{\hat{s}}{m_t^2} \right)$$

in the limit $\sqrt{\hat{s}} \gg m_t$

Precision for Higgs *beyond the Onshell*



A new physics which ...

- **A:** A Higgs portal light scalar

$$\mathcal{L} \ni \mathcal{L}_{\text{SM}} + \partial_\mu S \partial^\mu S^* - \mu^2 |S|^2 - \kappa |S|^2 |\Phi|^2$$

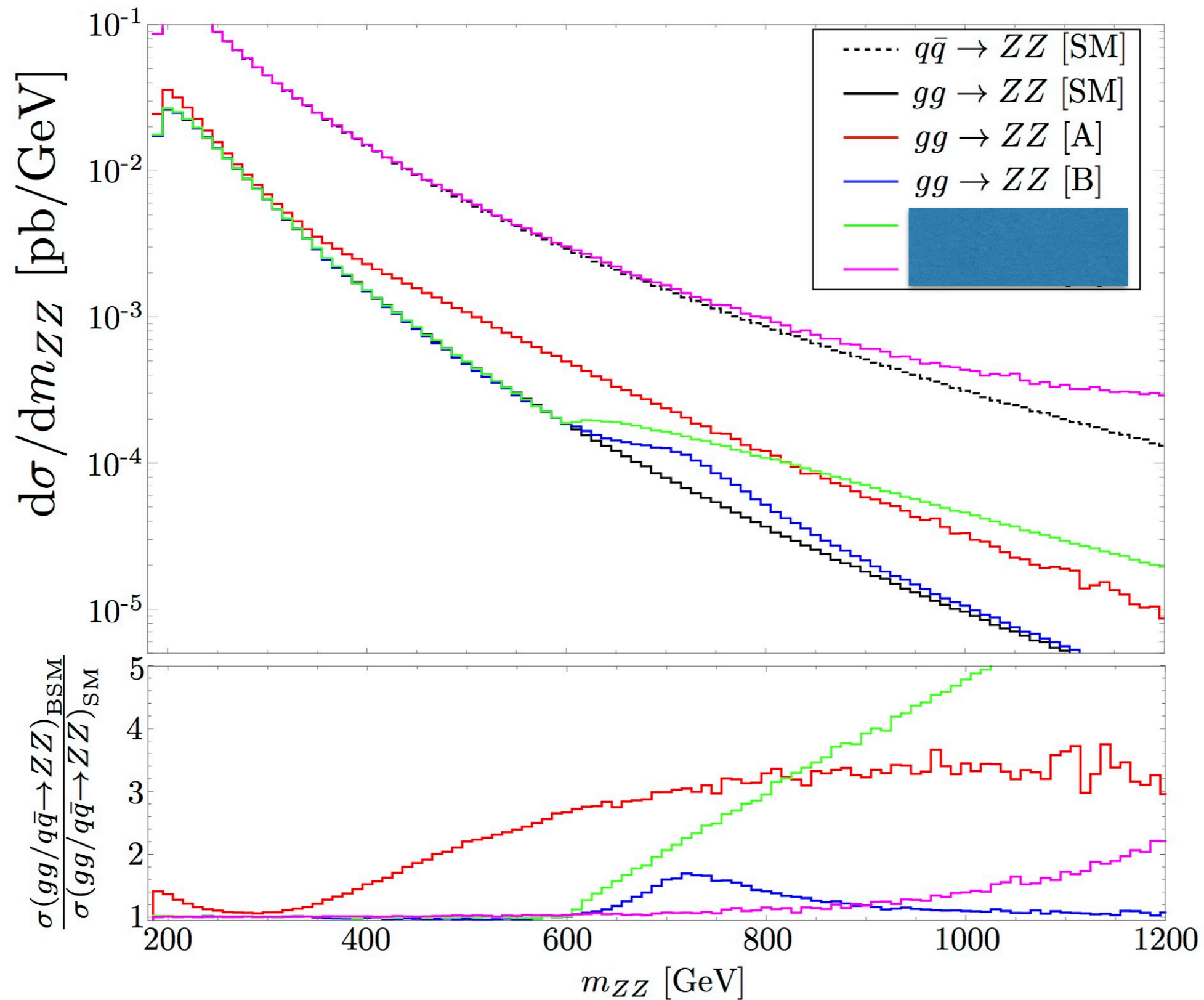
$$m_S = 80 \text{ GeV with } \kappa = 9$$

- **B:** Broad-width heavy scalar

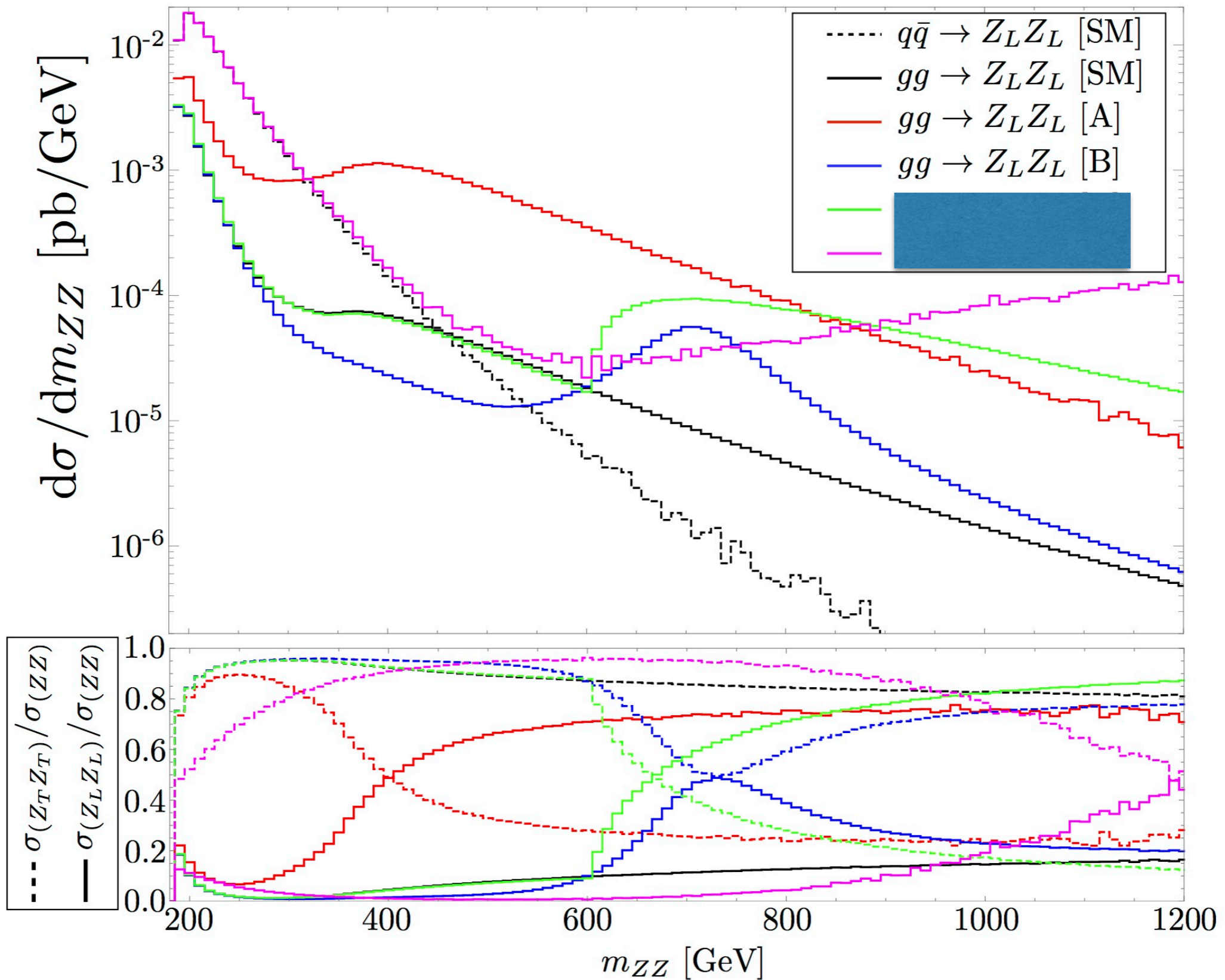
$$\mathcal{L} \ni \mathcal{L}_{\text{SM}} - \mu_S S |\Phi|^2$$

$$M_S = 700 \text{ GeV with } \Gamma_S = 20\%$$

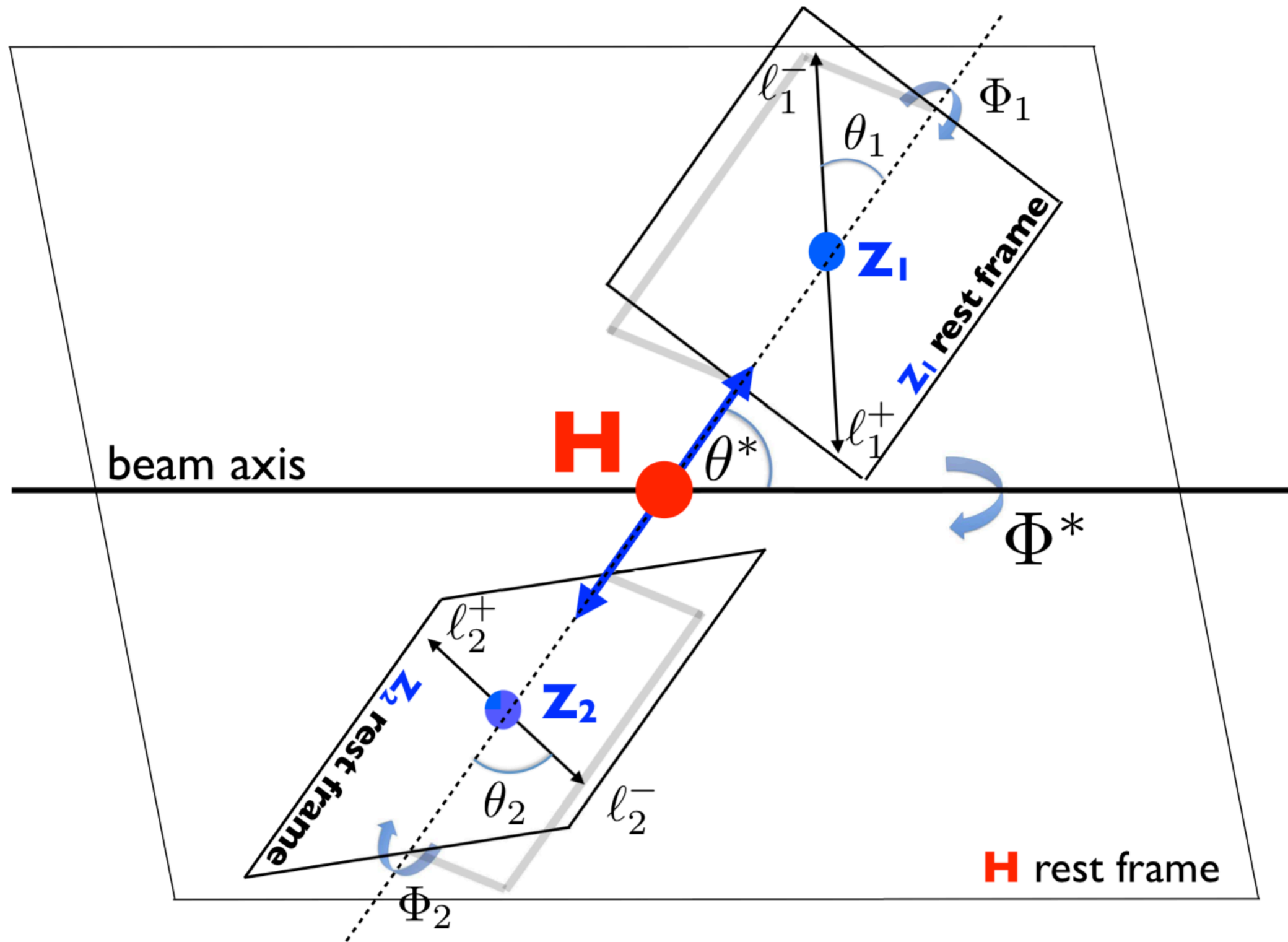
Precision for Higgs *beyond the Onshell*



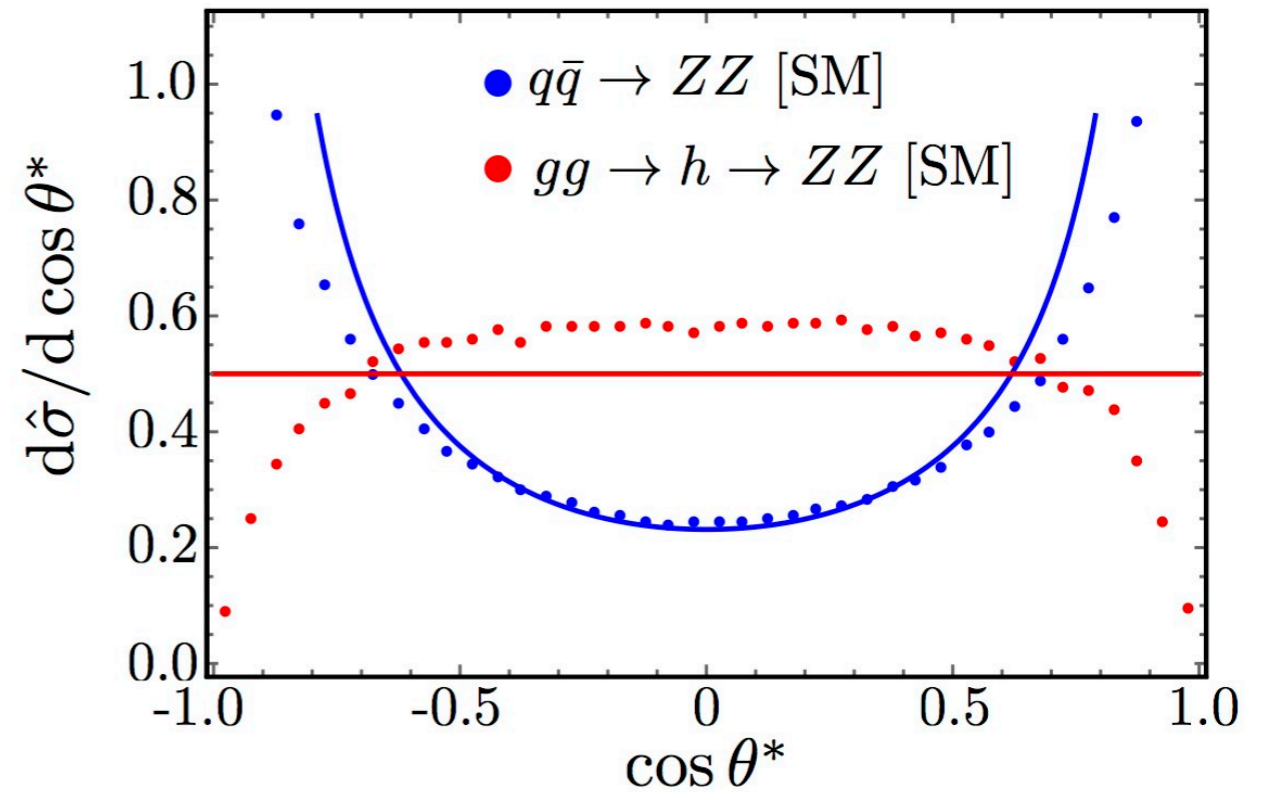
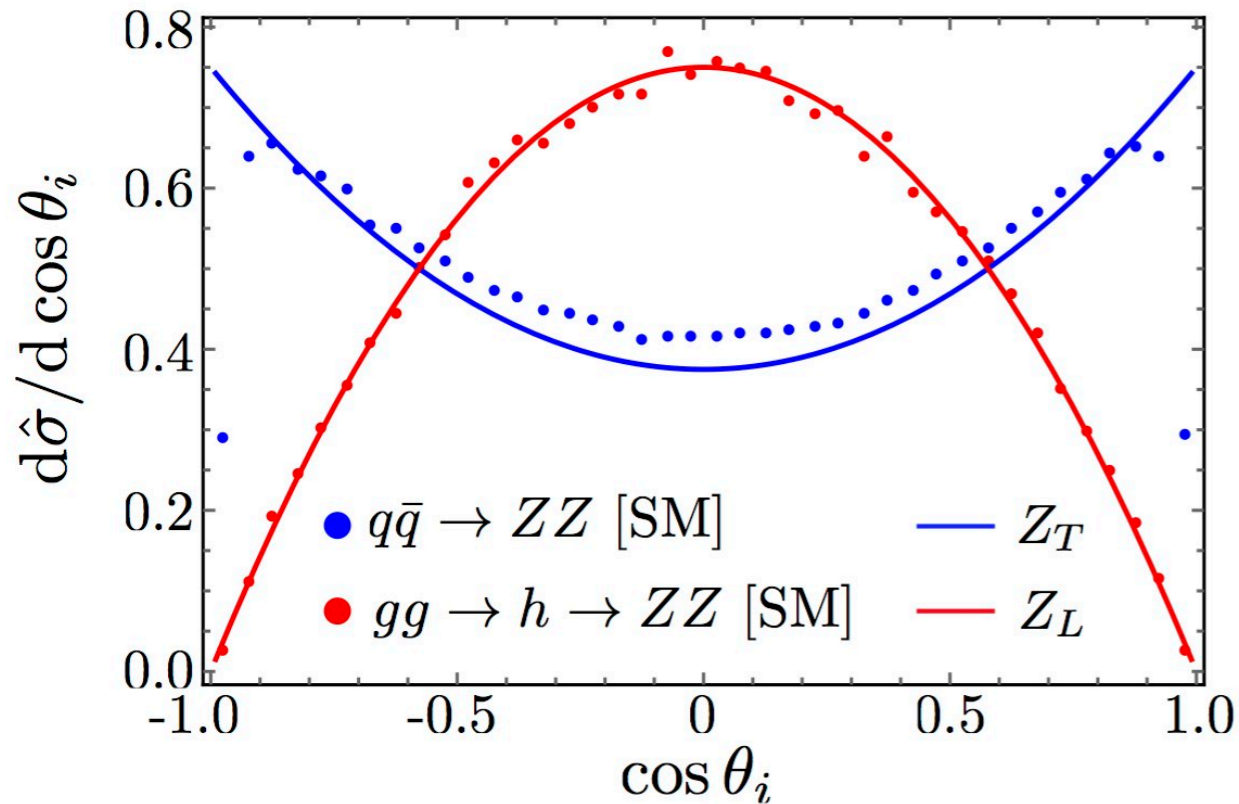
Precision for Higgs *beyond the Onshell*



Simple variables to utilize polarization



Simple variables to utilize polarization



Simple angle cuts: $|\cos \theta^*| < 0.7$, $|\cos \theta_1| < 0.68$

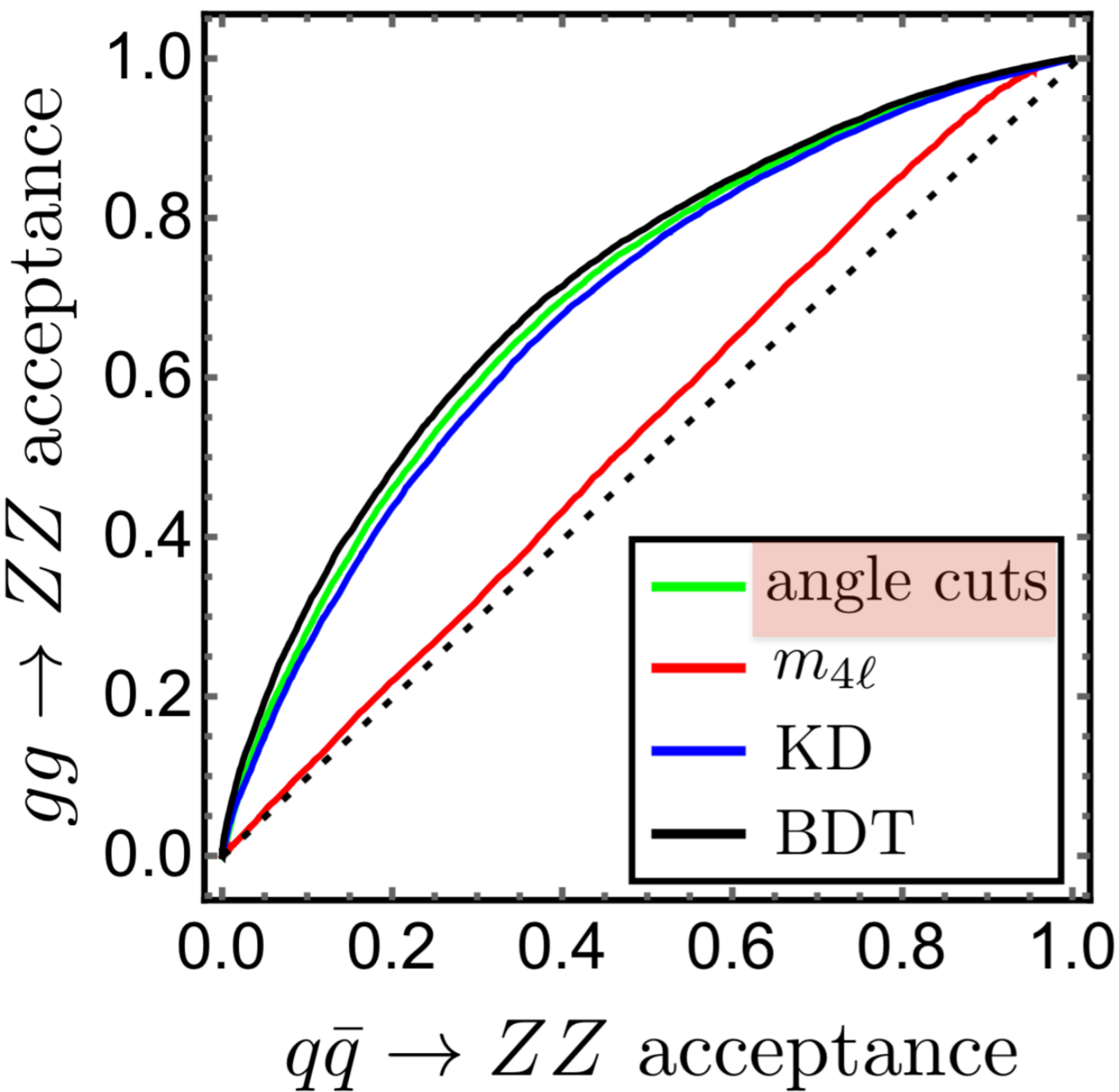
Conventional $m_{4\ell}$ cut: $m_{4\ell} > 700\text{GeV}$

Matrix Element Method: $\text{KD} = \ln \left(\frac{f_g(x_1)f_g(x_2)|\mathcal{M}(gg \rightarrow h^* \rightarrow 4\ell)|^2}{\sum f_q(x_1)f_{\bar{q}}(x_2)|\mathcal{M}(q\bar{q} \rightarrow 4\ell)|^2} \right)$

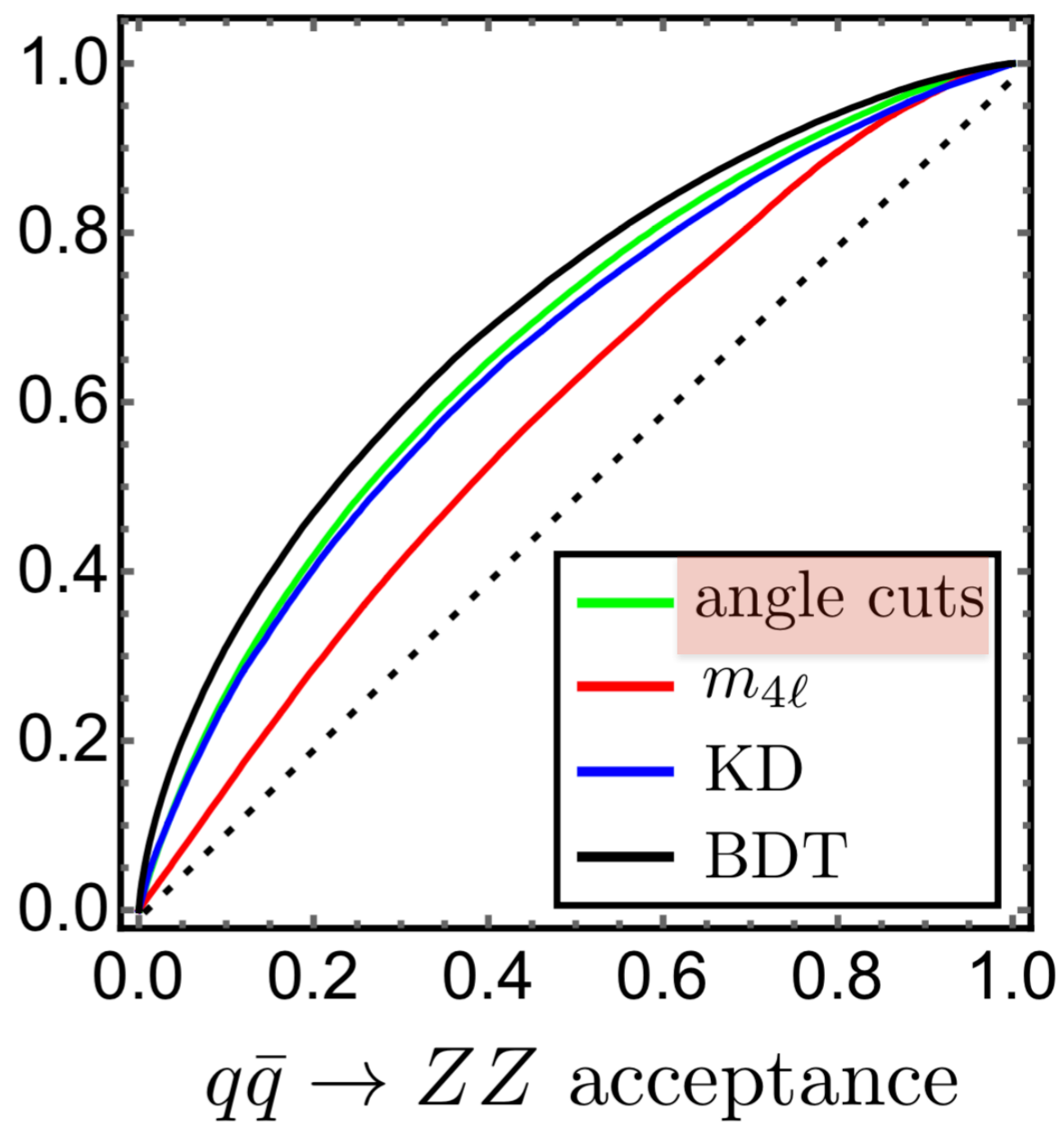
we set $m_{h^*} = m_{4\ell}$ for a generic case

Simple variables to utilize polarization

case A



case B



The Higgs Sector

"Mexican Hat" Higgs potential



Spontaneous Symmetry Breaking

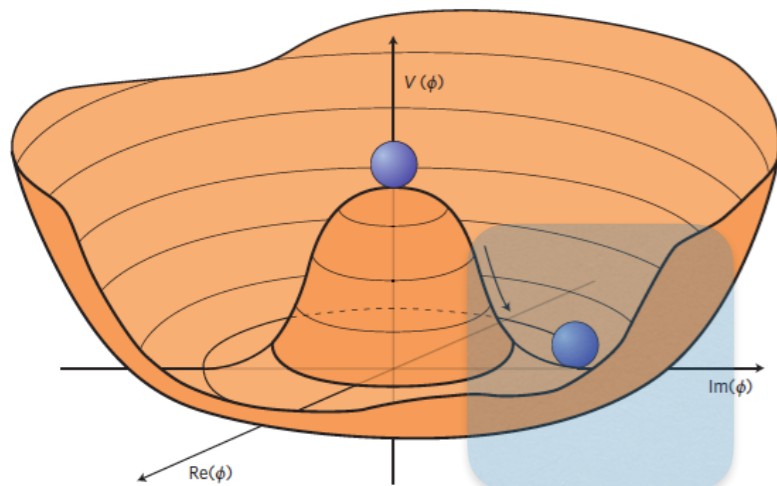
$$SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$$



Yukawa terms for massive fermions

Measuring Higgs potential

$$V(\phi) = -\mu^2\phi^2 + \lambda\phi^4$$



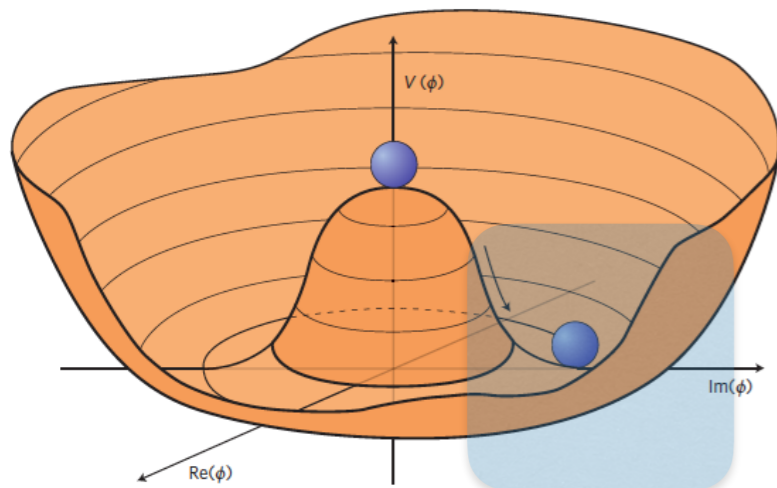
- From spontaneous symmetry breaking

$$V_h^{\text{SM}} = \frac{m_h^2}{2} h^2 + \frac{m_h^2}{2v^2} v h^3 + \frac{1}{4} \frac{m_h^2}{2v^2} h^4$$

- The shape of a potential around EW scale is determined by h^2 term and h^3 term.

Measuring Higgs potential

$$V(\phi) = -\mu^2\phi^2 + \lambda\phi^4$$



- From spontaneous symmetry breaking

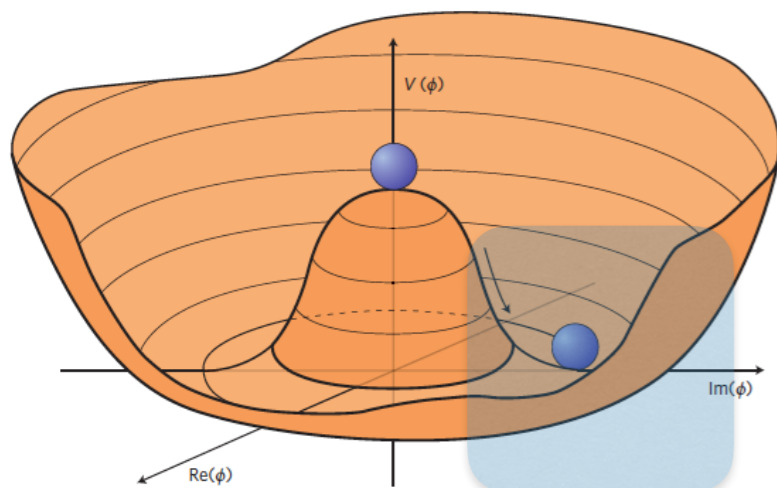
$$V_h^{\text{SM}} = \frac{m_h^2}{2} h^2 + \frac{m_h^2}{2v^2} v h^3 + \frac{1}{4} \frac{m_h^2}{2v^2} h^4$$

- The shape of a potential around EW scale is determined by h^2 term and h^3 term.
 - We know about the mass of Higgs from clean channels

Higgs Boson Pair

to measure Higgs potential

$$V(\phi) = -\mu^2\phi^2 + \lambda\phi^4$$



- From spontaneous symmetry breaking

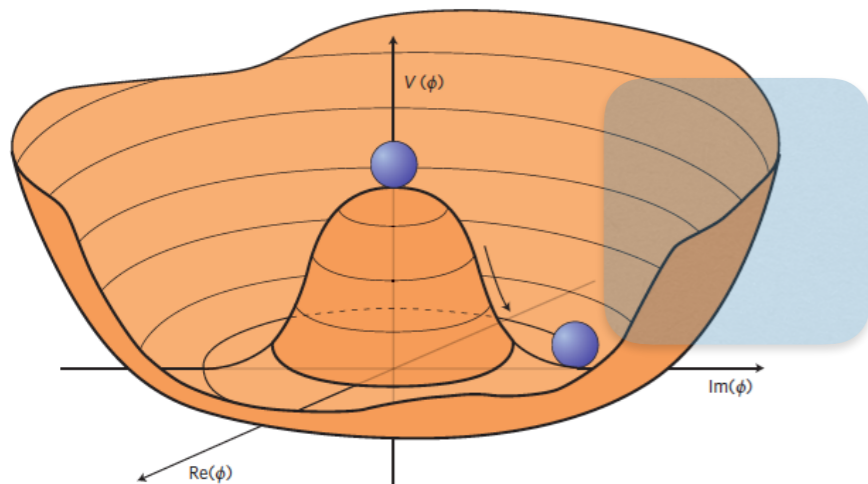
$$V_h^{\text{SM}} = \frac{m_h^2}{2} h^2 + \frac{m_h^2}{2v^2} v h^3 + \frac{1}{4} \frac{m_h^2}{2v^2} h^4$$

- The shape of a potential around EW scale is determined by h^2 term and h^3 term.

- This triple Higgs coupling can be determined by $h \rightarrow h, h$

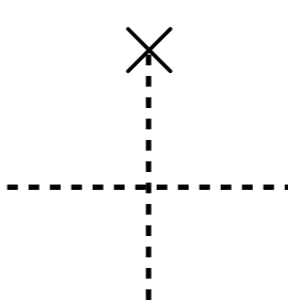
Multi-Higgs to measure Higgs potential

$$V(\phi) = -\mu^2\phi^2 + \lambda\phi^4$$



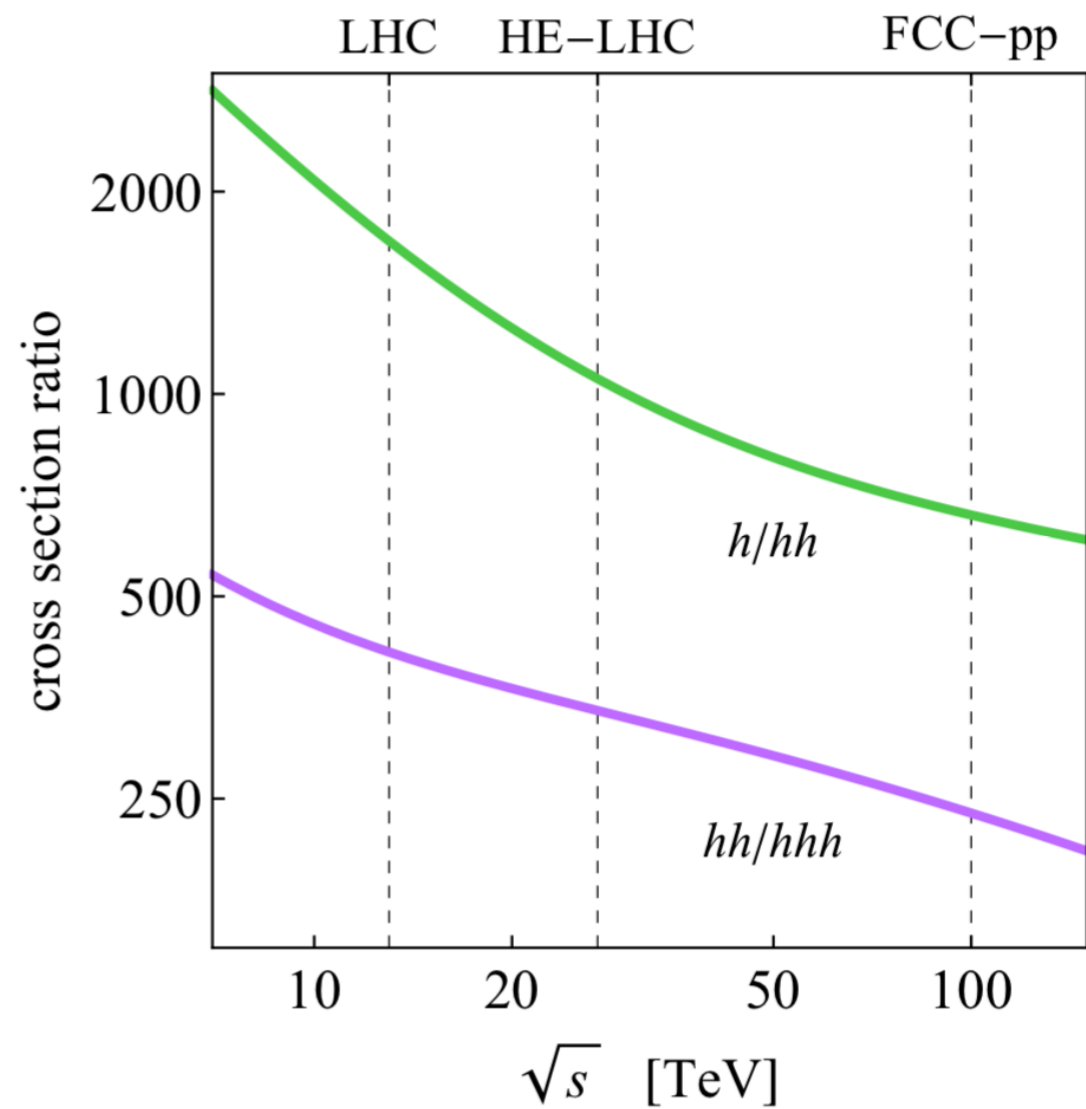
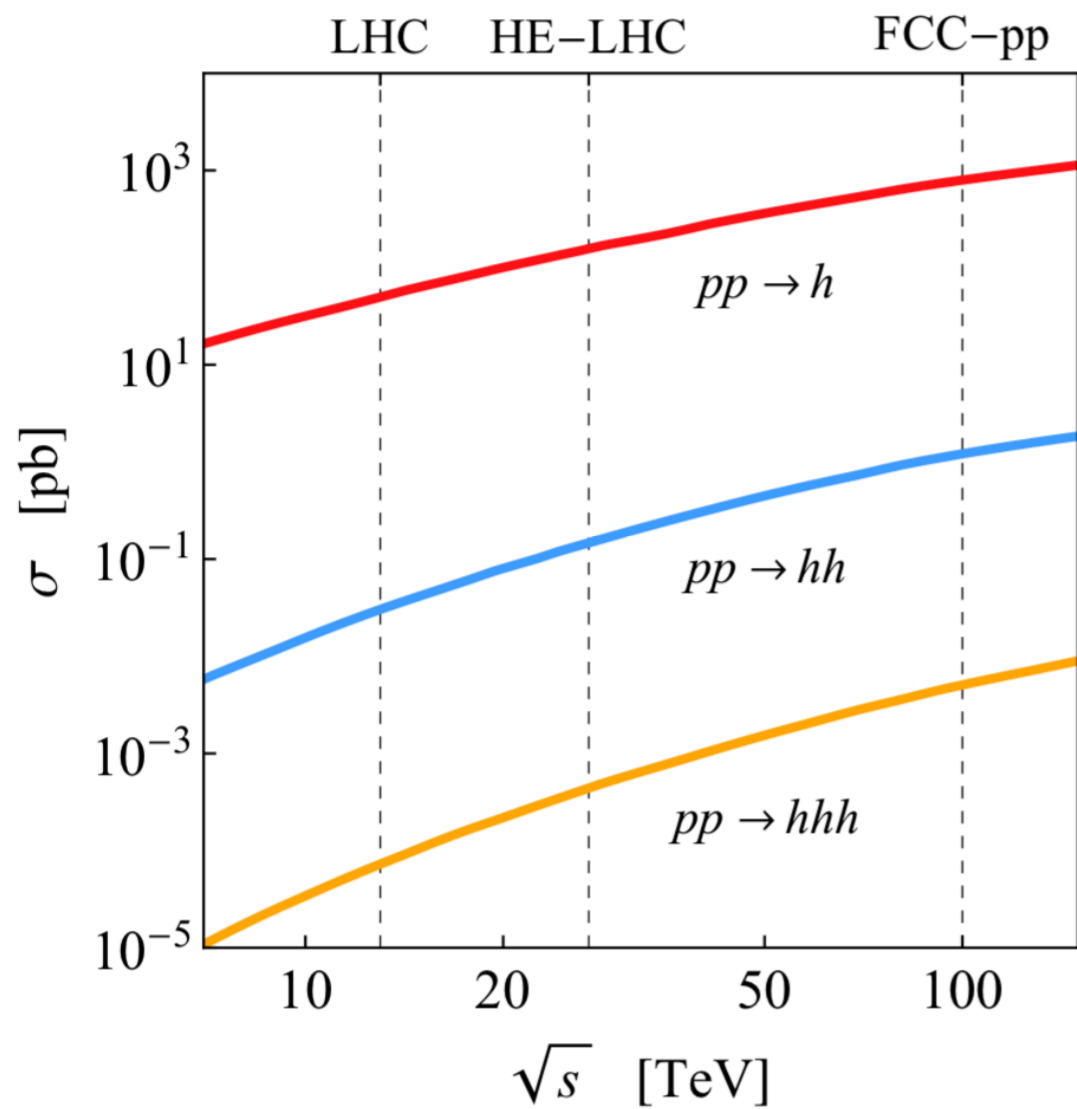
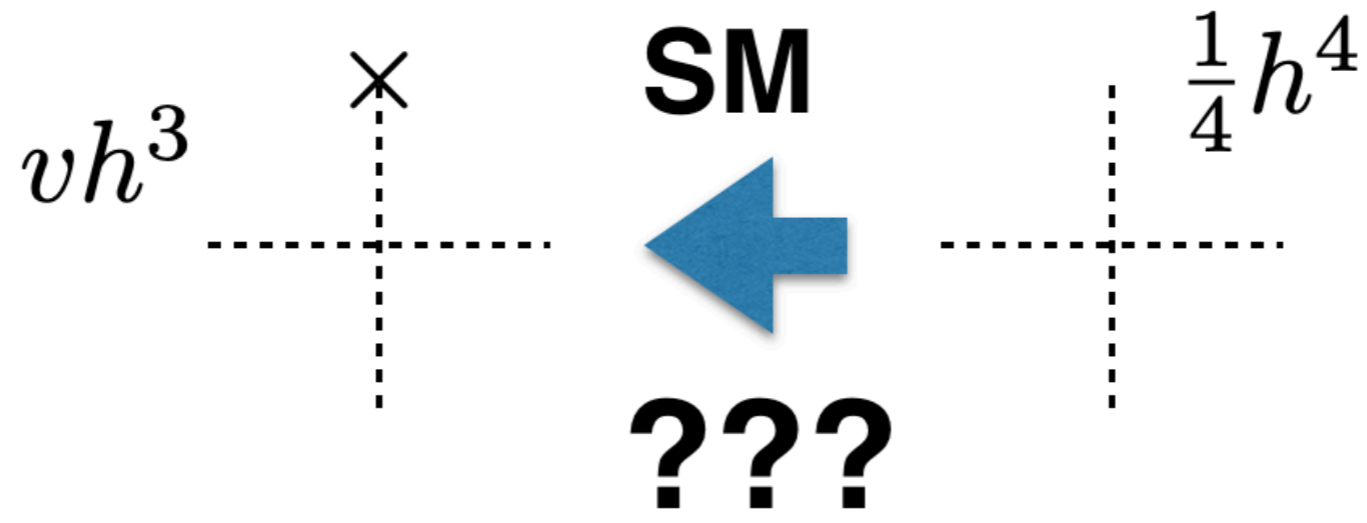
- From spontaneous symmetry breaking

$$V_h^{\text{SM}} = \frac{m_h^2}{2} h^2 + \frac{m_h^2}{2v^2} v h^3 + \frac{1}{4} \frac{m_h^2}{2v^2} h^4$$

- check vh^3  **SM** $\frac{1}{4} h^4$
???

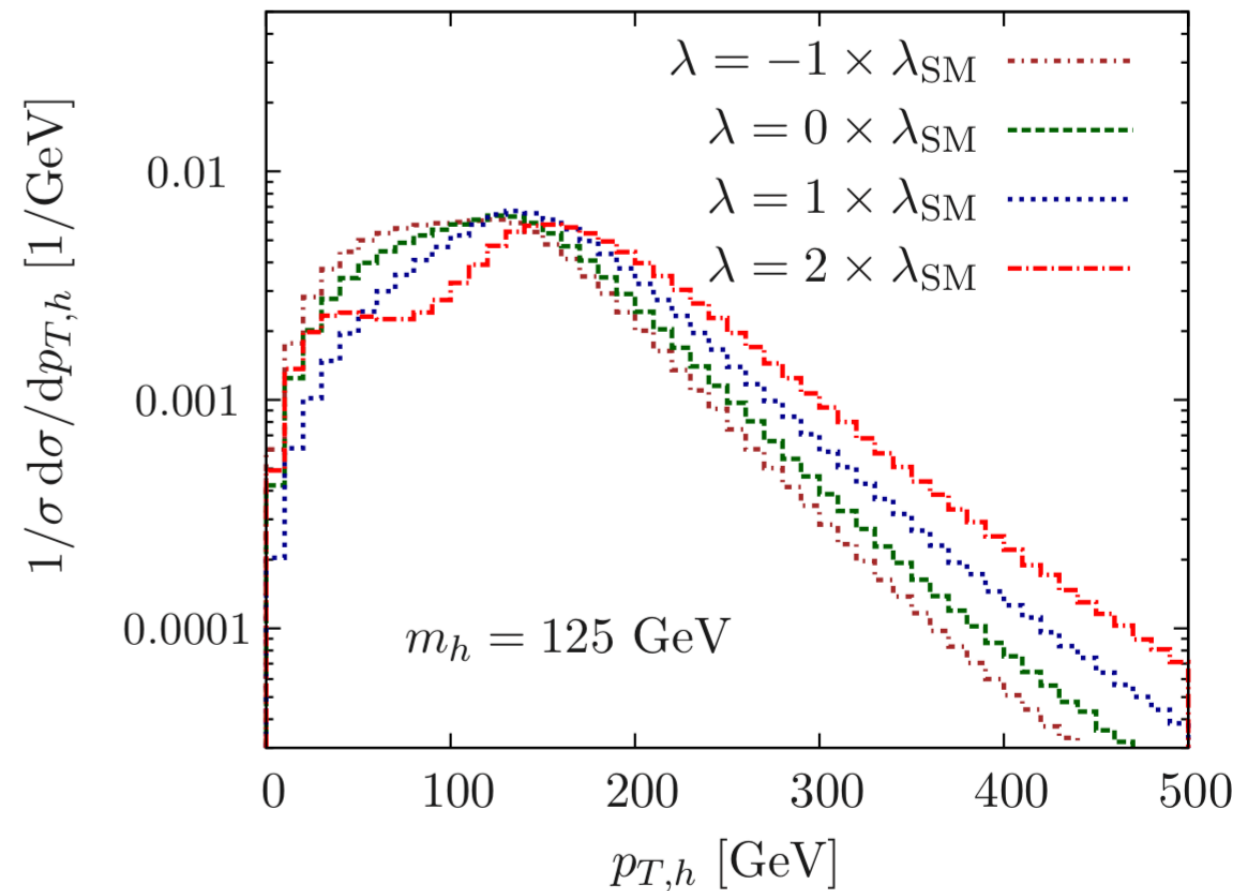
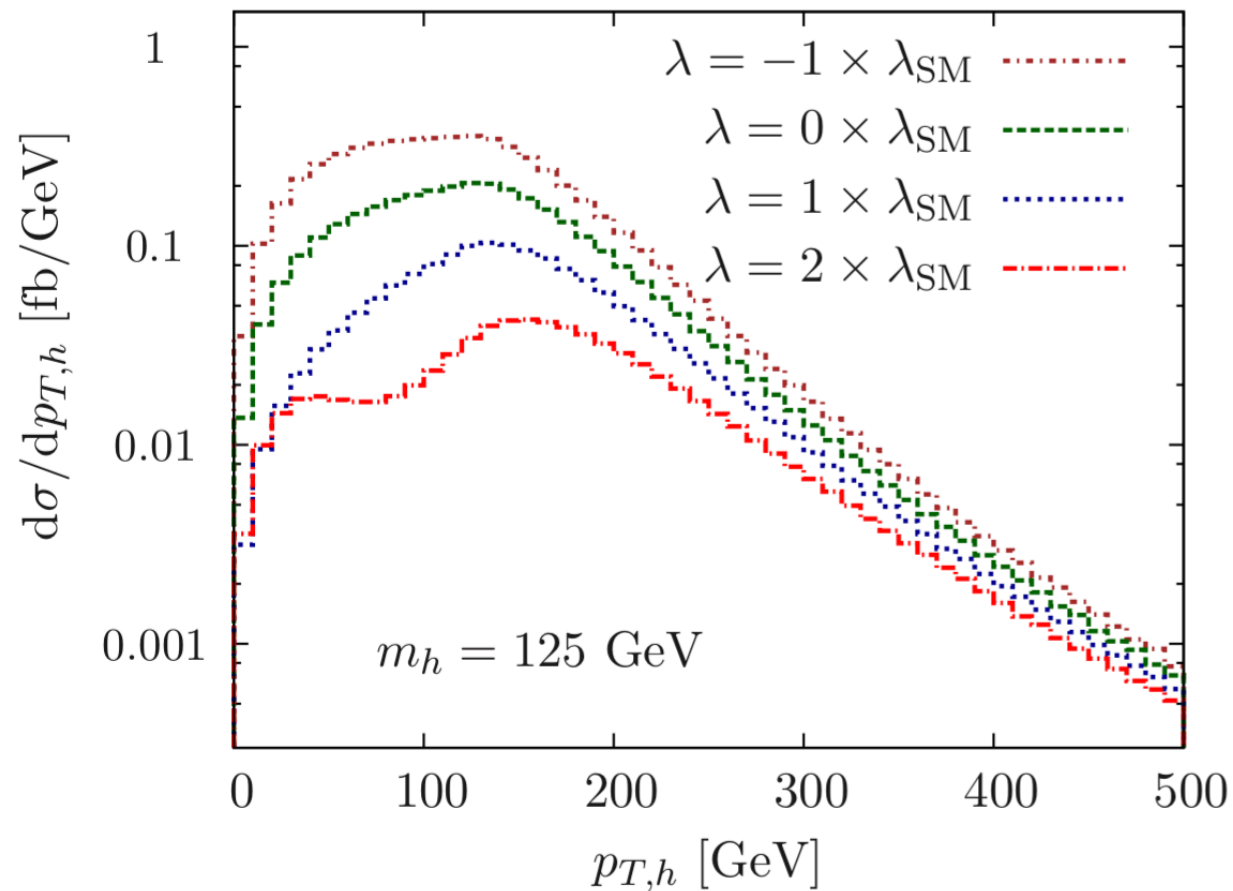
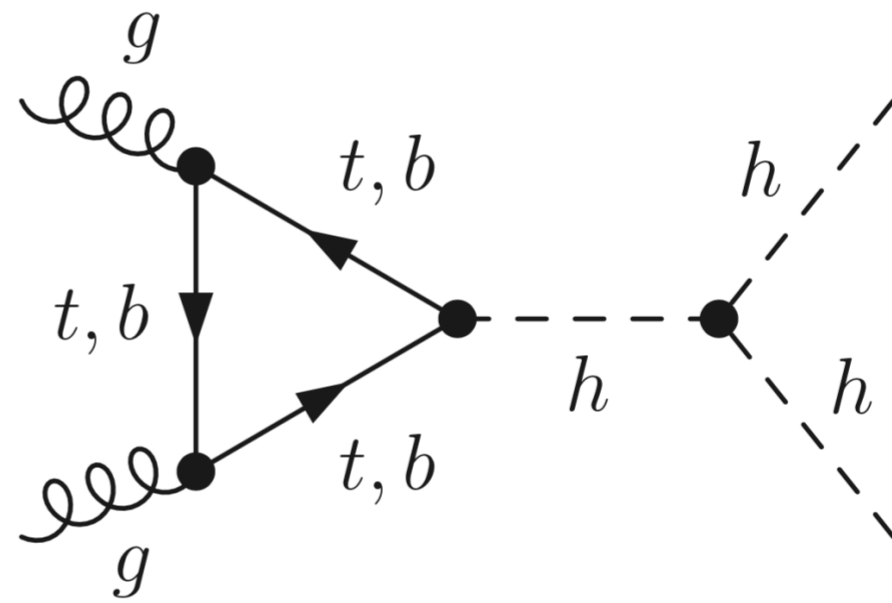
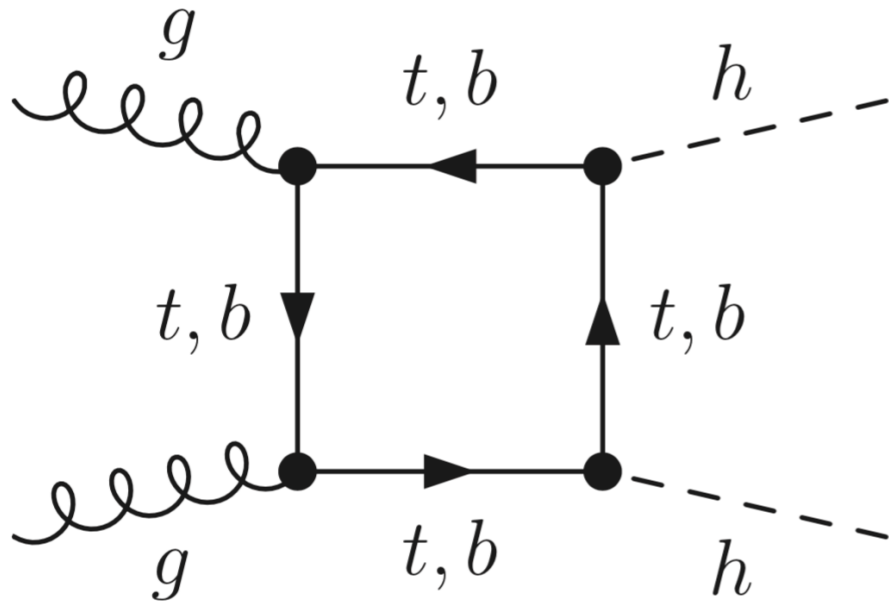
- Triple Higgs coupling can be determined by $h \rightarrow h, h$

- Quartic Higgs coupling can be determined by $h \rightarrow h, h, h$



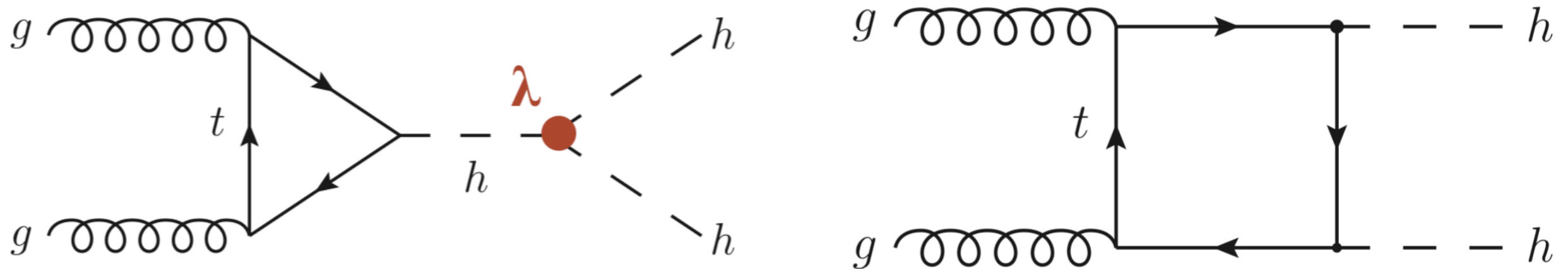
- So, let's focus on "**precise measurement**" of triple Higgs coupling at the LHC

Higgs Boson Pair @LHC

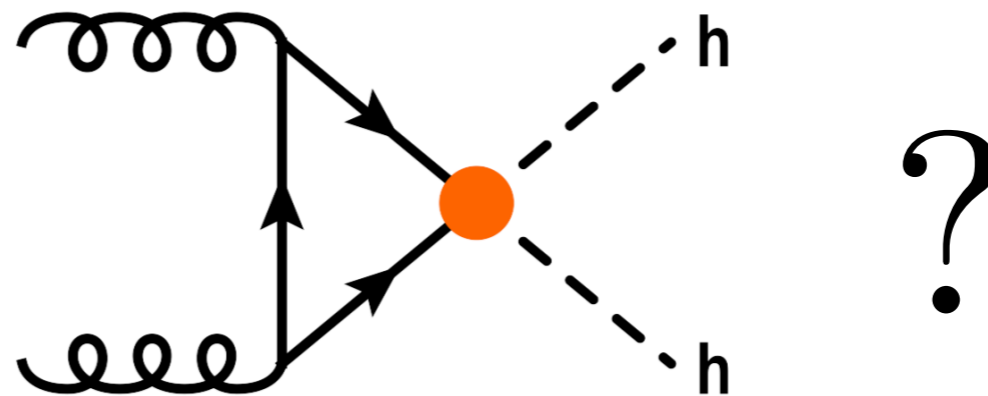


- This triple Higgs coupling can be determined by $h \rightarrow h, h$

Standard Model



+



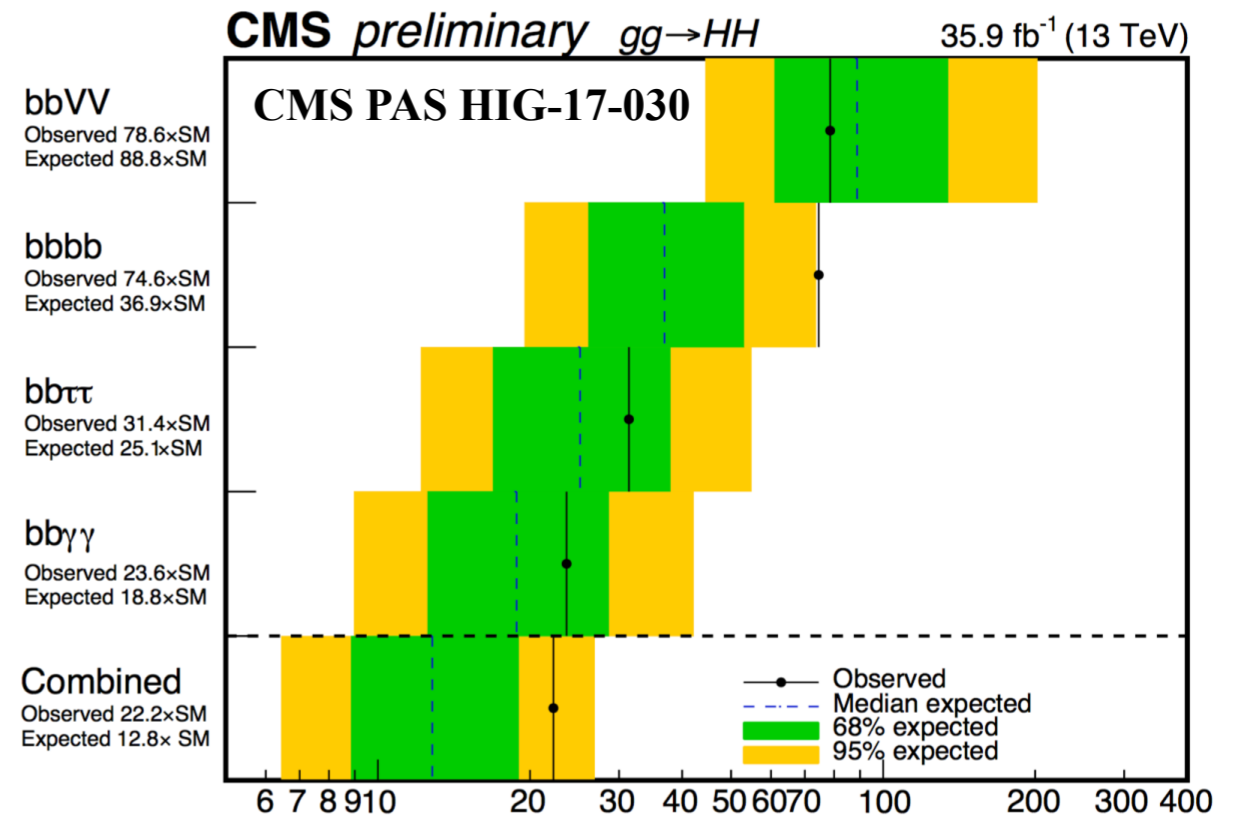
Non-resonant di-Higgs: can be probed

Higgs Boson Pair @LHC

- Current status from various channels

$h \rightarrow XX$

$XX \leftarrow h$	bb	WW^*	$\tau\tau$	ZZ^*	$\gamma\gamma$
bb	33%				
WW^*	25%	4.6%			
$\tau\tau$	7.3%	2.7%	0.39%		
ZZ^*	3.1%	1.1%	0.33%	0.069%	
$\gamma\gamma$	0.26%	0.1%	0.028%	0.012%	0.0005%

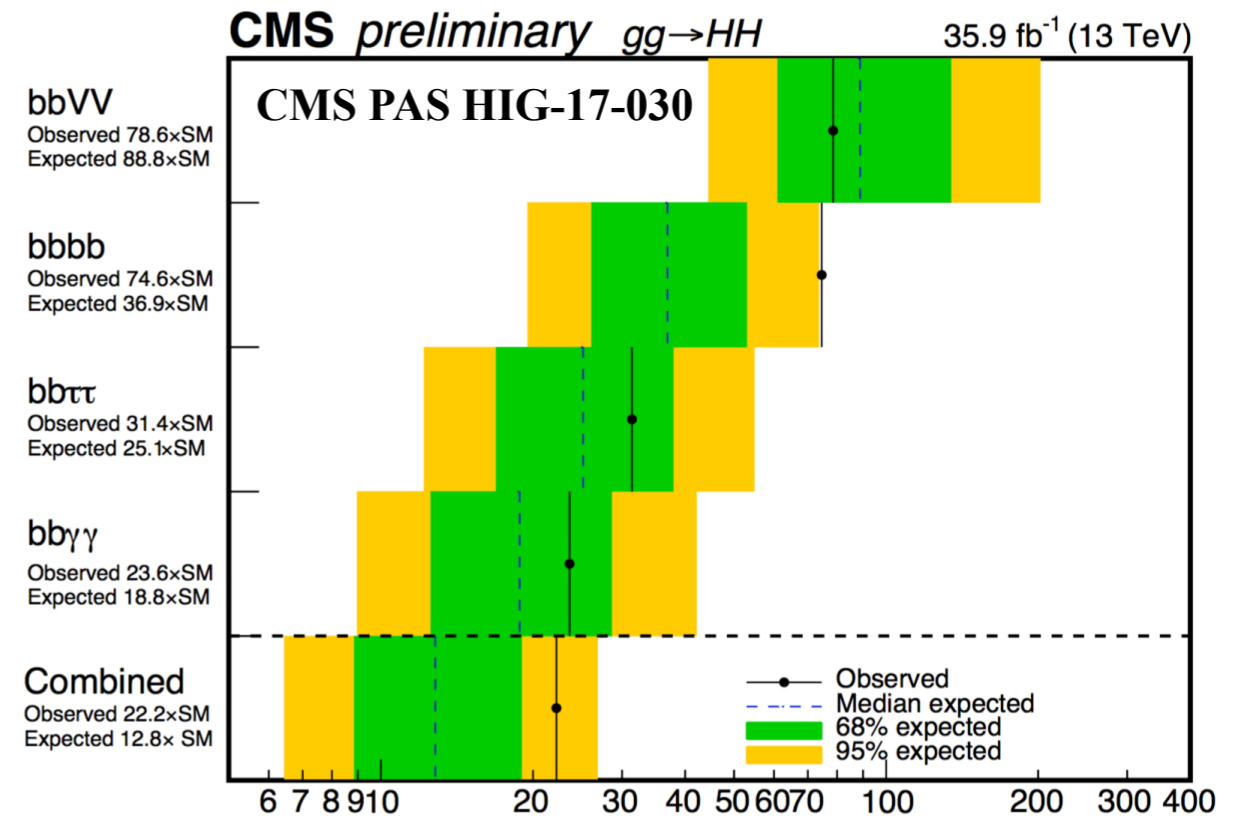


- The driven channel is the "compromised" clean channel.

Higgs Boson Pair @LHC

- Current status from various channels

		$N(hh)_{SM}$	N_{BKG}	
<i>ATLAS</i>	$bb\gamma\gamma$	8.4	47.1	1.2
<i>CMS</i>	$bb\gamma\gamma$	9	26.9	1.7
	$bb\tau\tau$ (fully-hadronic)	4.9	30.3	0.89
	$bb\tau\tau$ (semi-leptonic)	6.1	122	0.55
	$bbWW^*$ (di-leptonic)	37.1	3875	0.60

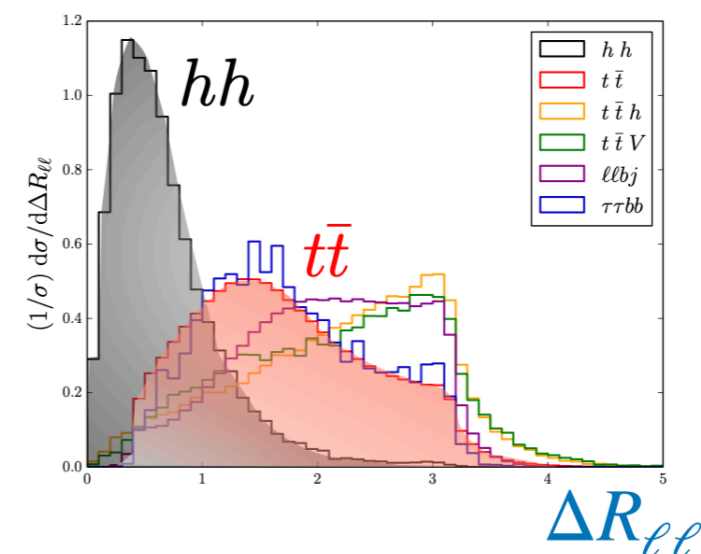
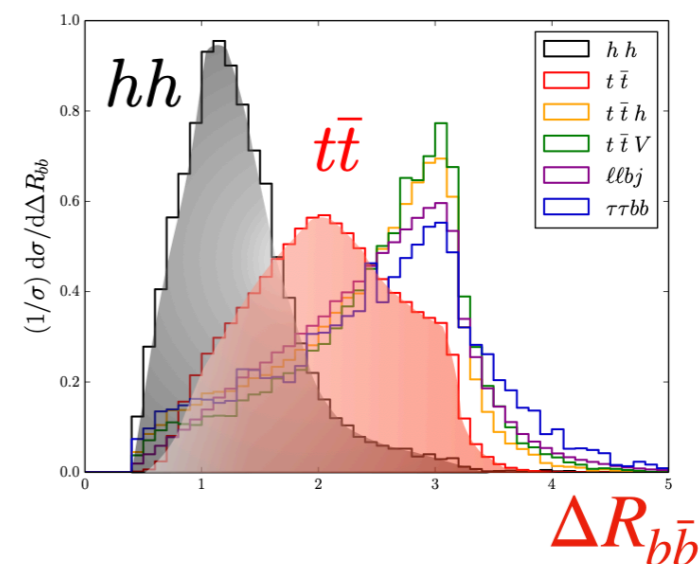
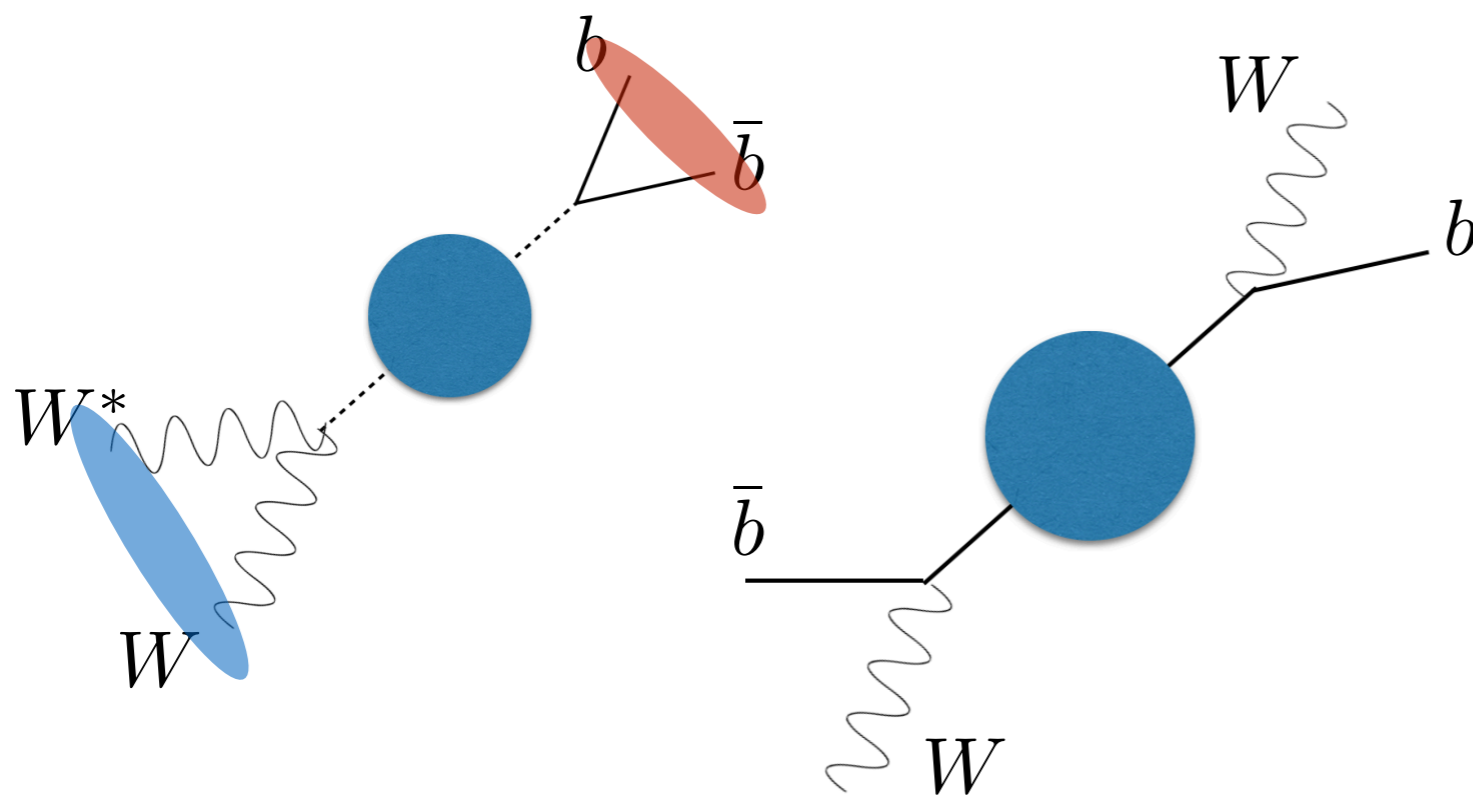


- **Why** is $bbVV$ **bad** ? **how** can one **improve** ?

- **Why is $bbVV$ bad ?** LHC is the **Top-factory**

$$\frac{\sigma(pp \rightarrow hh \rightarrow b\bar{b}VV^*)}{\sigma(pp \rightarrow t\bar{t} \rightarrow b\bar{b}VV)} \Bigg|_{13\text{TeV}} \simeq \frac{31\text{fb}*(25\%)}{215\text{pb}} \simeq \mathcal{O}(10^{-5})$$

- Applying "**low-level**" kinematic cuts based on event-topology



- Applying "**low-level**" kinematic cuts based on event-topology

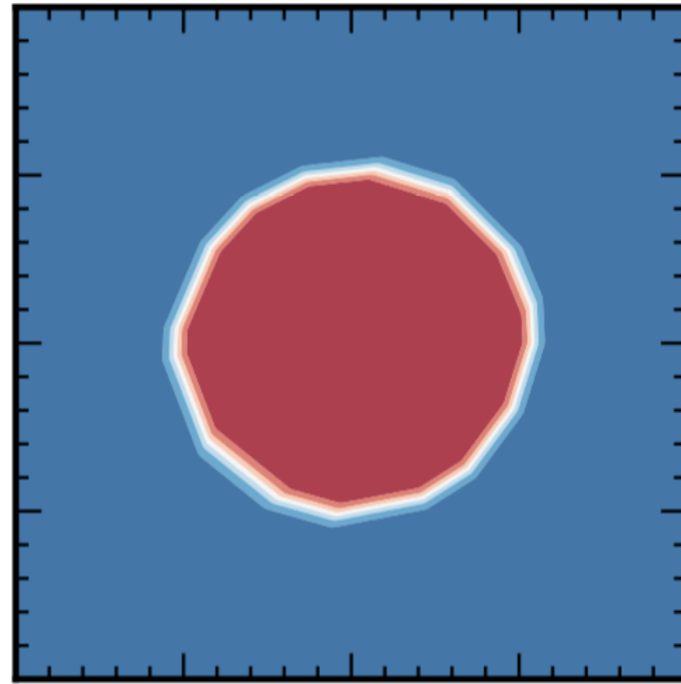
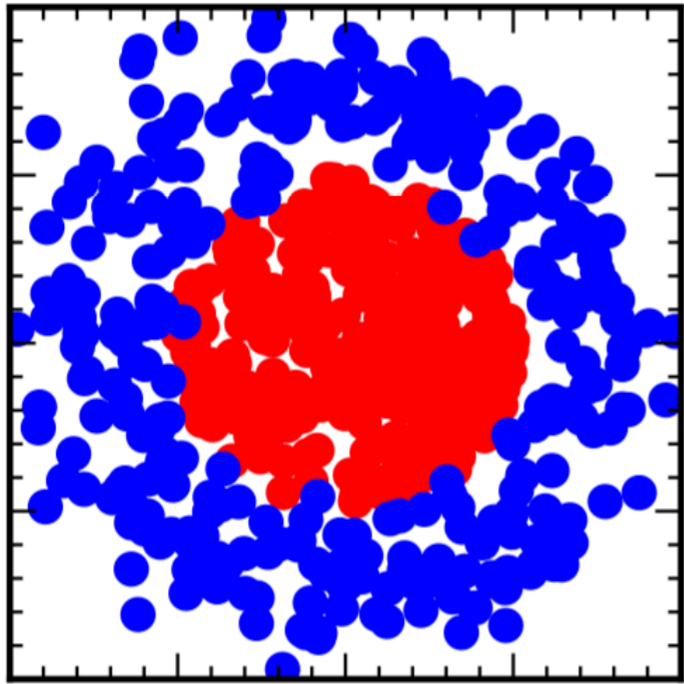
Baseline selections: $\cancel{E}_T > 20 \text{ GeV}$,
 $p_T^\ell > 20 \text{ GeV}$, $\Delta R_{\ell\ell} < 1.0$, $m_{\ell\ell} < 65 \text{ GeV}$,
 $\Delta R_{bb} < 1.3$, $95 < m_{bb} < 140 \text{ GeV}$

Signal	$t\bar{t}$	$t\bar{t}h$	$t\bar{t}V$	$\ell\ell bj$	$\tau\tau bb$	others	σ	$N_{\text{sig}}^{\text{SM}} / N_{\text{bknd}}$
0.0124	1.1724	0.0297	0.0246	0.0158	0.0379	0.00590	0.60	0.00964

jjllνν̄ backgrounds from QCD+EW

- We may apply the advanced statistical tools to see correlations among "low-level" kinematic variables.
- But the **efficiency based on "low-level cuts" is NOT GOOD** in that case.

what if we know BKG very well



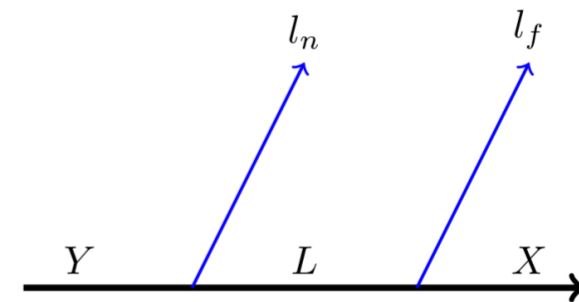
A cartoon is from a paper by Spencer Chang, Timothy Cohen and Bryan Ostdiek, Phys. Rev. D 97, 056009 (2018)

- If we know BKG is a circle, we can take BKG out very easily with **input variables which characterizes** BKG events

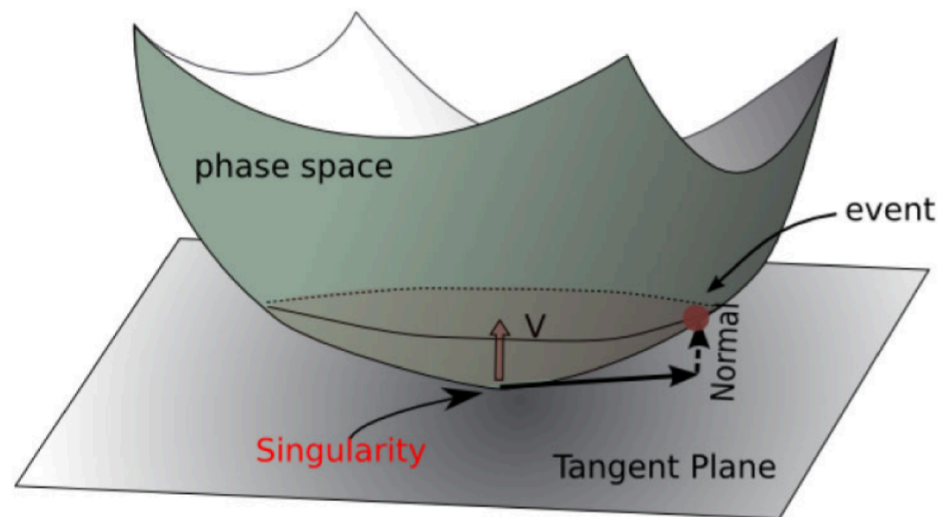
(for a circle, variables are; a location of center, radius)

- Considering "**high-level**" kinematic cuts based on event-topology
- What are the "**high-level**" kinematic variables?

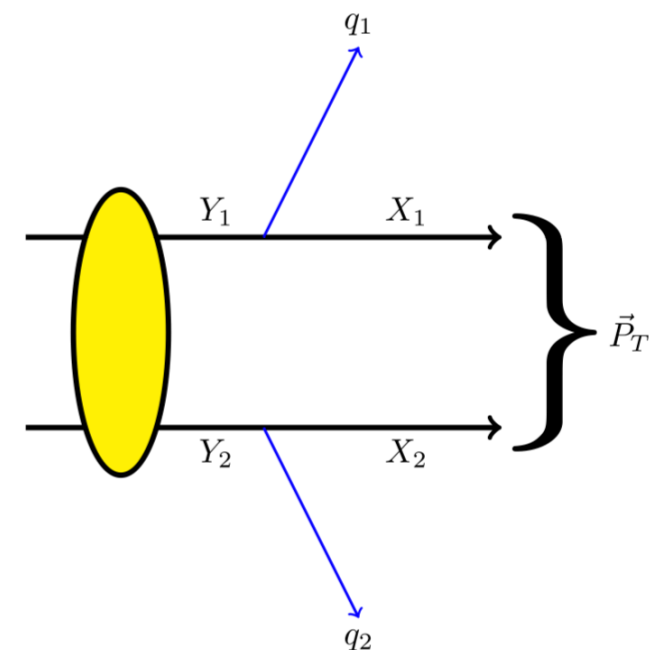
invariant mass



when one considers only visibles

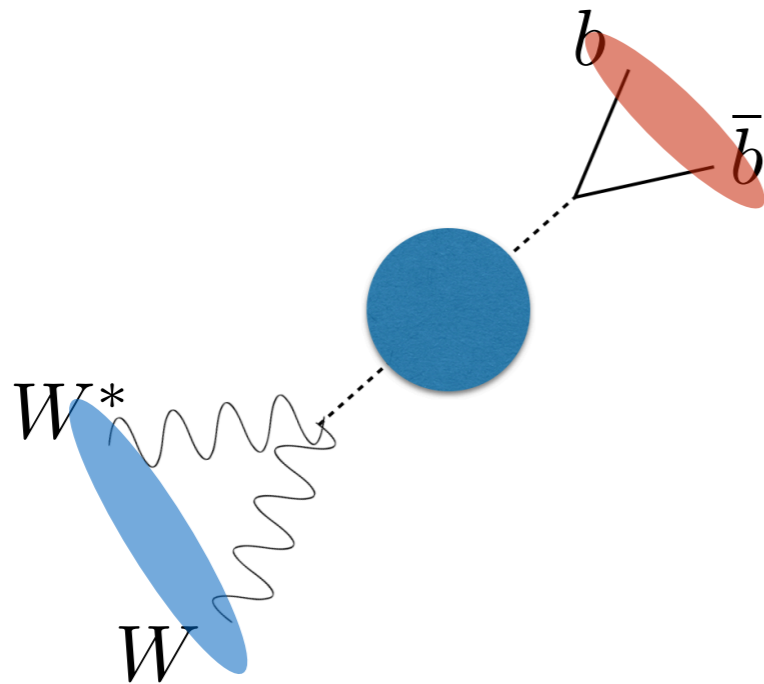


transverse mass



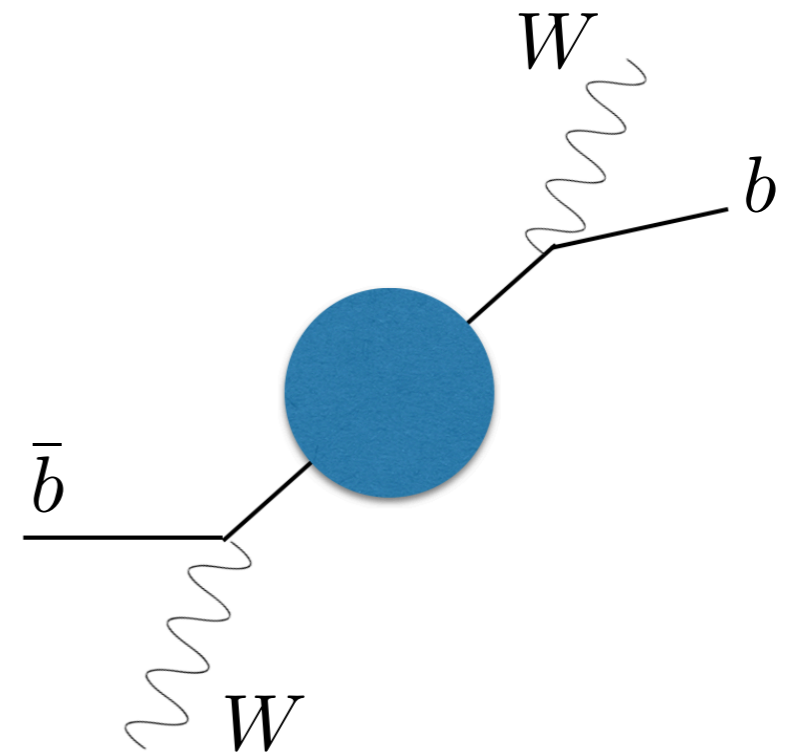
when one considers visibles+invisibles

- Considering "**high-level**" kinematic cuts based on event-topology **invariant mass**



$$\begin{aligned}
 H \equiv \min & \left[\frac{(m_{\ell^+\ell^-\nu\bar{\nu}}^2 - m_h^2)^2}{\sigma_{h\ell}^4} + \frac{(m_{\nu\bar{\nu}}^2 - m_{\nu\bar{\nu},peak}^2)^2}{\sigma_{\nu}^4} \right. \\
 & + \min \left(\frac{(m_{\ell^+\nu}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell^-\bar{\nu}}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4}, \right. \\
 & \left. \left. \frac{(m_{\ell^-\bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} + \frac{(m_{\ell^+\nu}^2 - m_{W^*,peak}^2)^2}{\sigma_{W^*}^4} \right) \right]
 \end{aligned}$$

Higgs(to WW*), W, and off-shell W

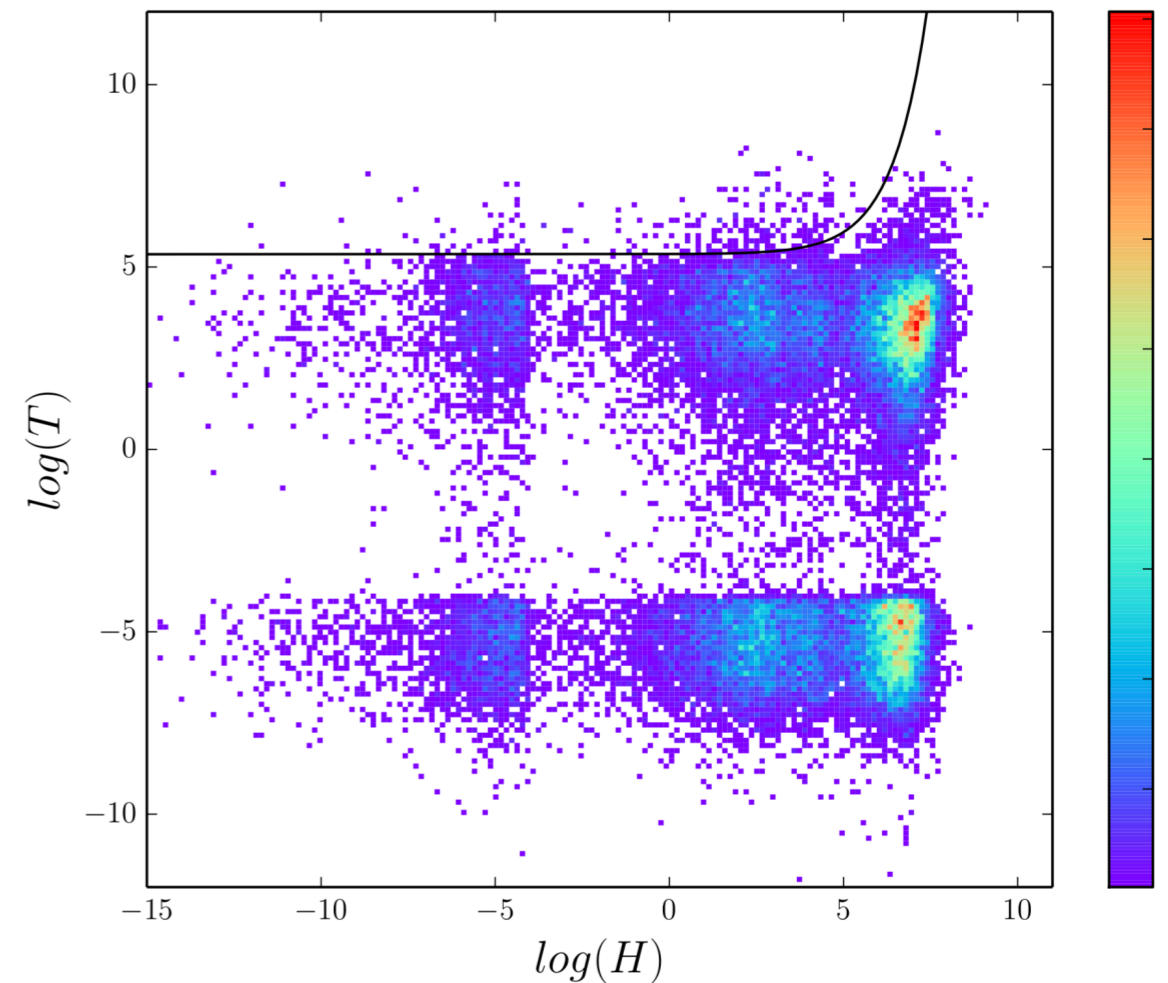
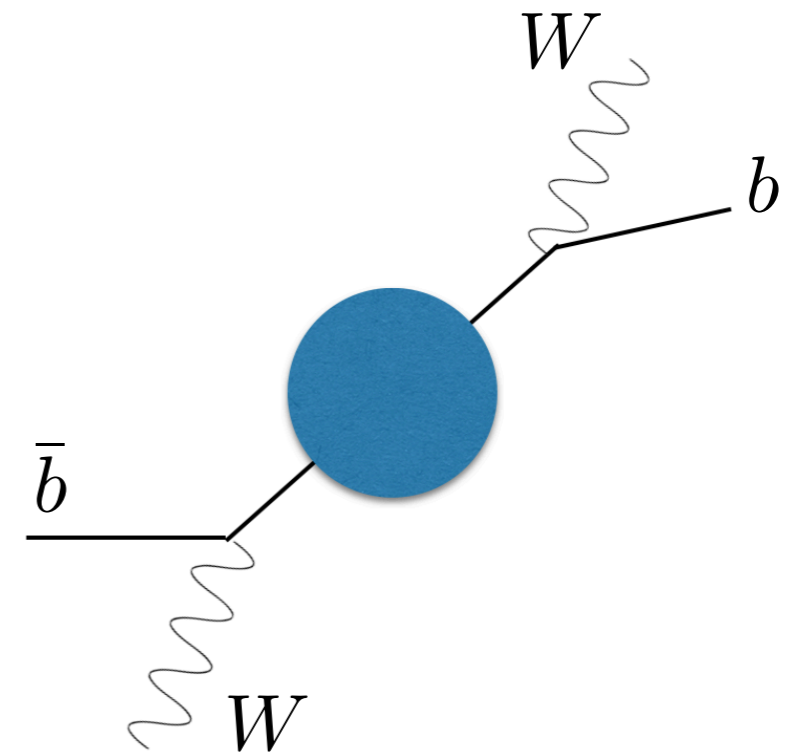
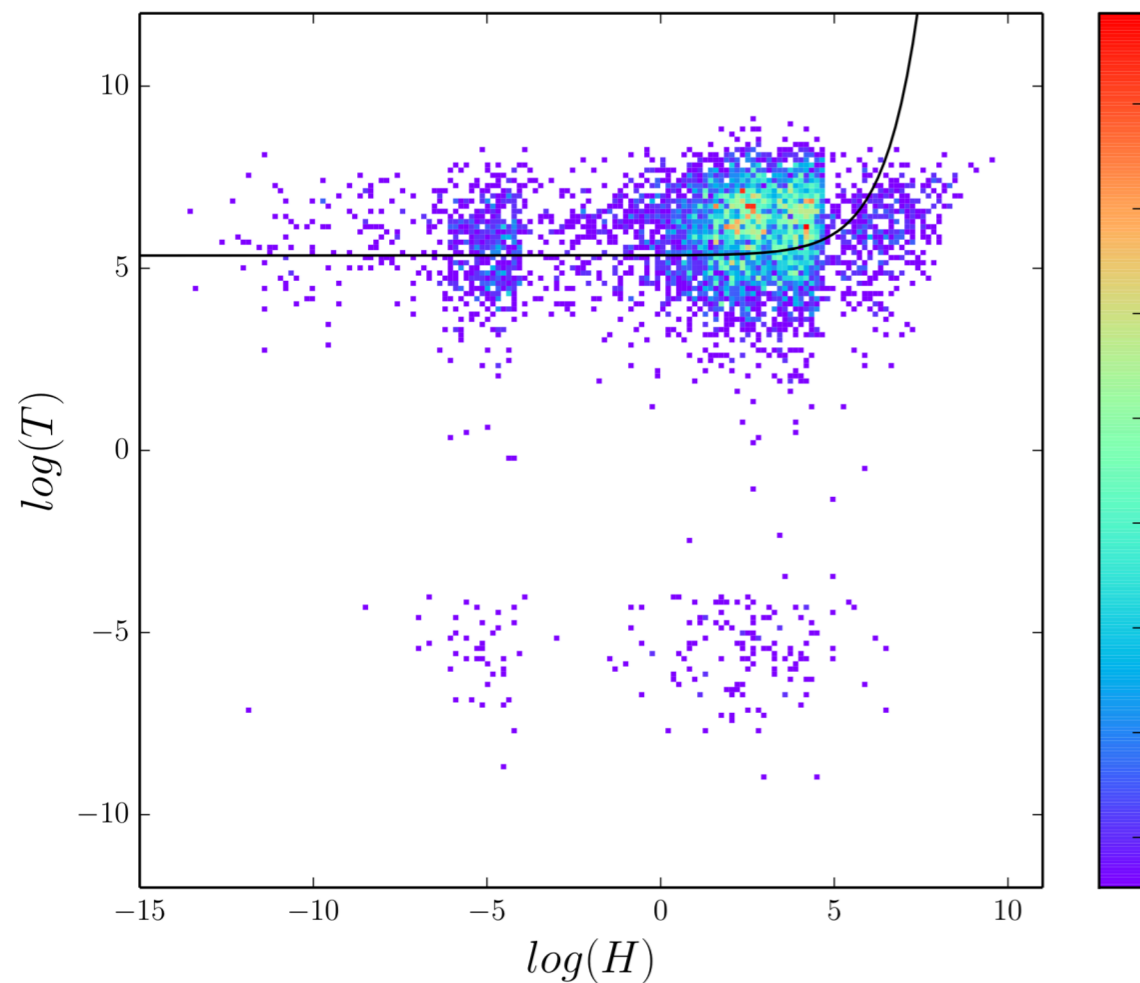
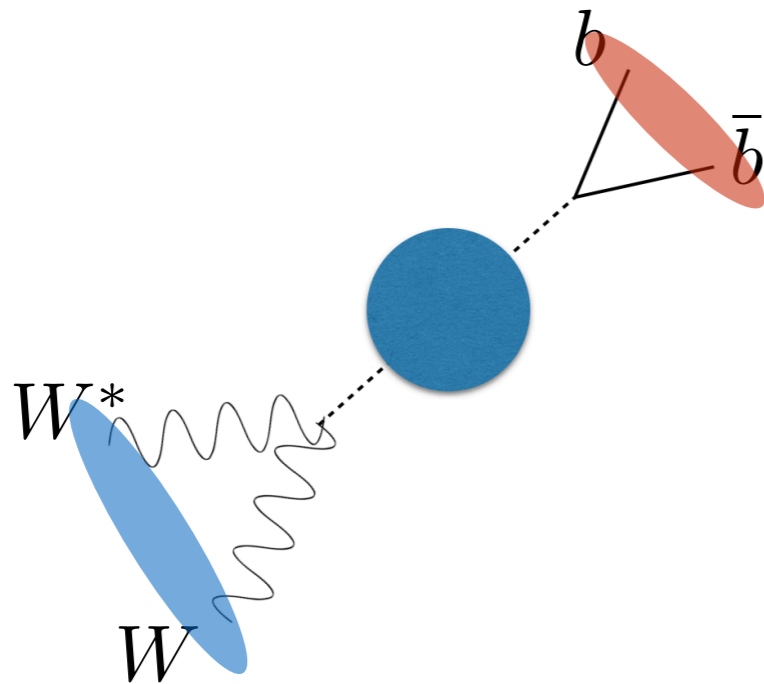


$$\begin{aligned}
 \chi_{ij}^2 \equiv \min_{\vec{p}_T = \vec{p}_{\nu T} + \vec{p}_{\bar{\nu} T}} & \left[\frac{(m_{b_i\ell^+\nu}^2 - m_t^2)^2}{\sigma_t^4} + \frac{(m_{\ell^+\nu}^2 - m_W^2)^2}{\sigma_W^4} \right. \\
 & \left. + \frac{(m_{b_j\ell^-\bar{\nu}}^2 - m_t^2)^2}{\sigma_t^4} + \frac{(m_{\ell^-\bar{\nu}}^2 - m_W^2)^2}{\sigma_W^4} \right]
 \end{aligned}$$

two top quarks and two W

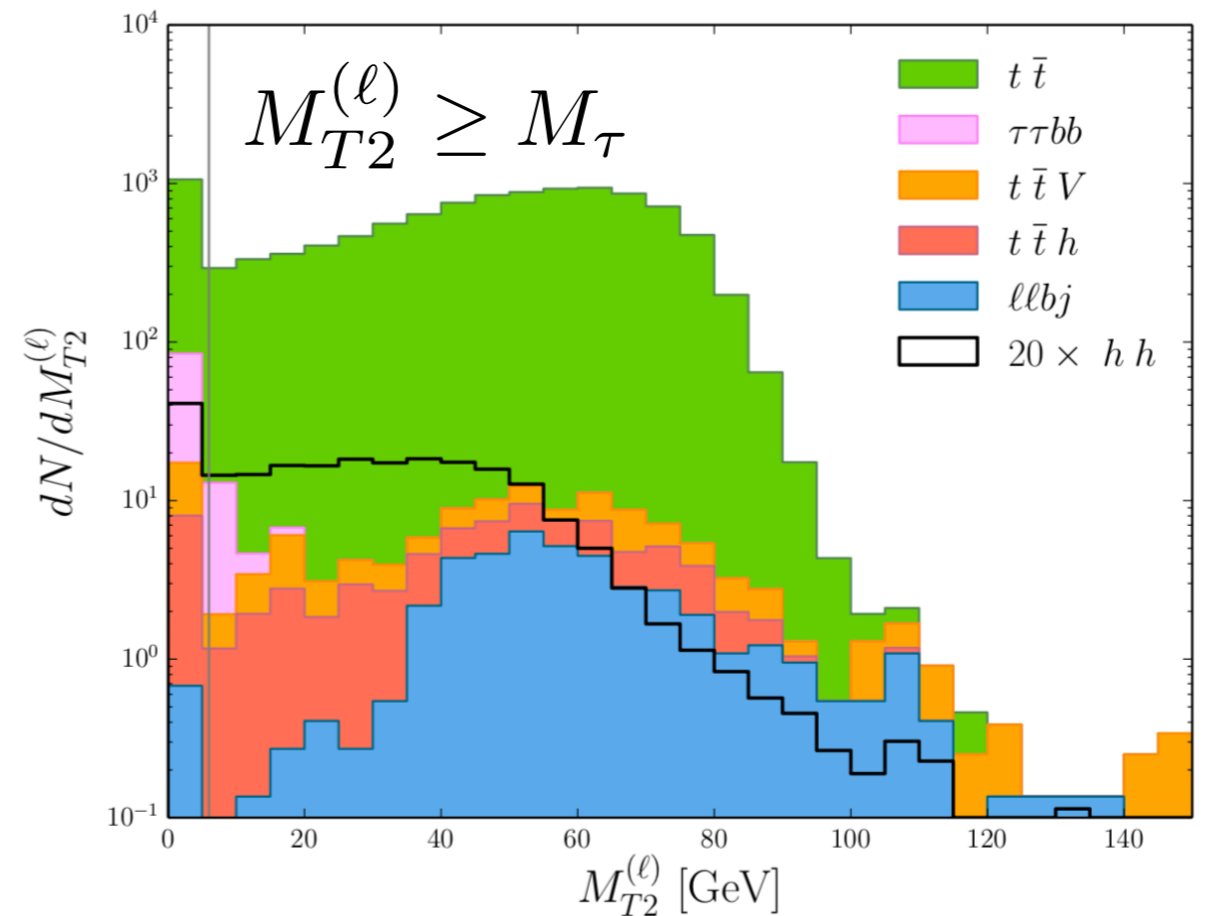
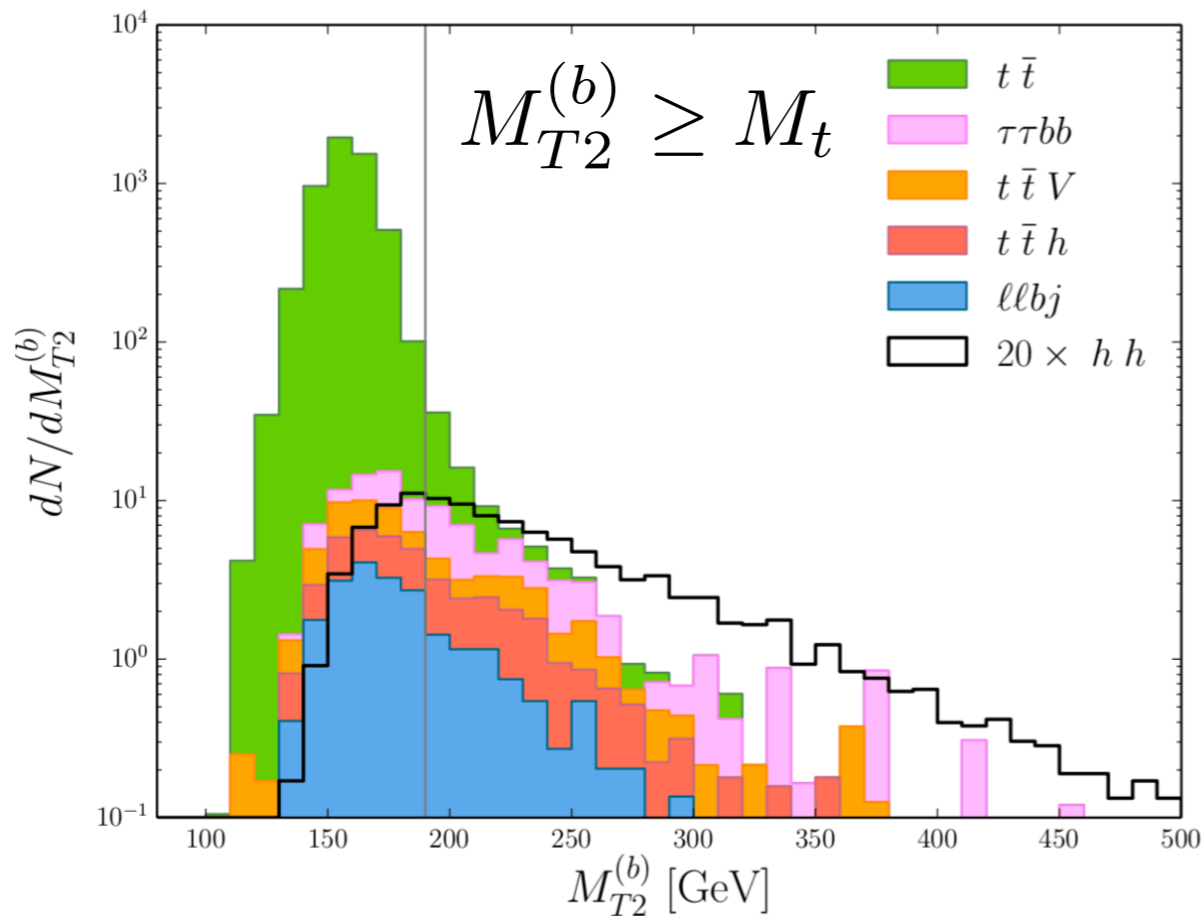
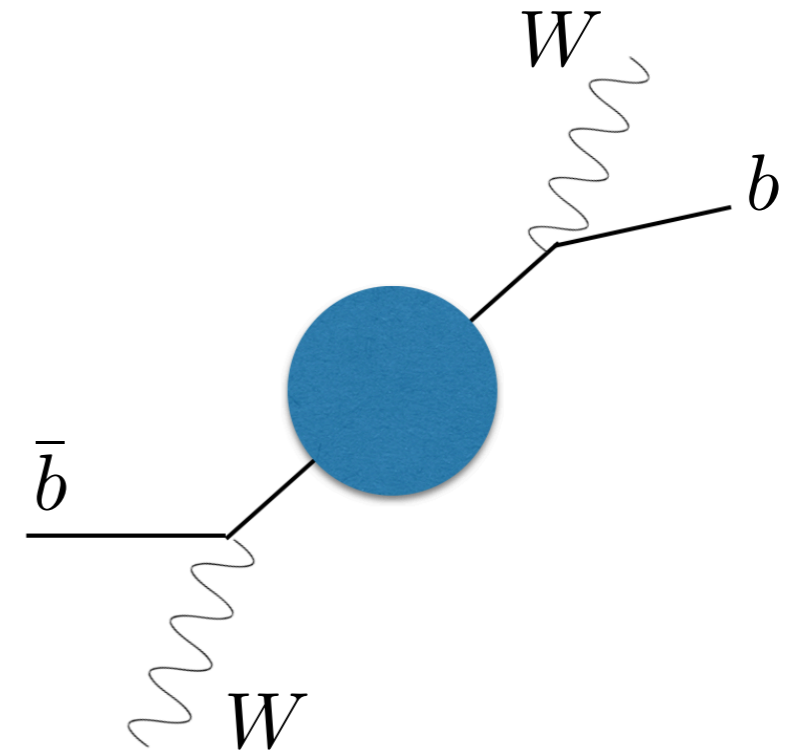
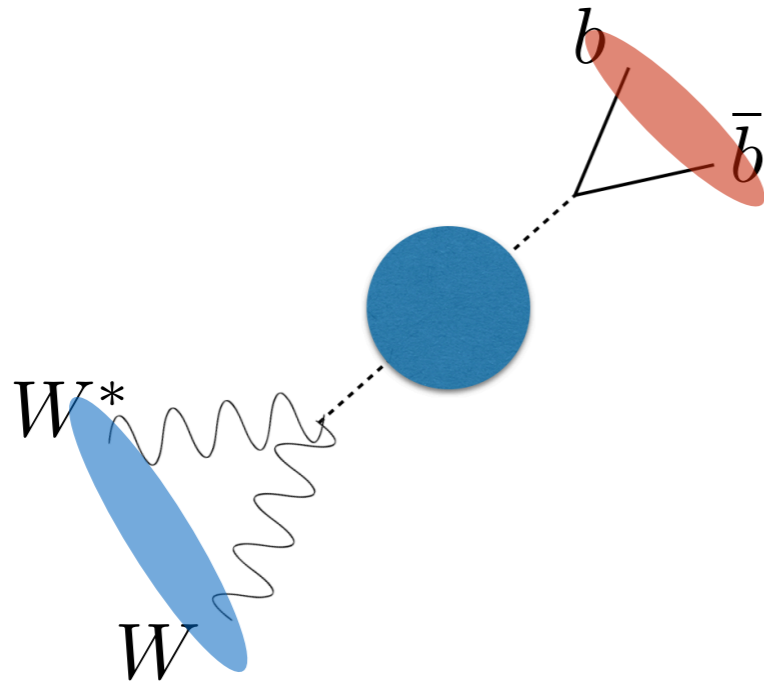
- Considering "**high-level**" kinematic cuts based on event-topology

invariant mass

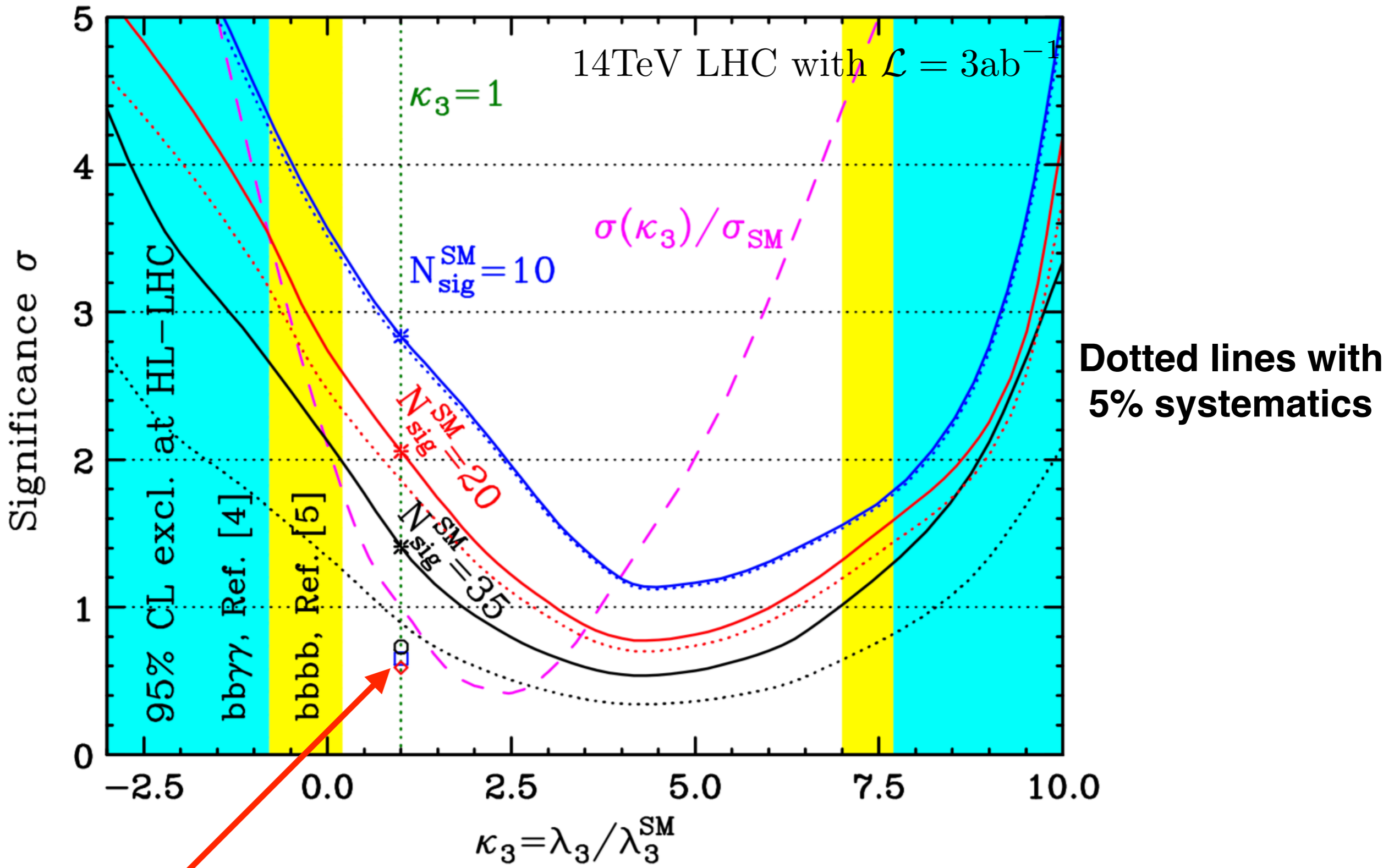


- Considering "**high-level**" kinematic cuts based on event-topology

Transverse mass

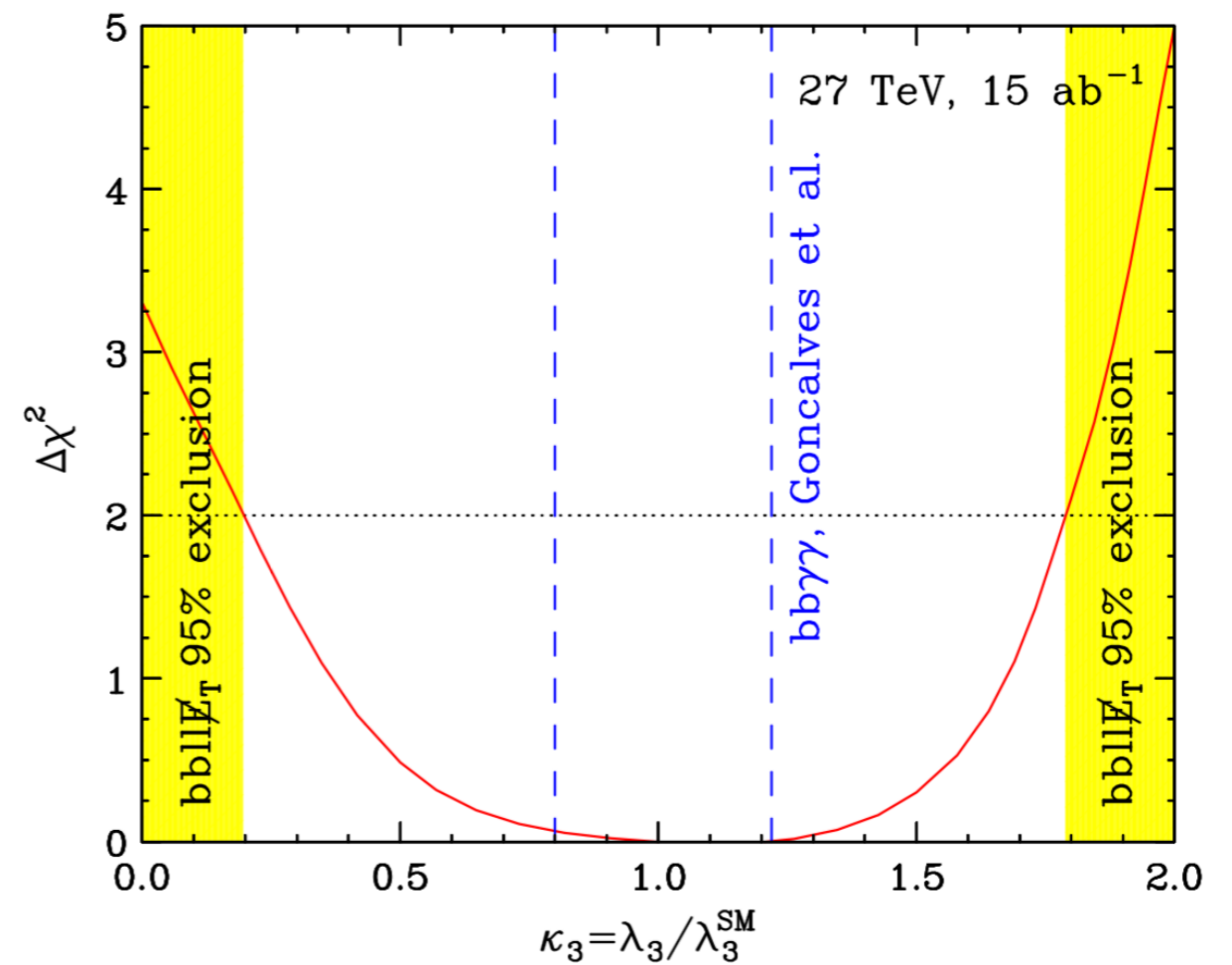
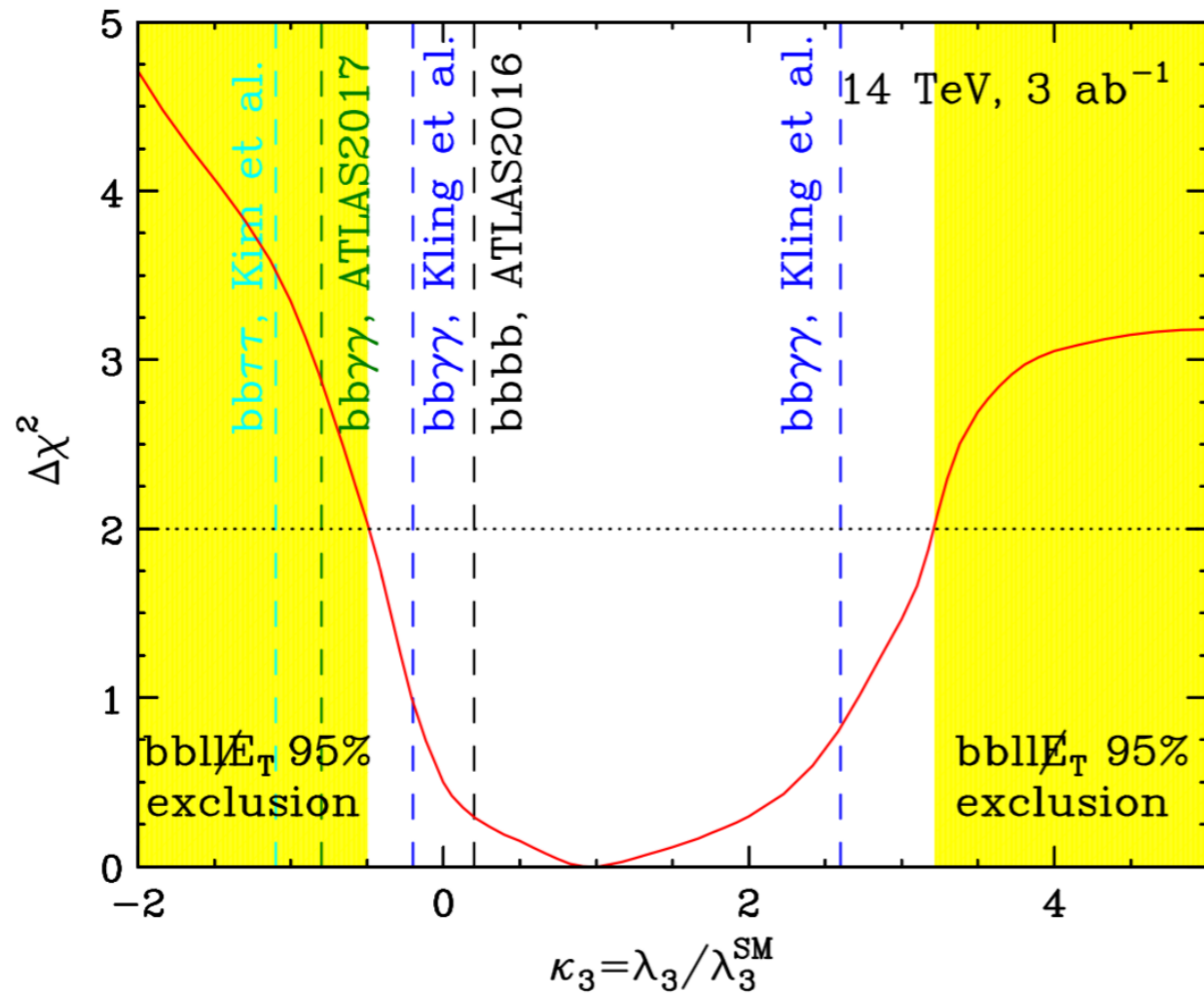


- Discovery chance for triple Higgs coupling
 - Optimization for cut is set for SM case



CMS-NN, CMS-BDT, BDT with basic observables

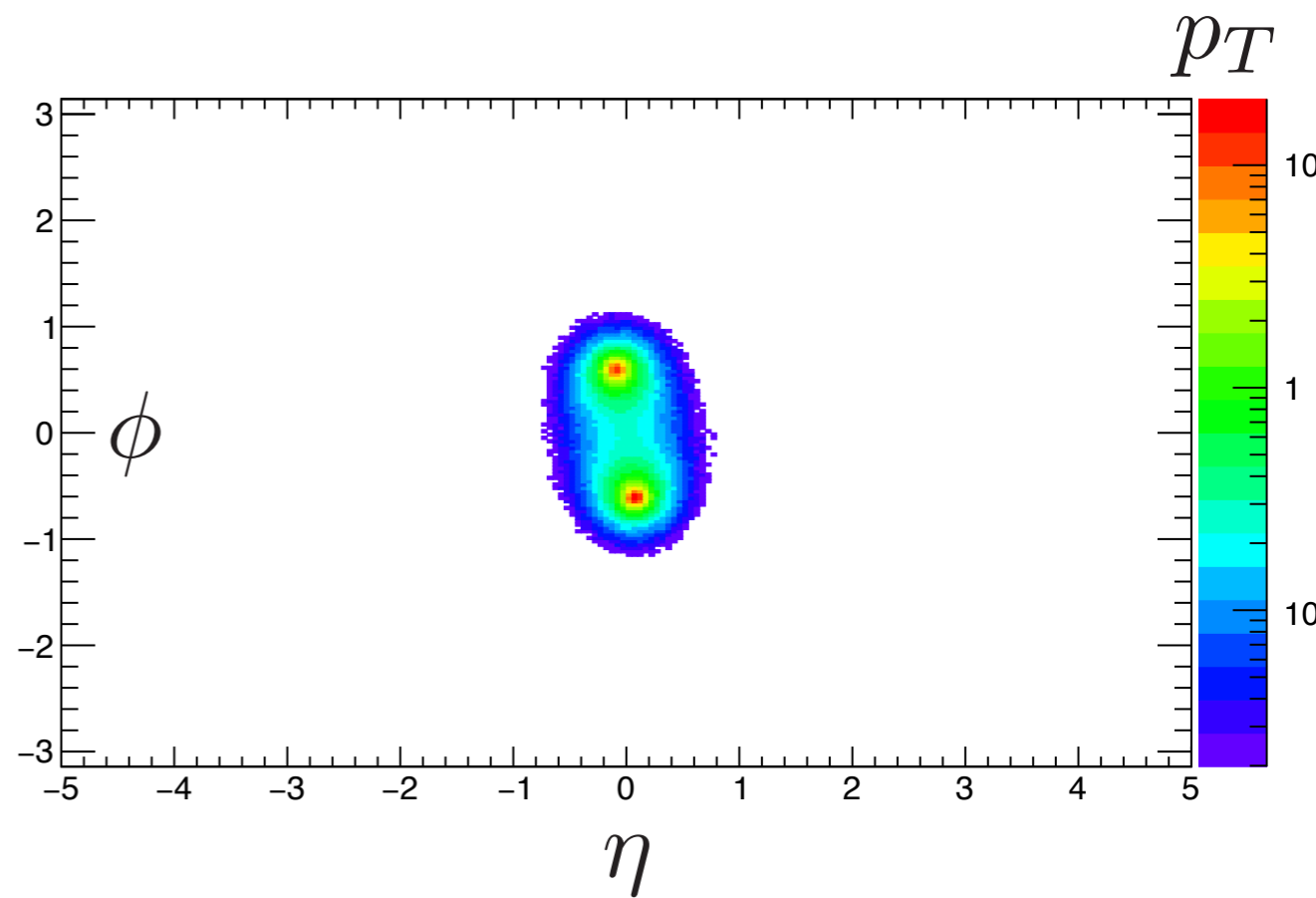
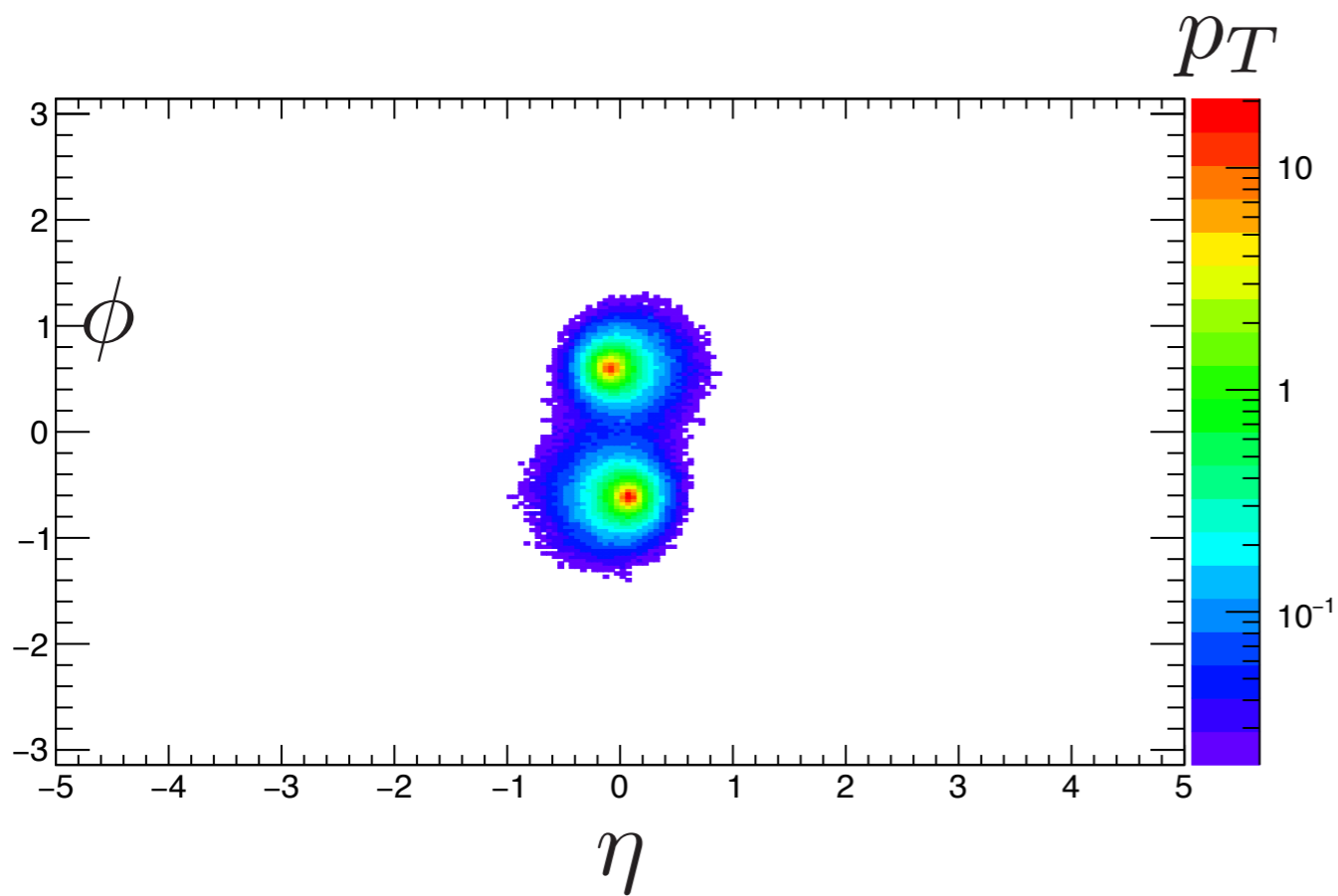
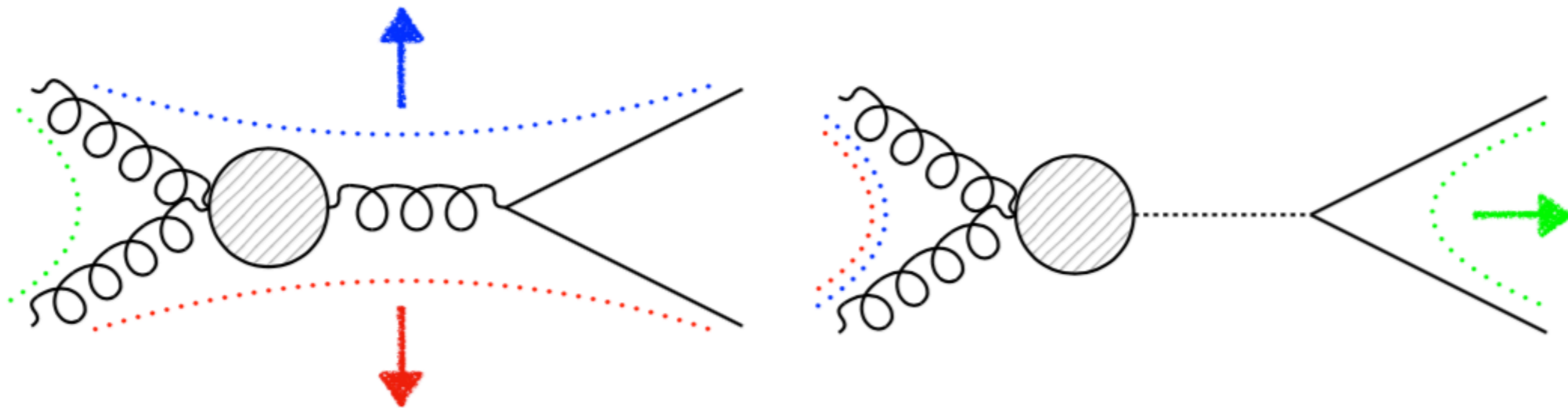
- Significance for observing "**anomalous**" Higgs triple coupling



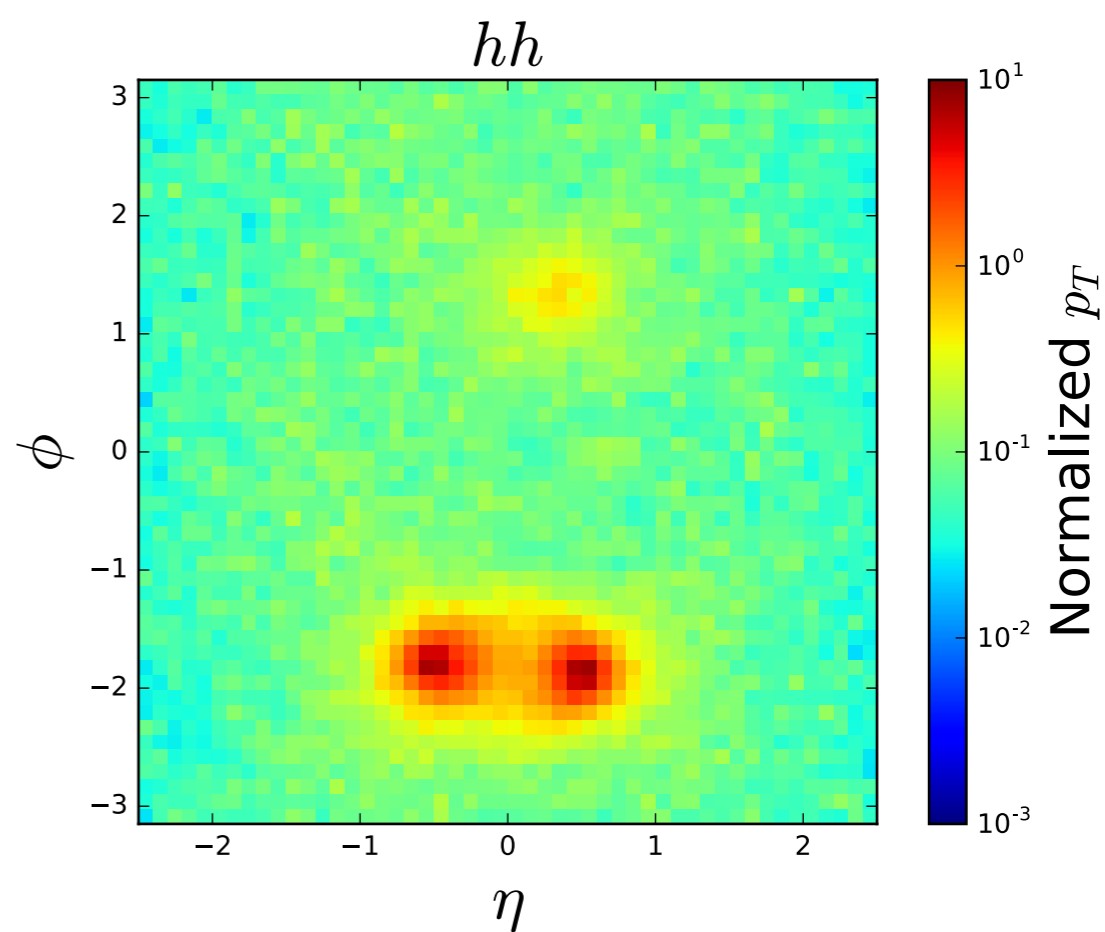
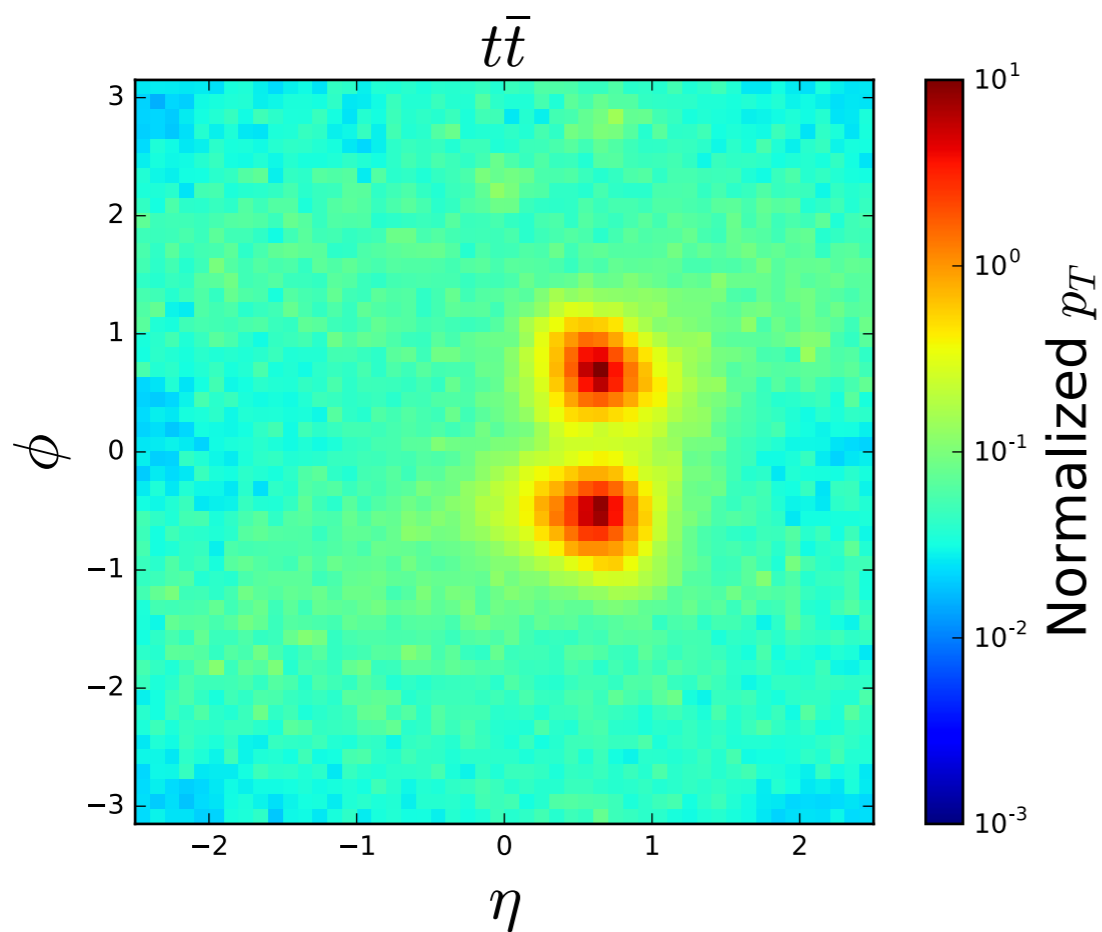
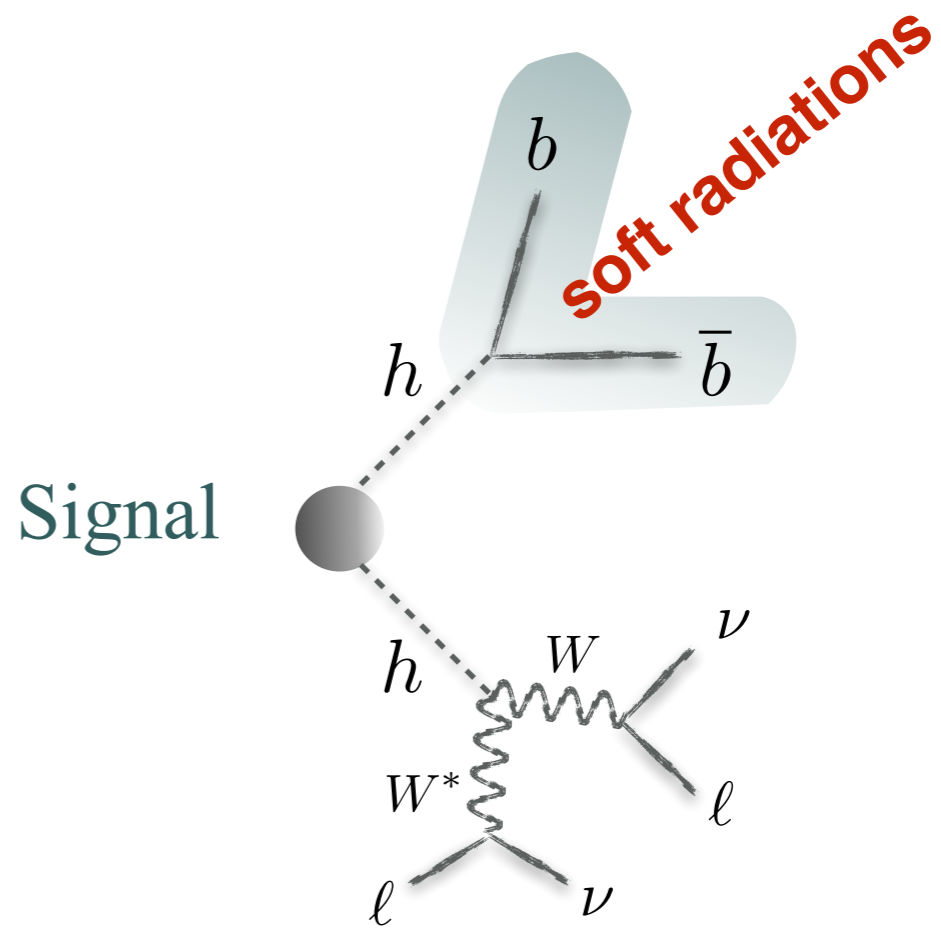
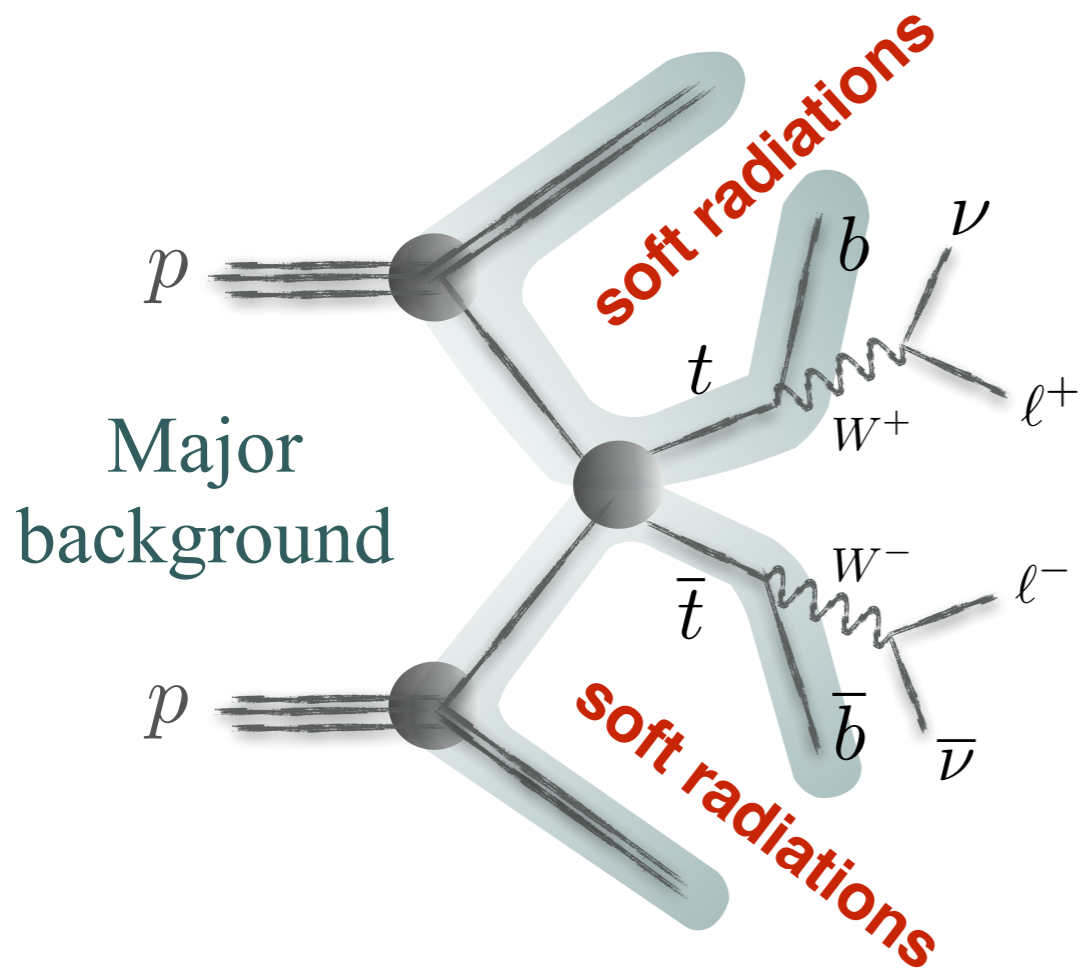
- Here the hypothesis is SM Higgs (triple coupling = SM case)

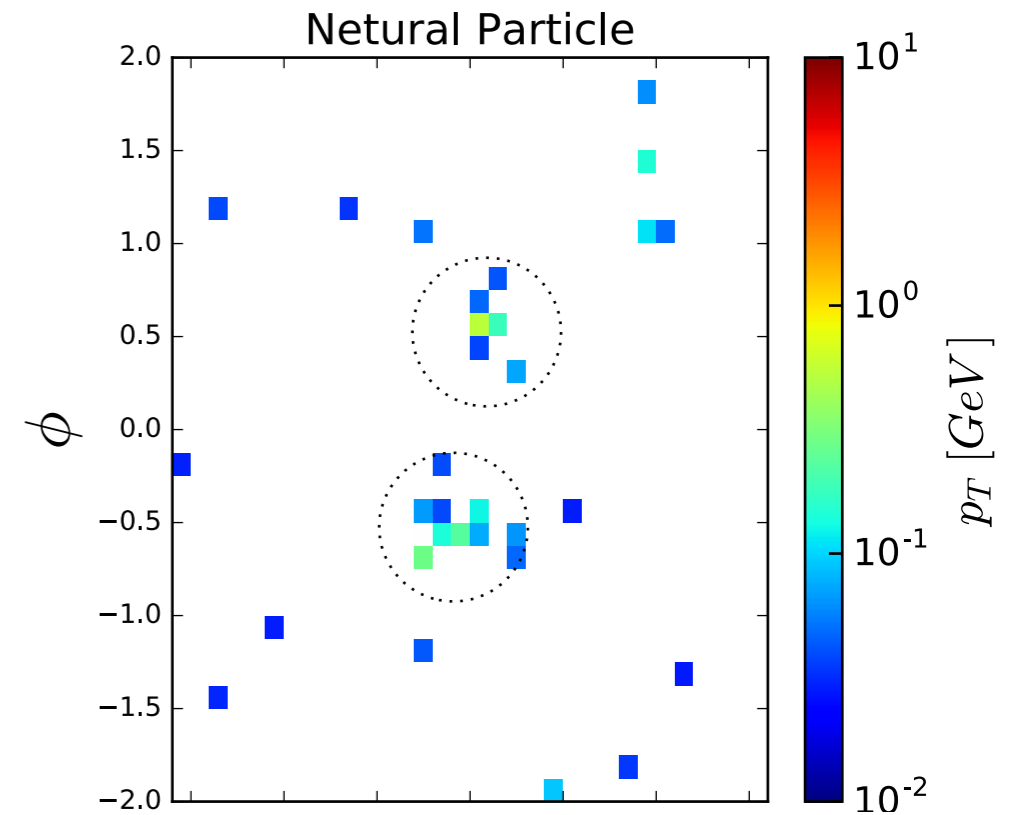
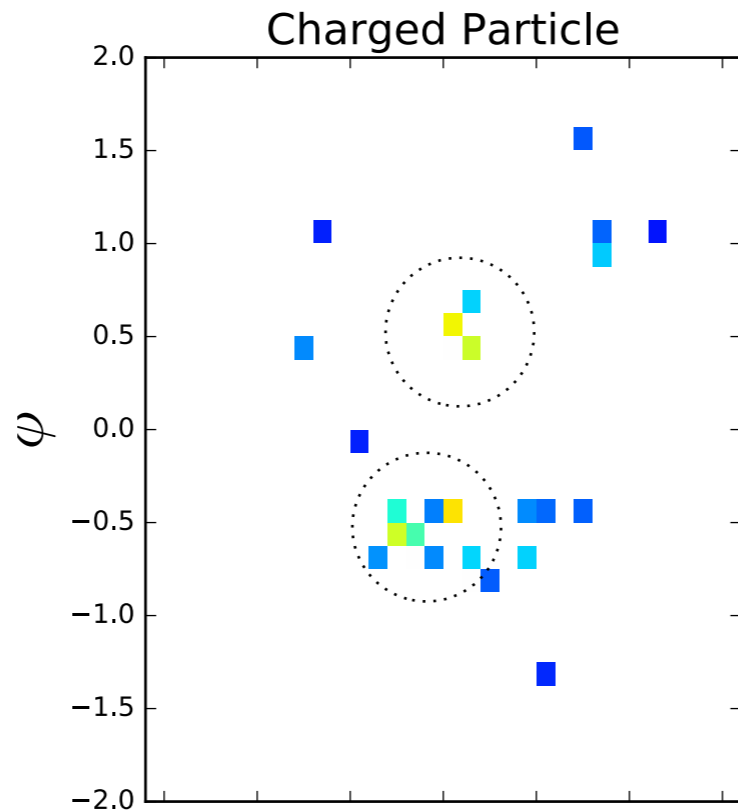
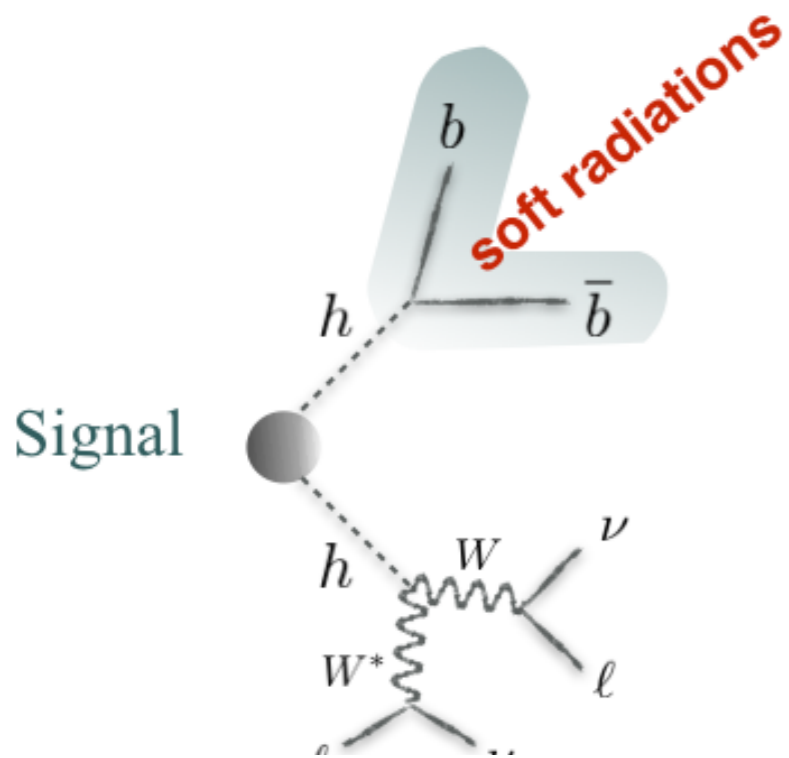
Can we do
more than this?

- Consider "**orthogonal**" method to kinematics; **Color-flow**

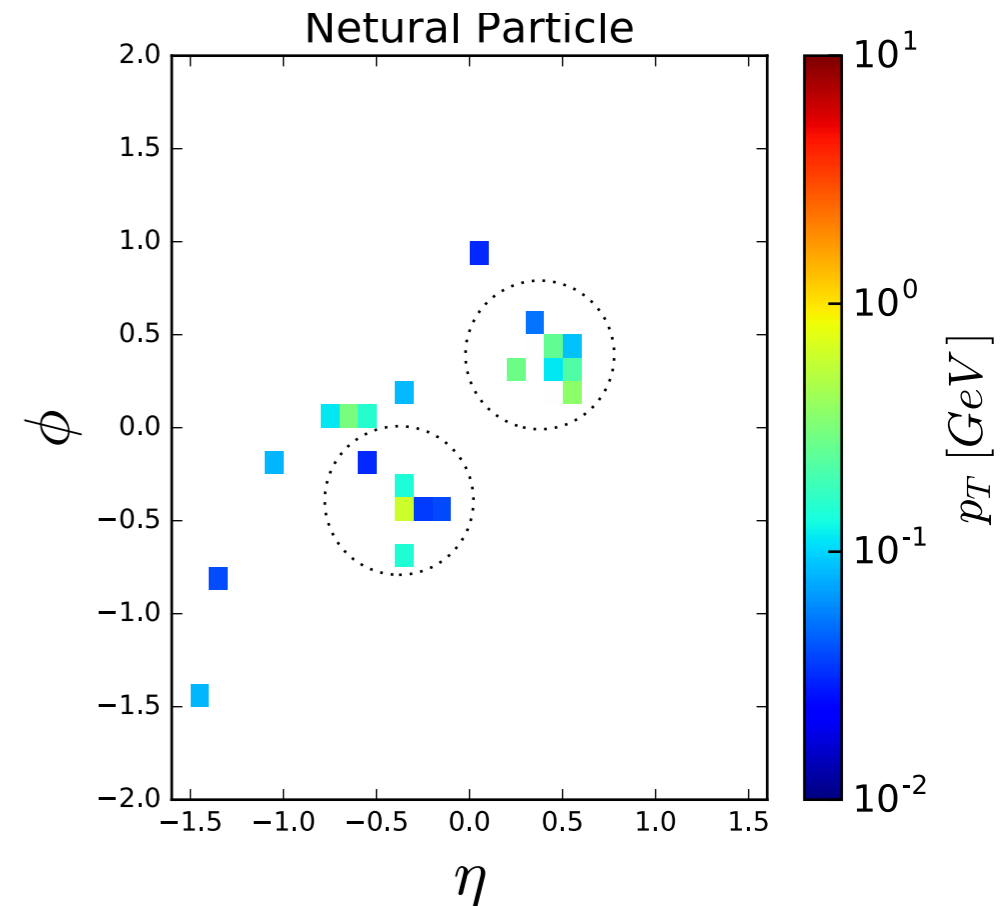
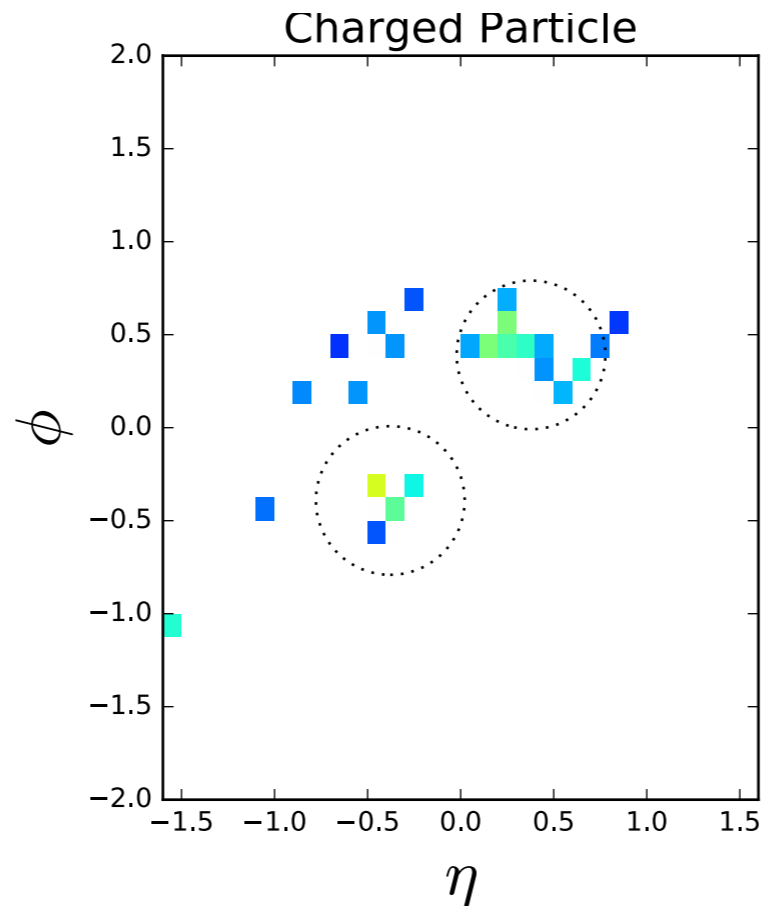
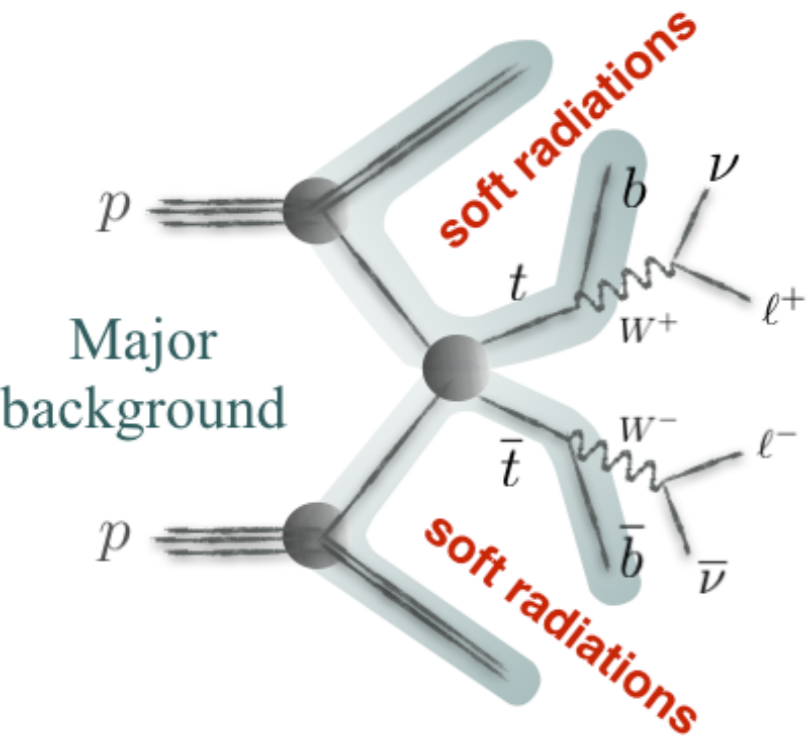


Energy deposits

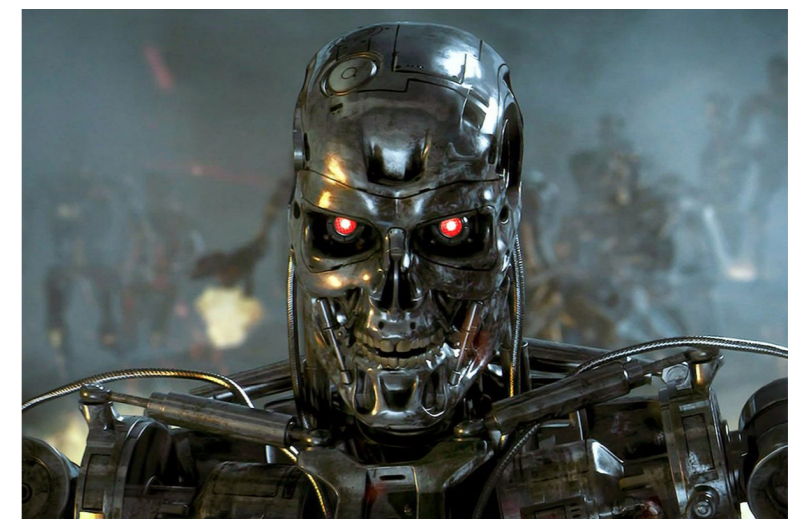




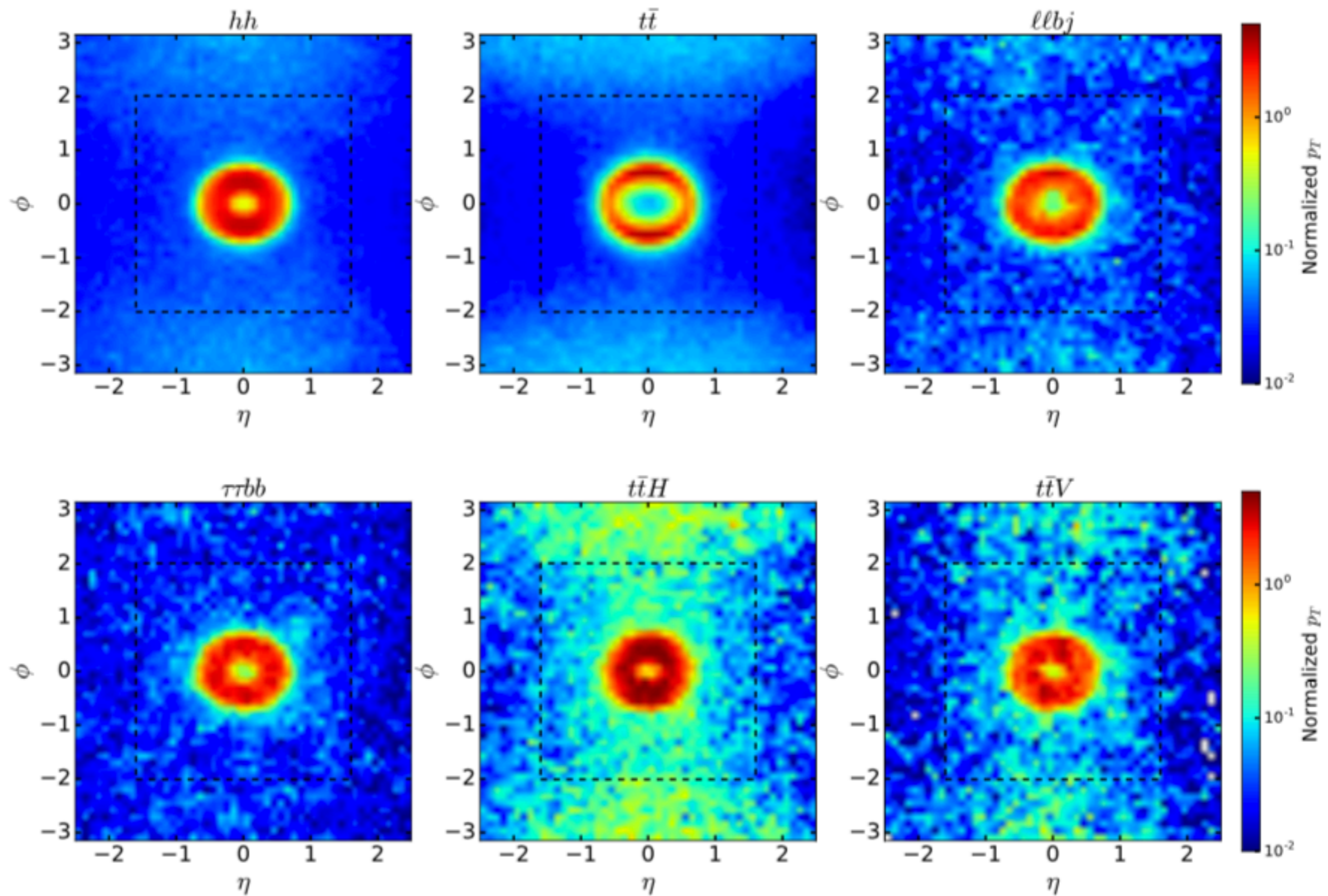
Can you see a pattern?



Let's leave simple and mundane jobs to a machine



- Whole events with $b\bar{b}$ (fat-jet) axis as a center

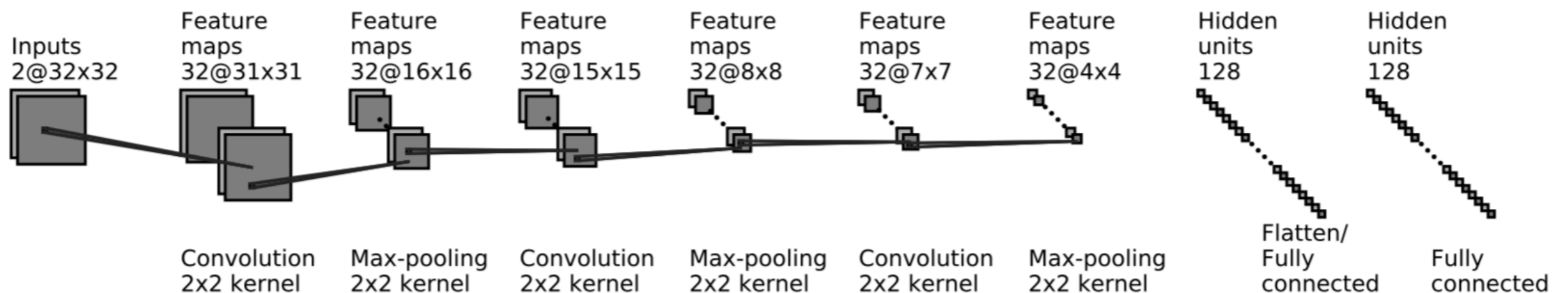


Maximizing information with Deep Neural Network

- For **jet-image**, we use 32x32 pixels for $-2.5 < \eta < 2.5$, $-\pi < \phi < \pi$.

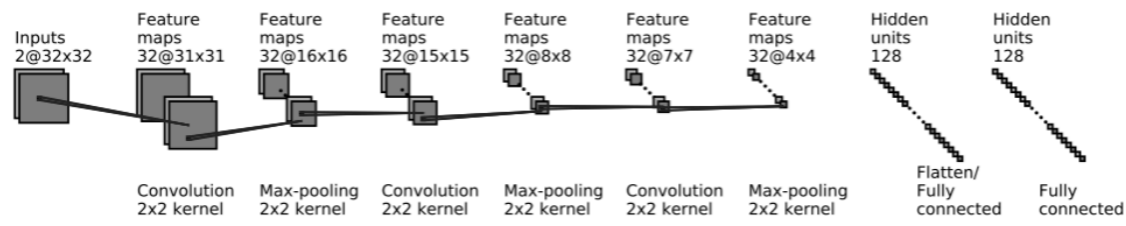
Input channels for **CNN** are divided into two with particle flow:

- Neutral particles
- Charged particles

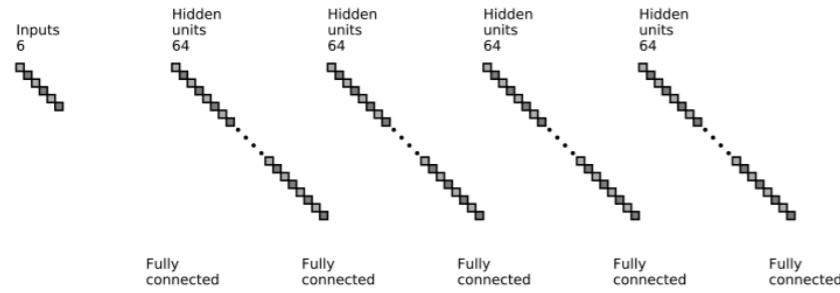


For technical details about DNN, please contact **Minho Kim**
(kmhmon@postech.ac.kr)

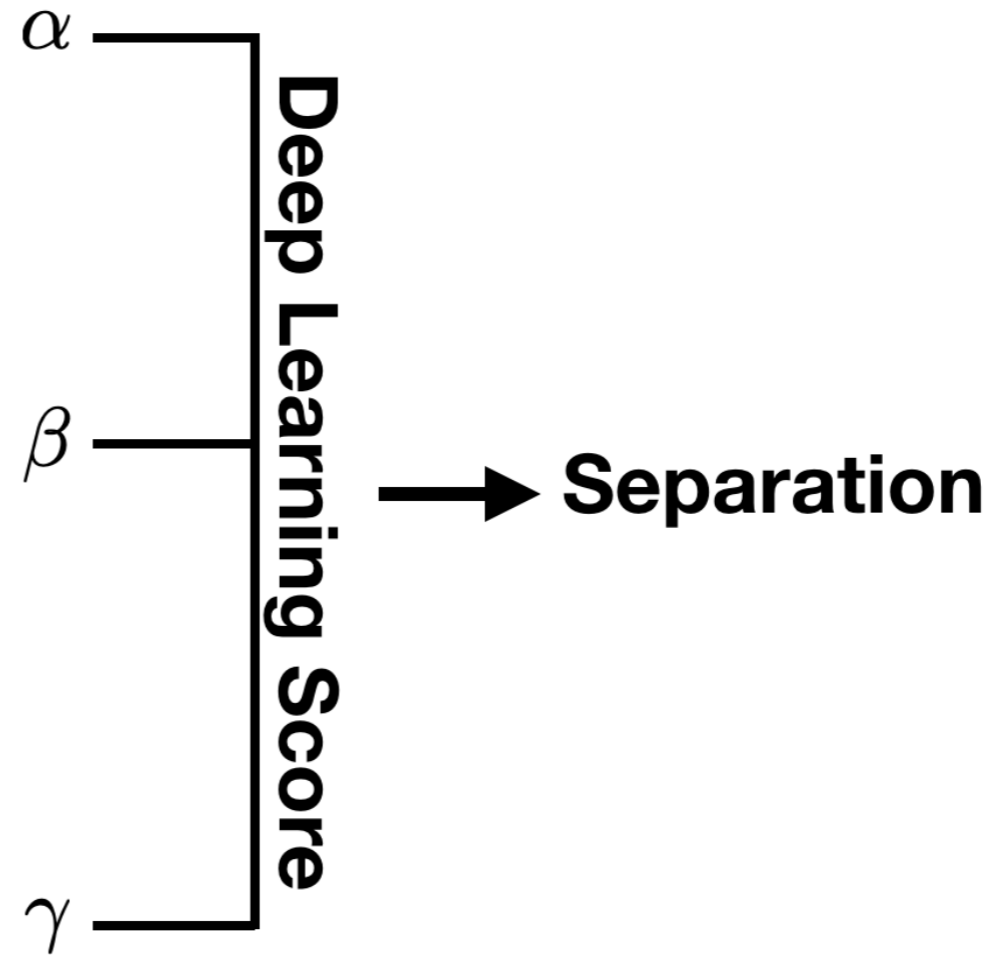
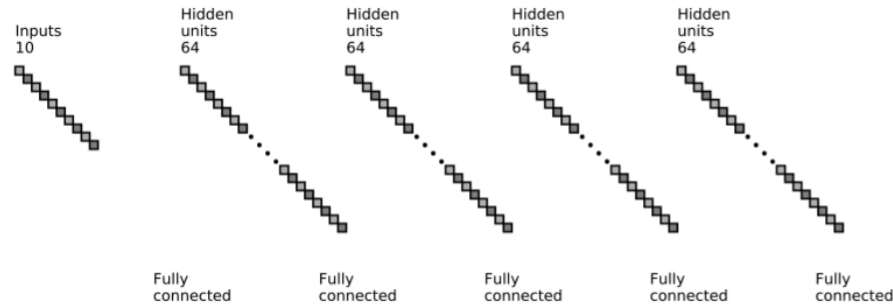
QCD observable



High level Kinematics-observables



Low level Kinematics-observables

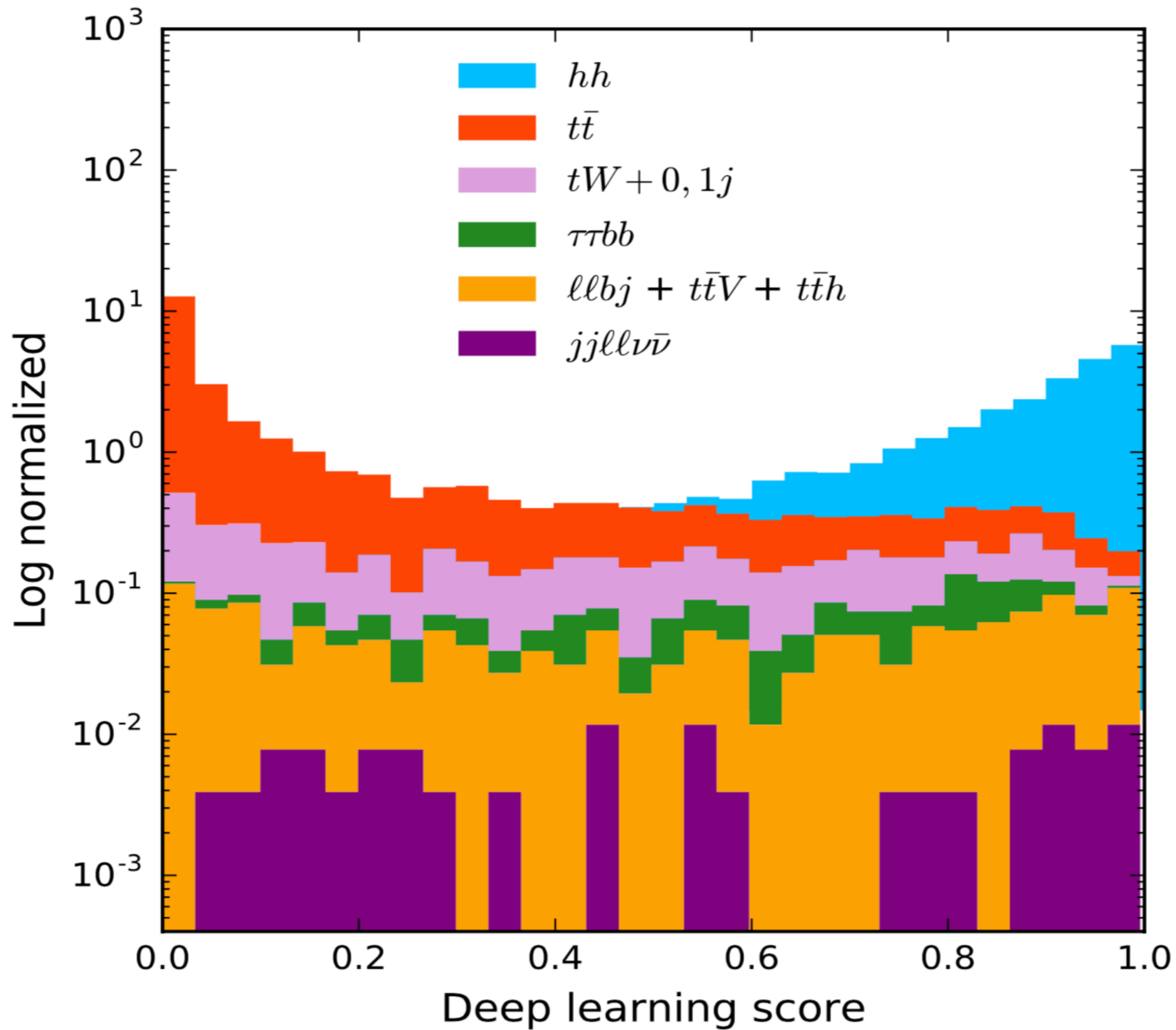


6 High Level Variables

Input data: $\sqrt{\hat{s}}_{\min}^{(\bar{b}, b, \bar{\ell}, \ell)}$ $\sqrt{\hat{s}}_{\min}^{(\bar{\ell}, \ell)}$ M_{T2}^b M_{T2}^{ℓ} Higgsness, Topness

10 Low Level Variables

Input data: MET $p_{\bar{\ell}}^t$ p_{ℓ}^t $\Delta R_{(\bar{\ell}, \ell)}$ $M_{(\bar{b}, b)}$ $p_{(\bar{b}, b)}^t$ $\Delta R_{(\bar{b}, b)}$ $M_{(\bar{\ell}, \ell)}$ $p_{(\bar{\ell}, \ell)}^t$, $\Delta\phi_{\{(\bar{\ell}, \ell), (\bar{b}, b)\}}$



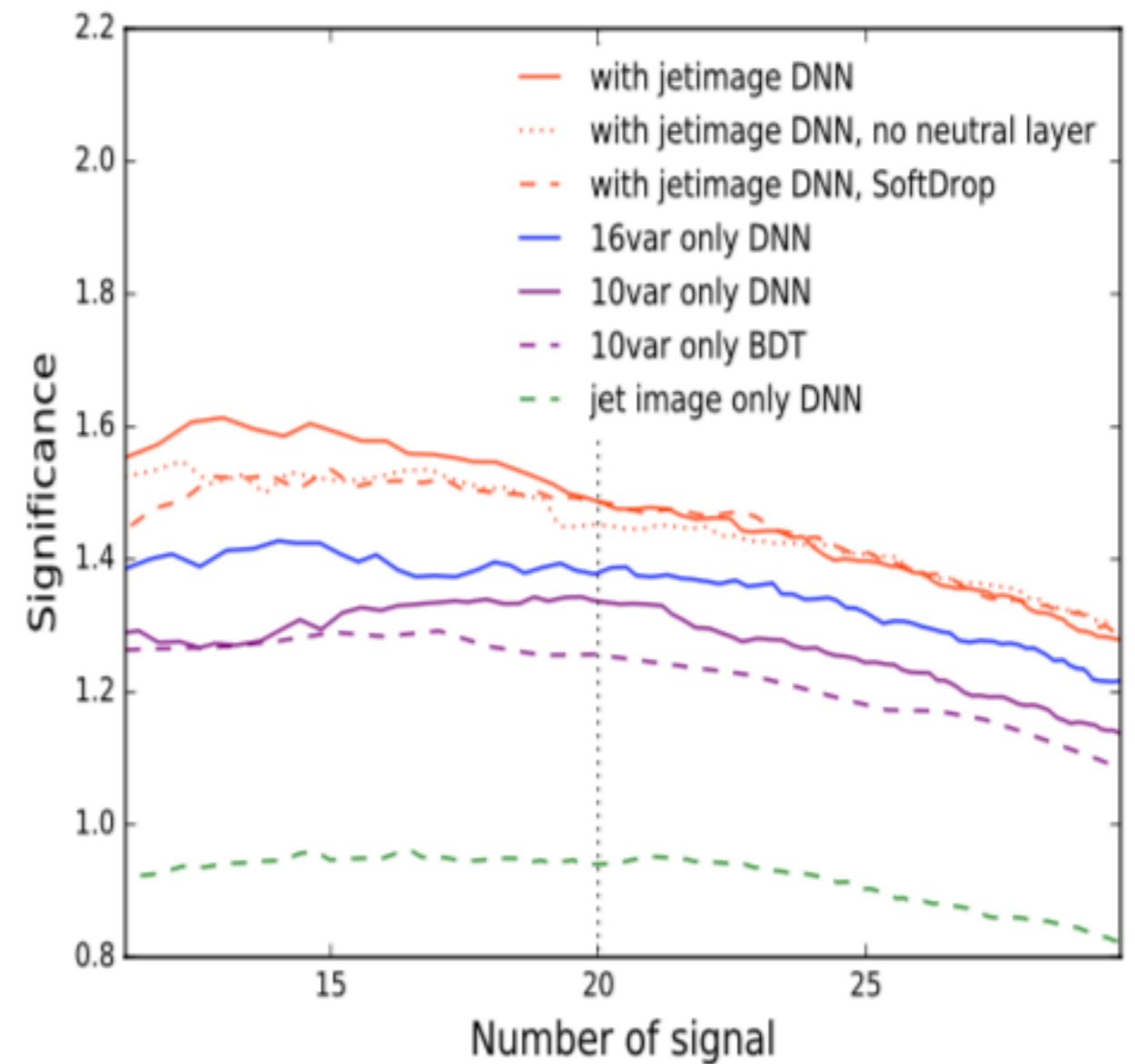
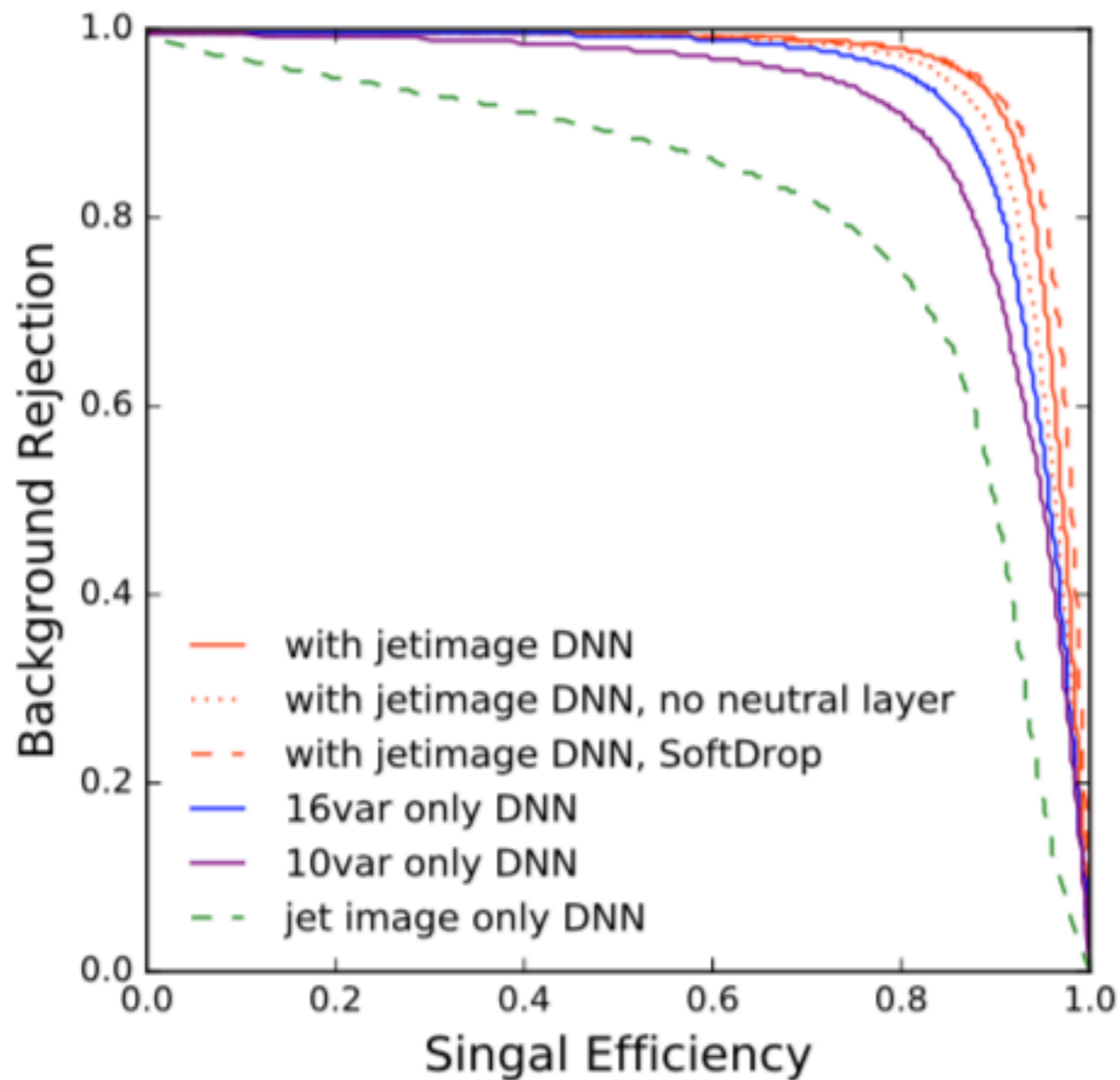
"Backgrounds as stacked Histogram"

- To estimate effects from **pileup** removal,

0. No additional processes.

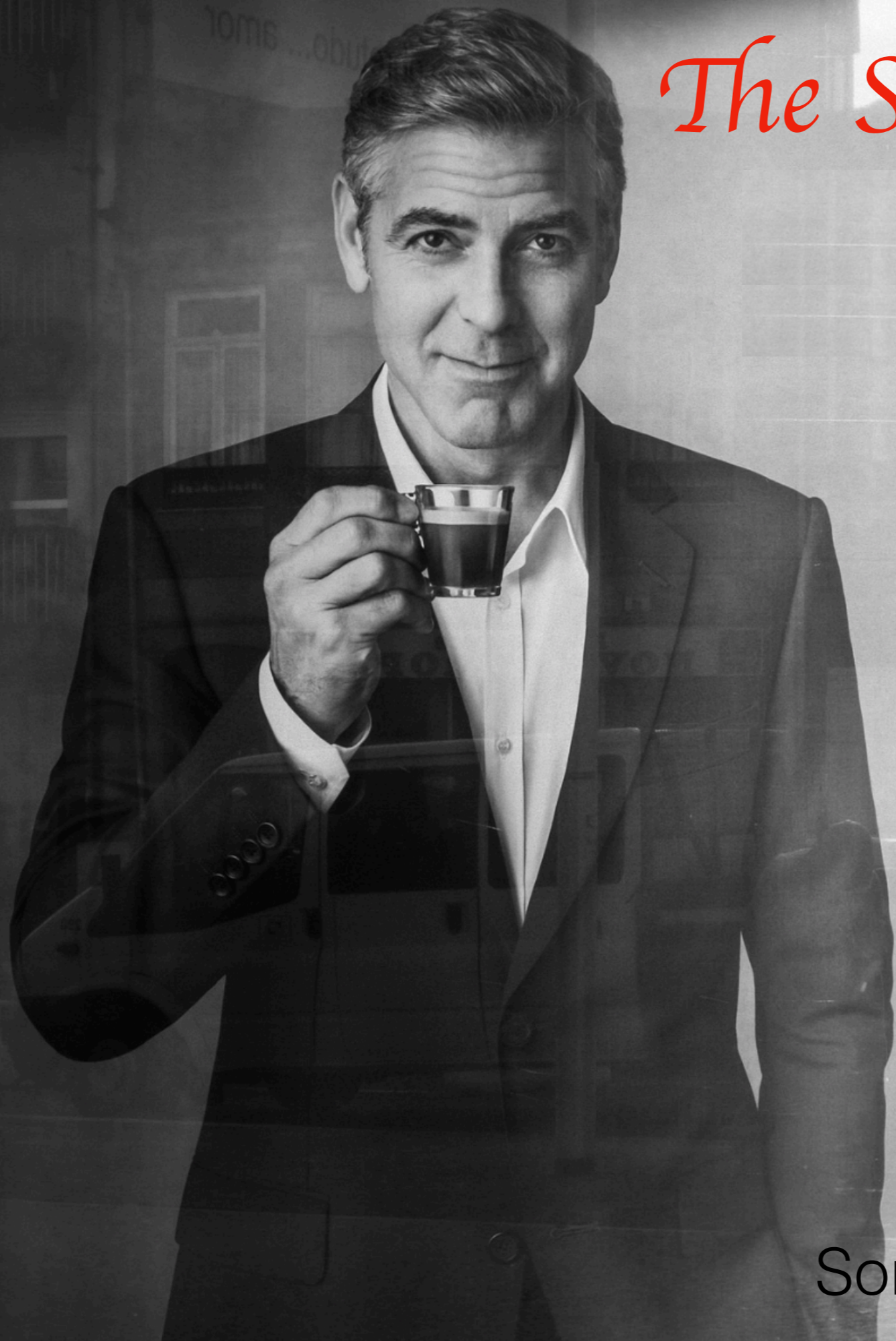
1. we apply SoftDrop to a fat-jet ($R= 1.2$ anti-Kt)

2. we use only charged layer



Conclusions

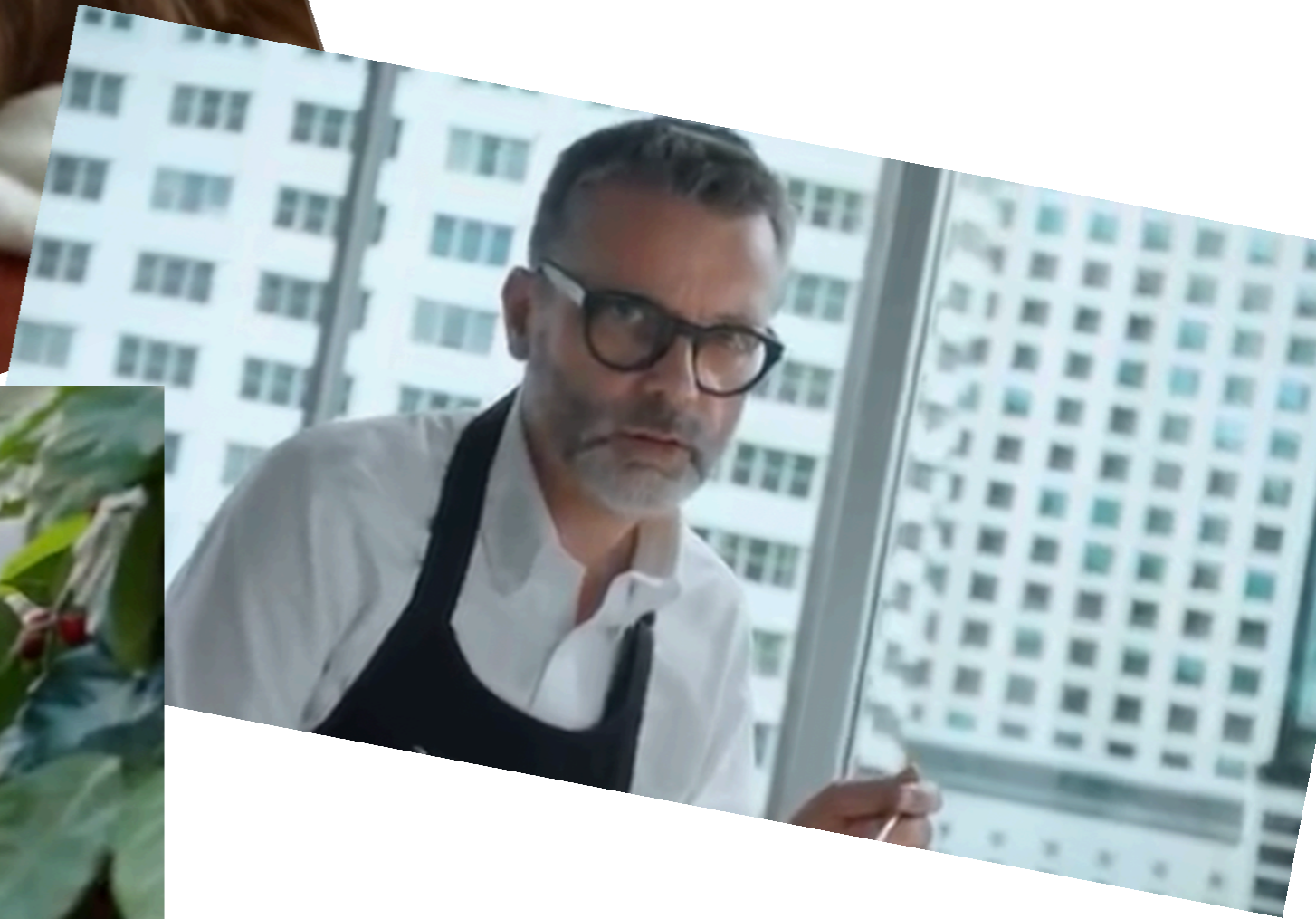
- For HL-LHC, Offshell Higgs can be a good probe for BSM.
- High Precision and maximization of **background rejection** are required.
- For a new computational tool, there are things that we need to look into and "**design**" analysis methods.
 - High level kinematical variables is introduced.



*The Standard Model ,
What else ?*

Sorry George for a copy right !

Really George?



Precision to a New Physics!!!